

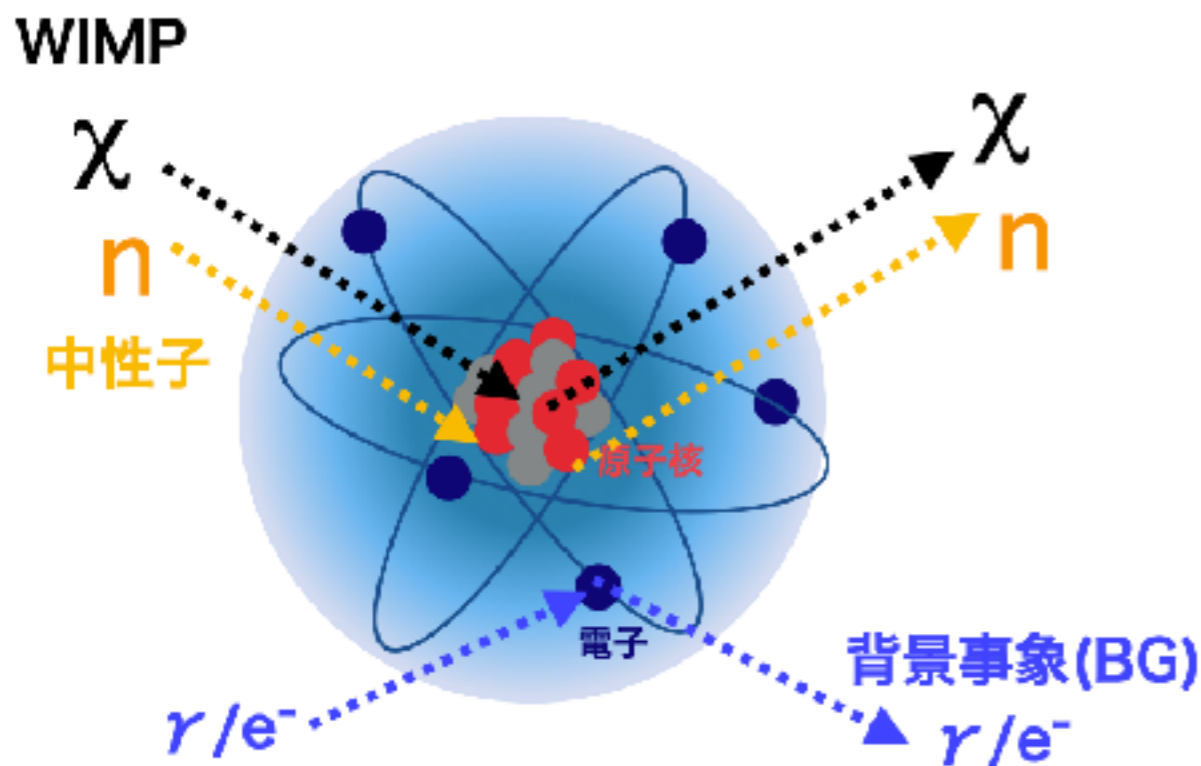
# 電子散乱事象の超過 @XENON1T実験

風間慎吾 (名古屋大学 KMI & 高等研究院)

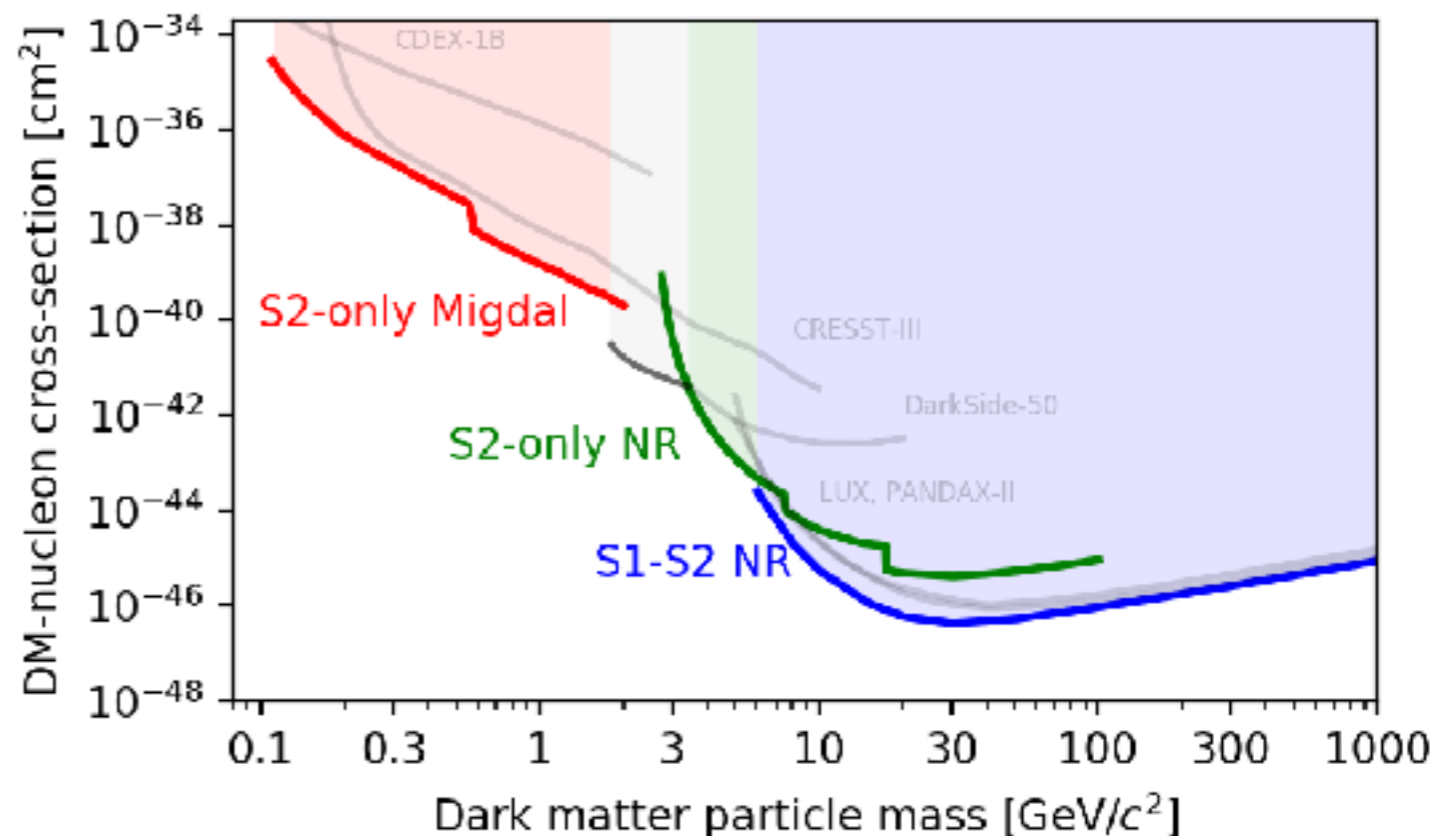
@ダークマター懇談会2020

## WIMPの探索方法(直接探索)

確率は非常に小さいがWIMPも身の回りの物質(原子核)と相互作用をする(原子核反跳)。  
原子核が受け取る反跳エネルギーを検出する(光 & 電子 for XENON1T実験)。



## XENON1T実験のWIMP探索結果



## XENON1T実験

- ・ 液体キセノンを用いた直接探索実験 (3.2トン(有効体積~1トン))
- ・ 低質量&高質量の両極限(100MeV - TeV)で、世界で最も厳しい制限を与えている。
- ・ 実験自体は既に終了していて、現在XENONnT実験へとアップグレード中(小林くんのトーク)

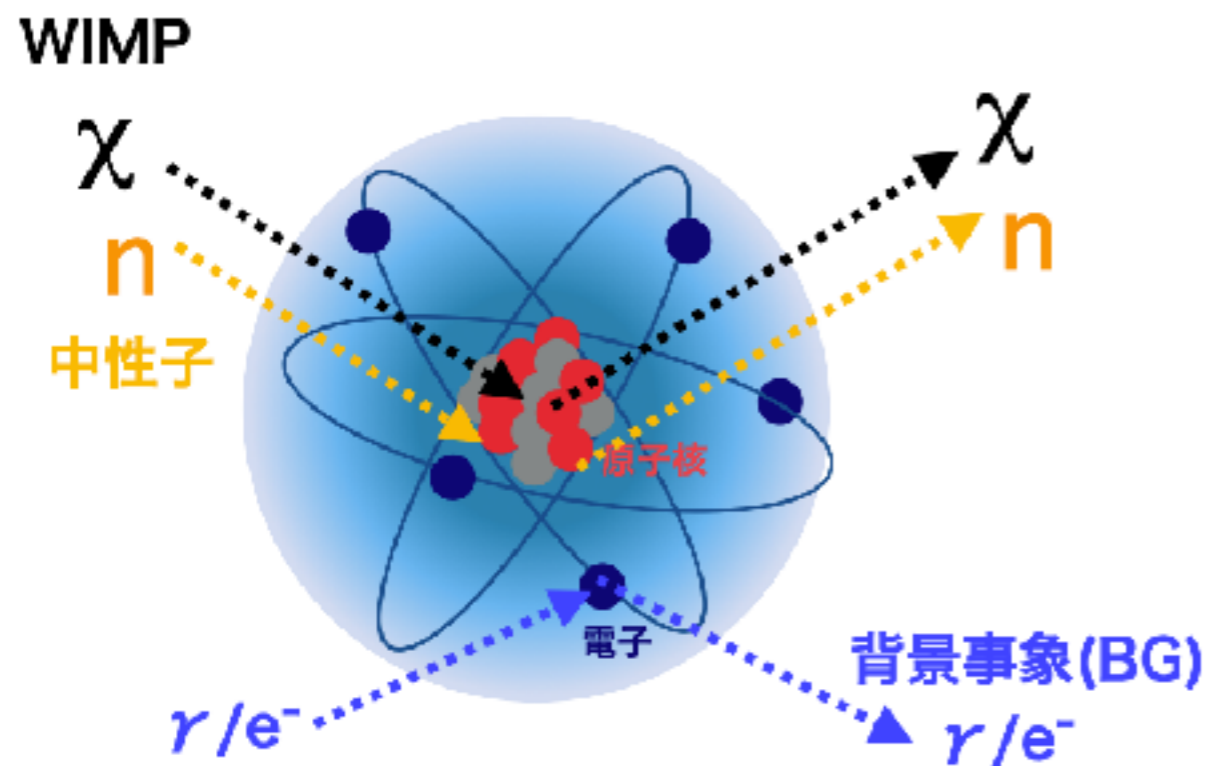
## 太陽アクシオンや太陽ニュートリノの探索方法: 電子反跳

太陽アクシオン: Axio-electric effect (光電効果と似た効果,  $g_{ae}$ )

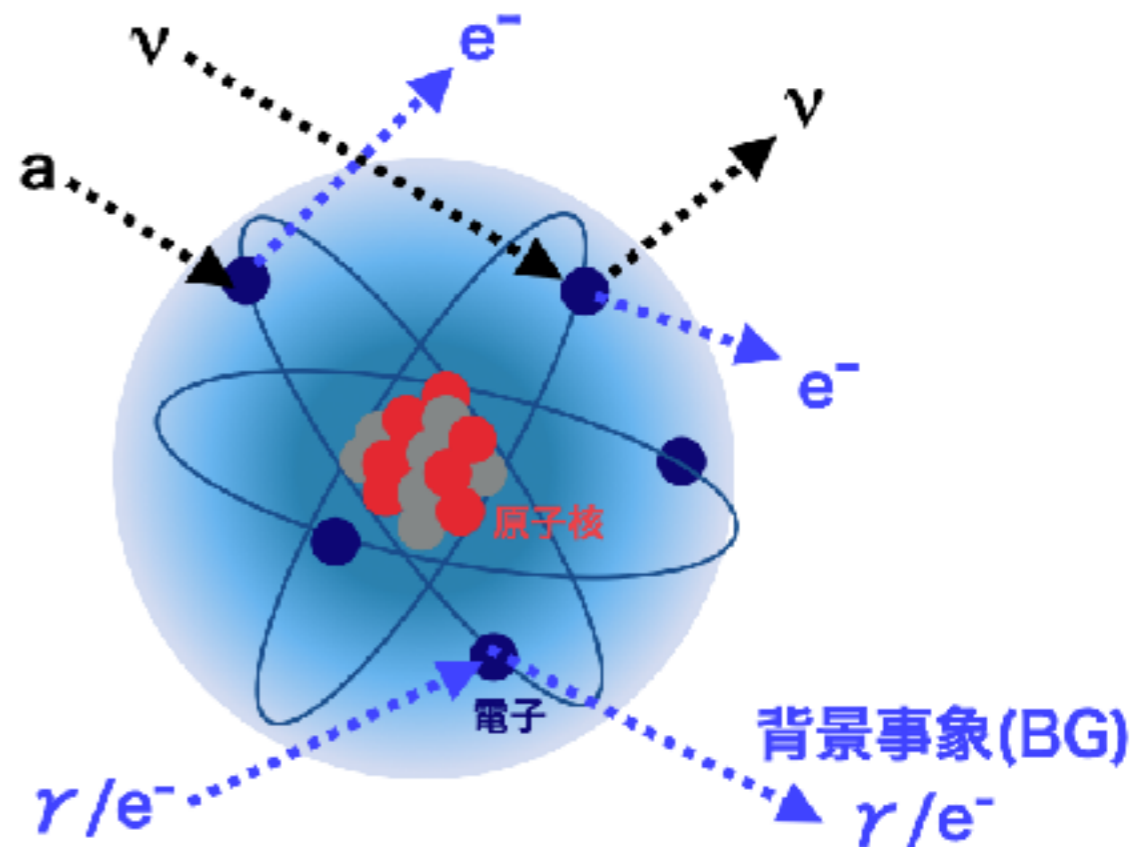
太陽ニュートリノ: 電子散乱 (elastic scattering)



## WIMPと原子核の相互作用 (原子核反跳)



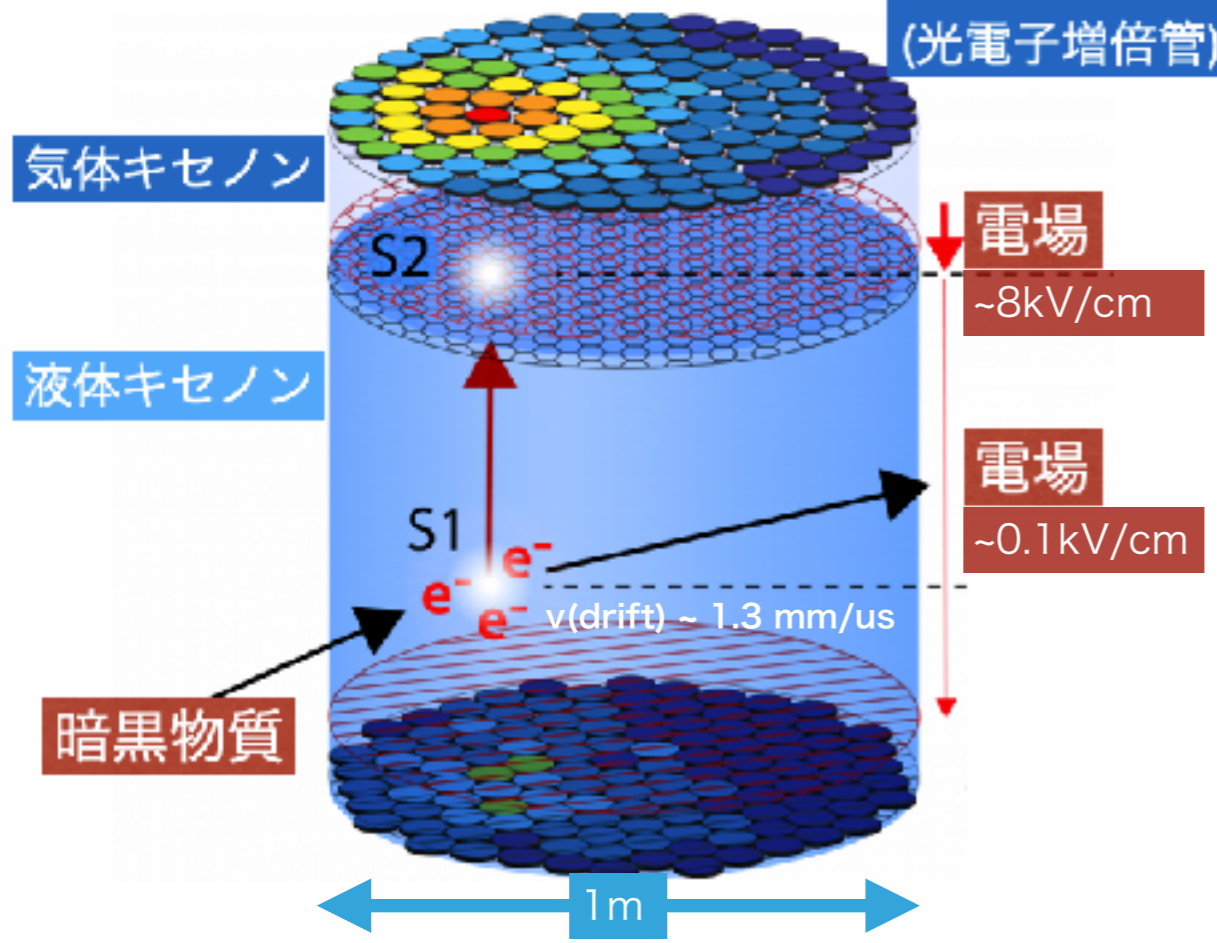
## 太陽アクシオン・太陽ニュートリノ (今回はこちら!)



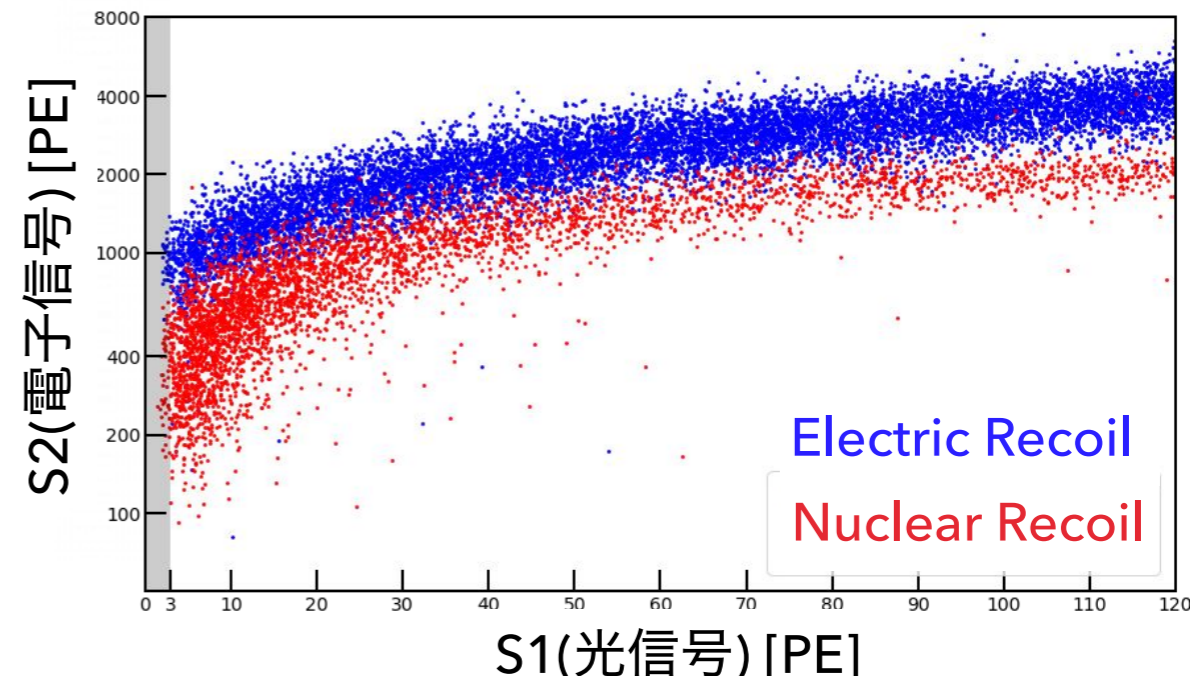
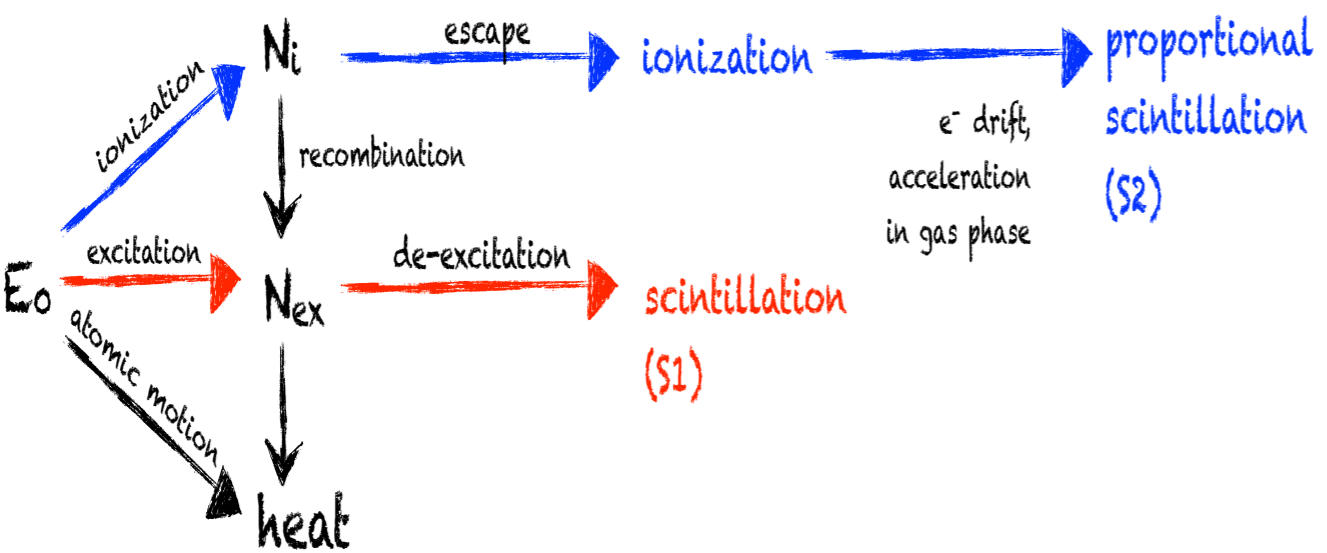
## 電子反跳事象の探索

- ・ 通常、電子反跳事象はWIMP探索の背景事象(BG)
- ・ WIMP searchと比べてBG量(ex: 放射性ラドン・クリプトンなど)が多いので、BGをより精密に評価し、そこからの超過を探す

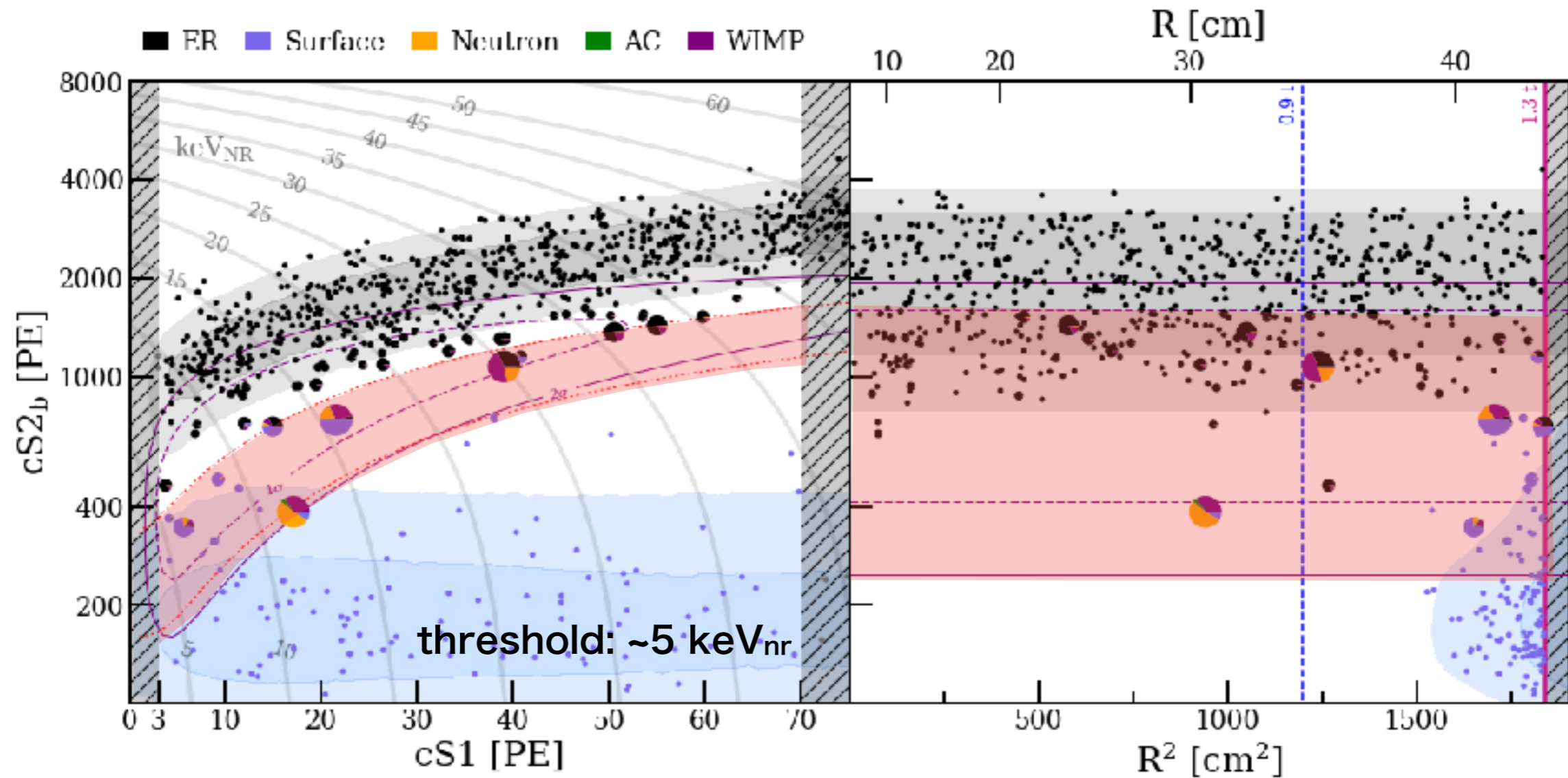
## 液体キセノン検出器



- Primary scintillation light (**S1**) is produced promptly at the interaction site
- Ionization electrons drift up through the LXe in the applied electric field
- Some recombine with ions → more scintillation light (**S1**)
- Others are extracted above the liquid surface into gas phase region, where they form secondary proportional scintillation light (**S2**)
- Event vertex reconstruction in 3D space
  - X,Y position: **S2 hit-pattern in top PMT array**
  - Z position: **electron drift time,  $\Delta t$  (s1, s2)**
- Particle type discrimination:  **$(S2/S1)_{\gamma,e} > (S2/S1)_{WIMP}$**   
     Electric Recoil                      Nuclear Recoil

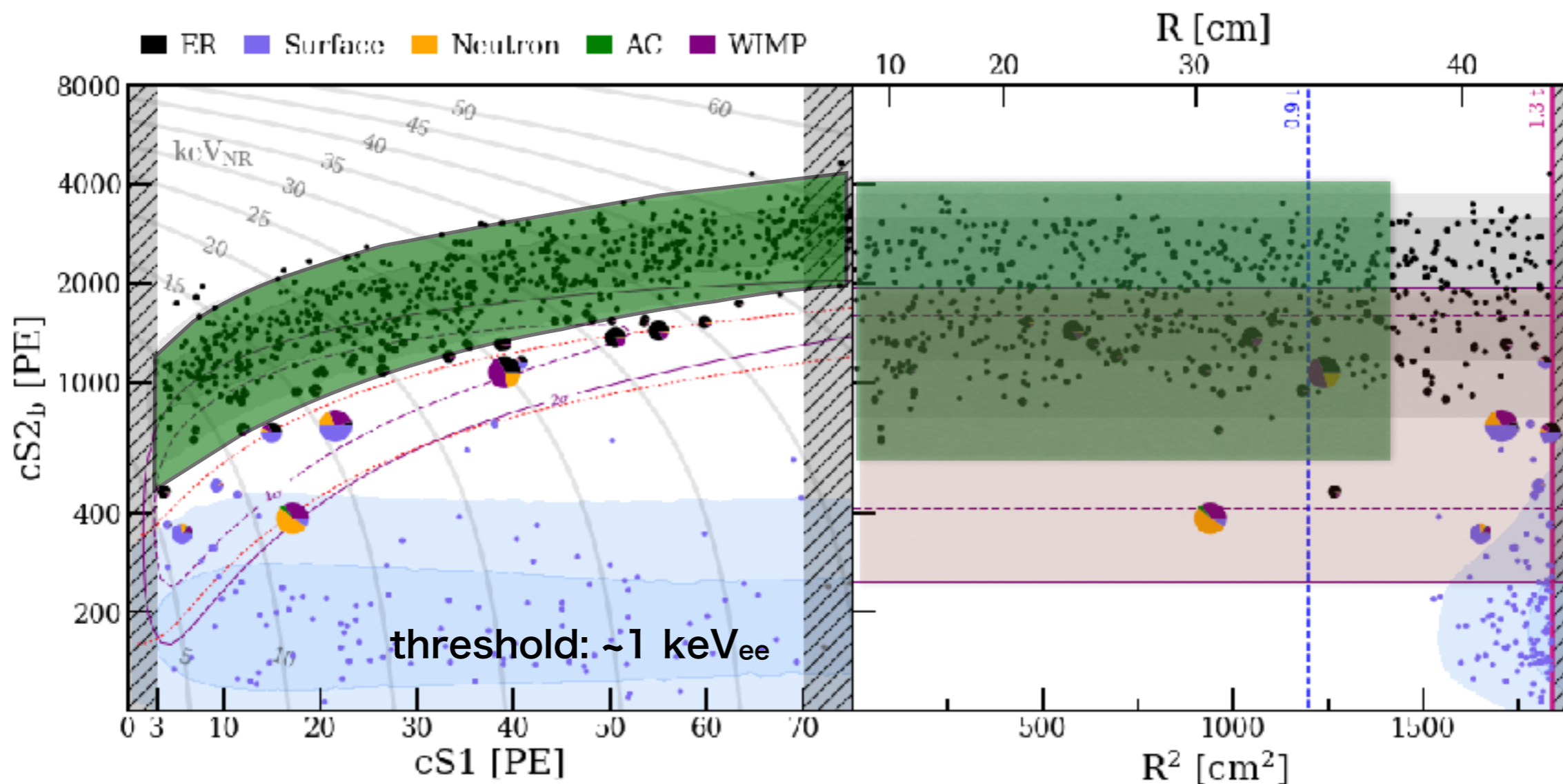


One ton-year of search for WIMPs induced nuclear recoils



Most stringent result on WIMP Dark Matter down to  $3 \text{ GeV}/c^2$  masses

検出器部材からの放射線(ガンマ線)を除くため、有効体積はWIMPより小さい。



太陽アクシオンやALPs, Dark Photon探索における戦略

- 電子反跳BGの絶対量を減らす
- 既知のBG(放射性ラドン・クリプトンなど)を精密に評価し、超過を探す

# XENON1T Detector

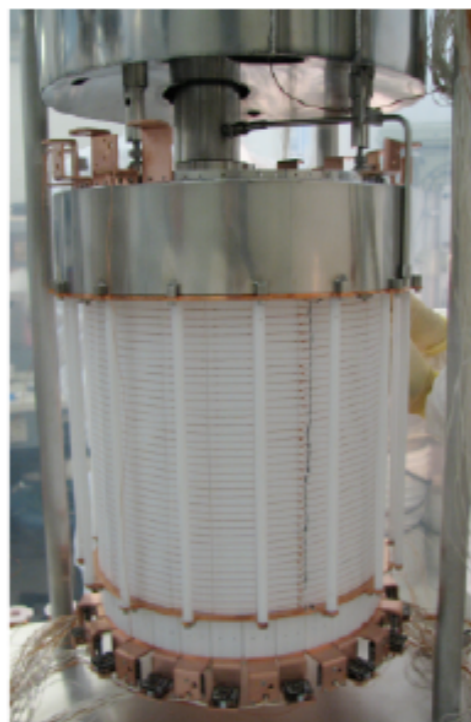


# The XENON + DARWIN Program

XENON10



XENON100



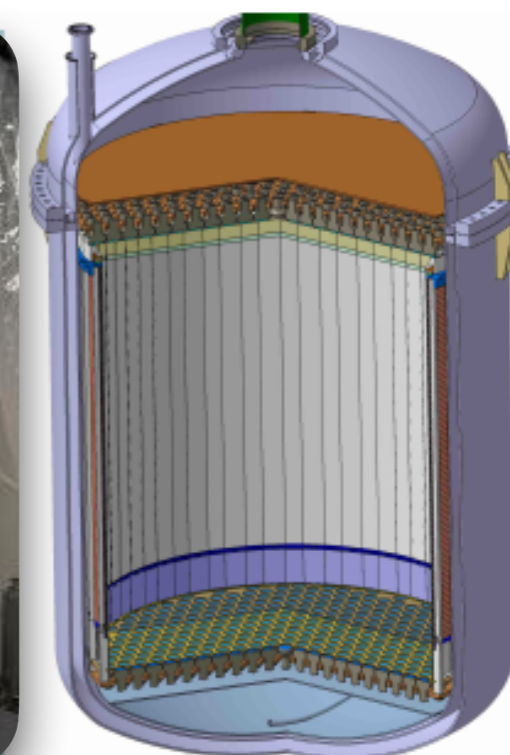
XENON1T



XENONnT



DARWIN



2005-2007

2008-2016

2012-2018

2019-2023

2020+

15 kg

161 kg

3200 kg

8200 kg

50 tonnes

$\sim 10^{-43} \text{ cm}^2$

$\sim 10^{-45} \text{ cm}^2$

$\sim 10^{-47} \text{ cm}^2$

$\sim 10^{-48} \text{ cm}^2$

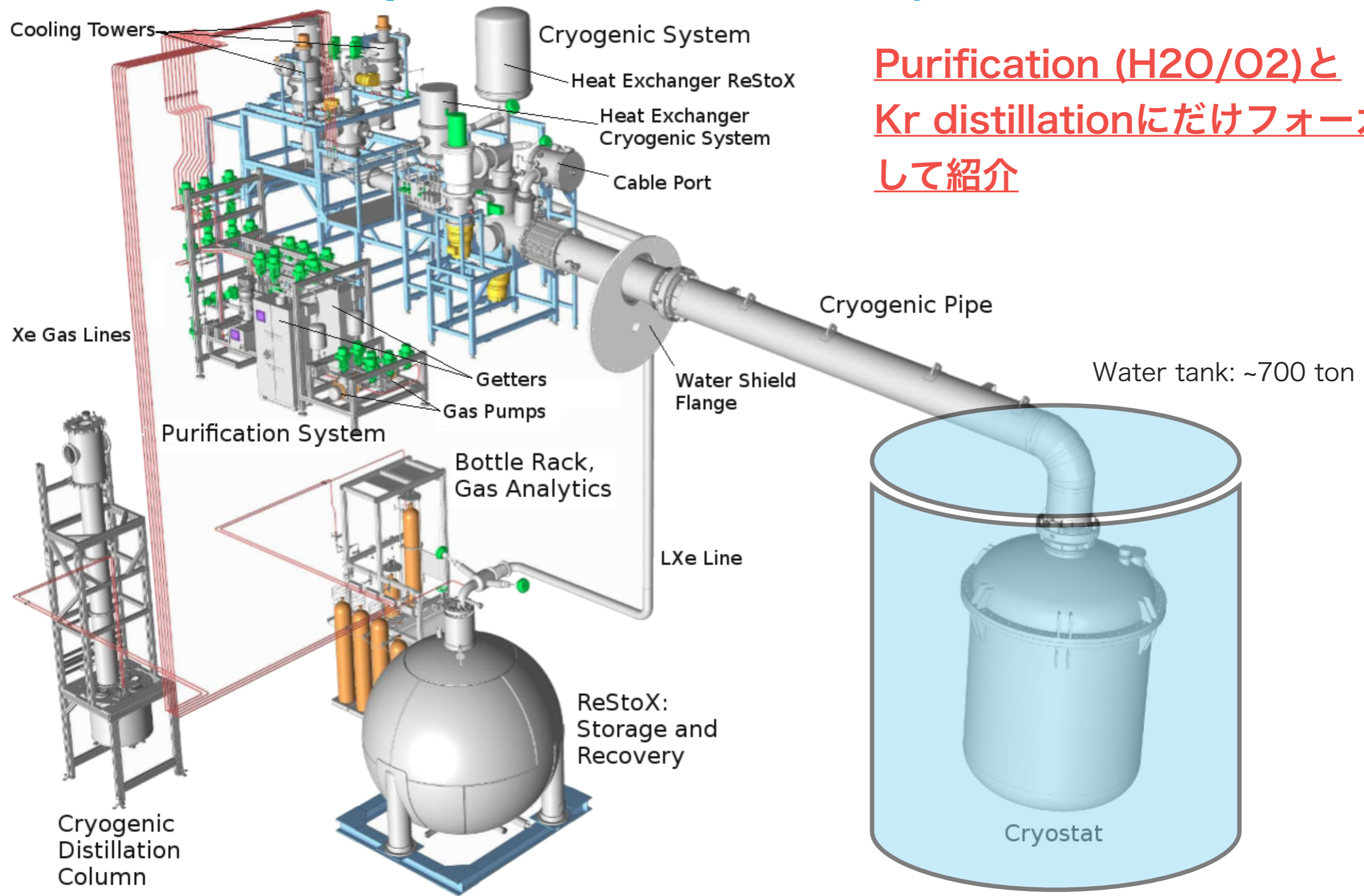
$\sim 10^{-49} \text{ cm}^2$

先週ちょうど液体キセノンをfillし始めた！

名大、神戸大、IPMUは先月末にDARWINにも参加！



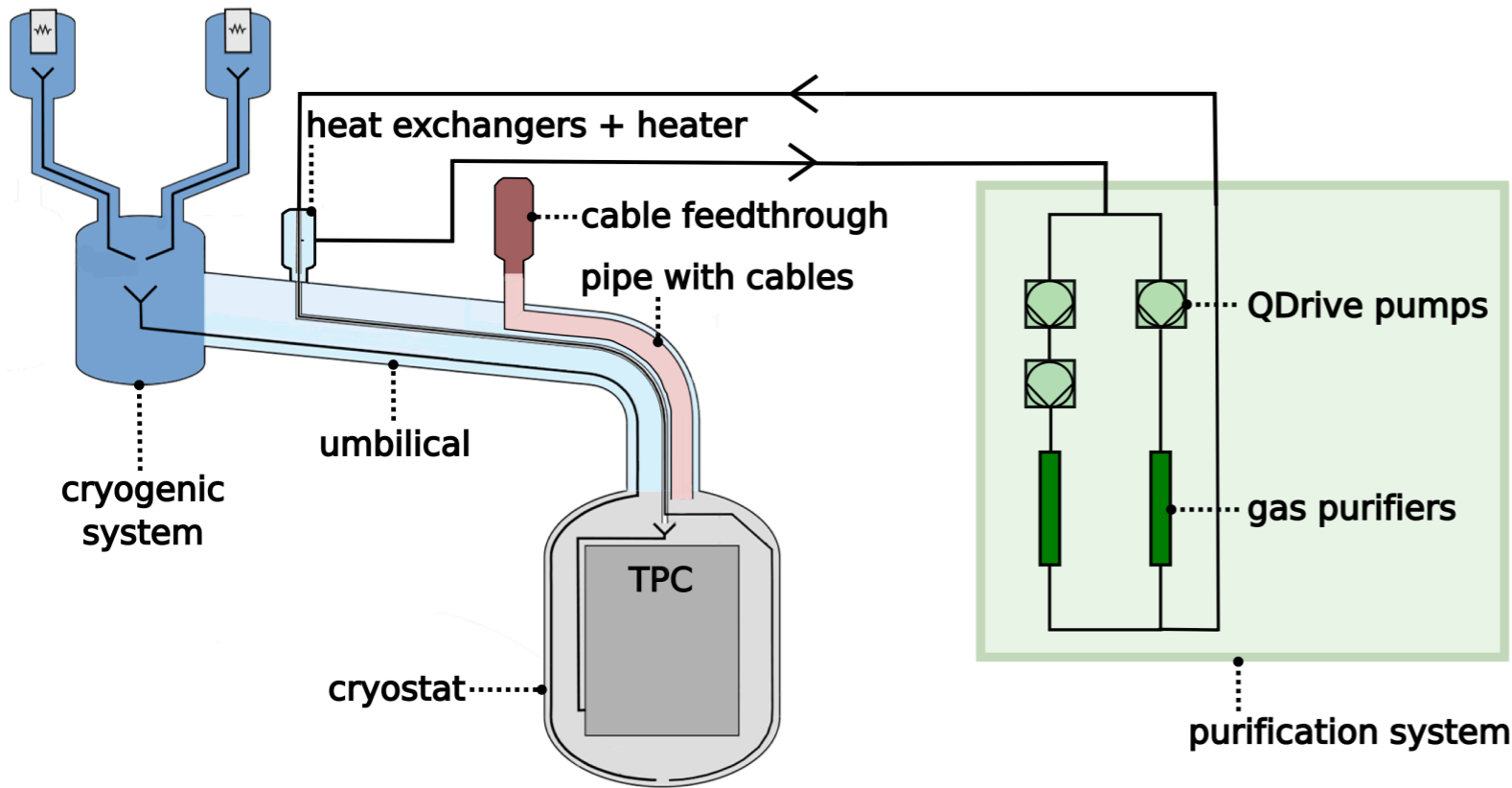
# The XENON1T Experiment @ LNGS in Italy



Purification (H<sub>2</sub>O/O<sub>2</sub>)と  
Kr distillationにだけフォーカス  
して紹介

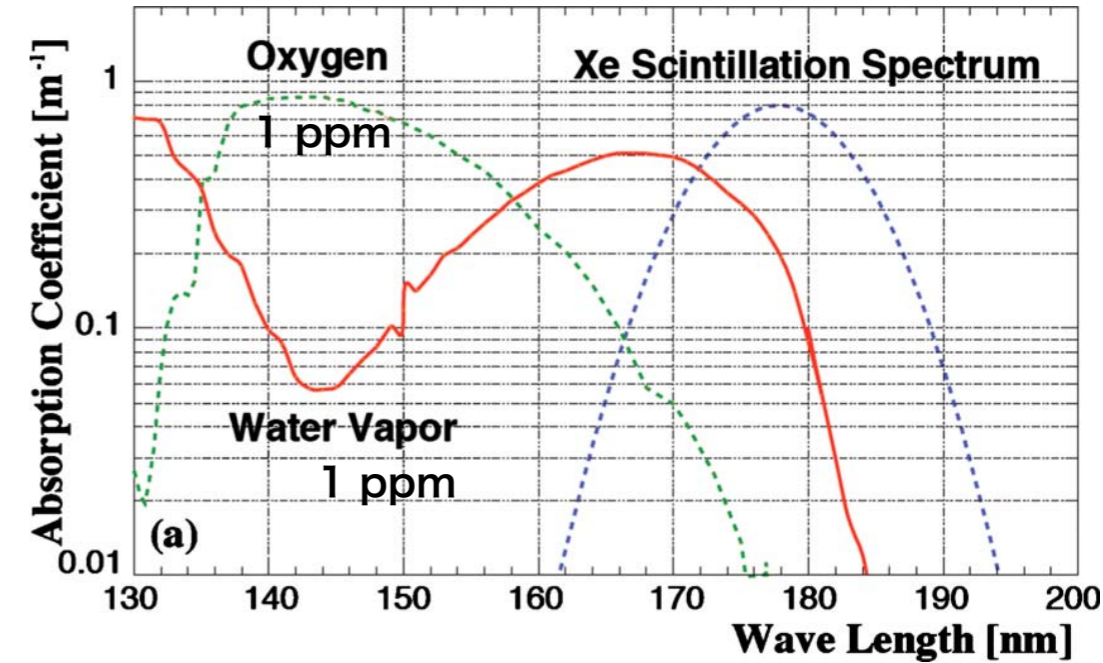
248 3-inch PMTs (R11410-21, QE~34%@178nm)

LXe mass: 3.2 t(total), 2.0t (active)



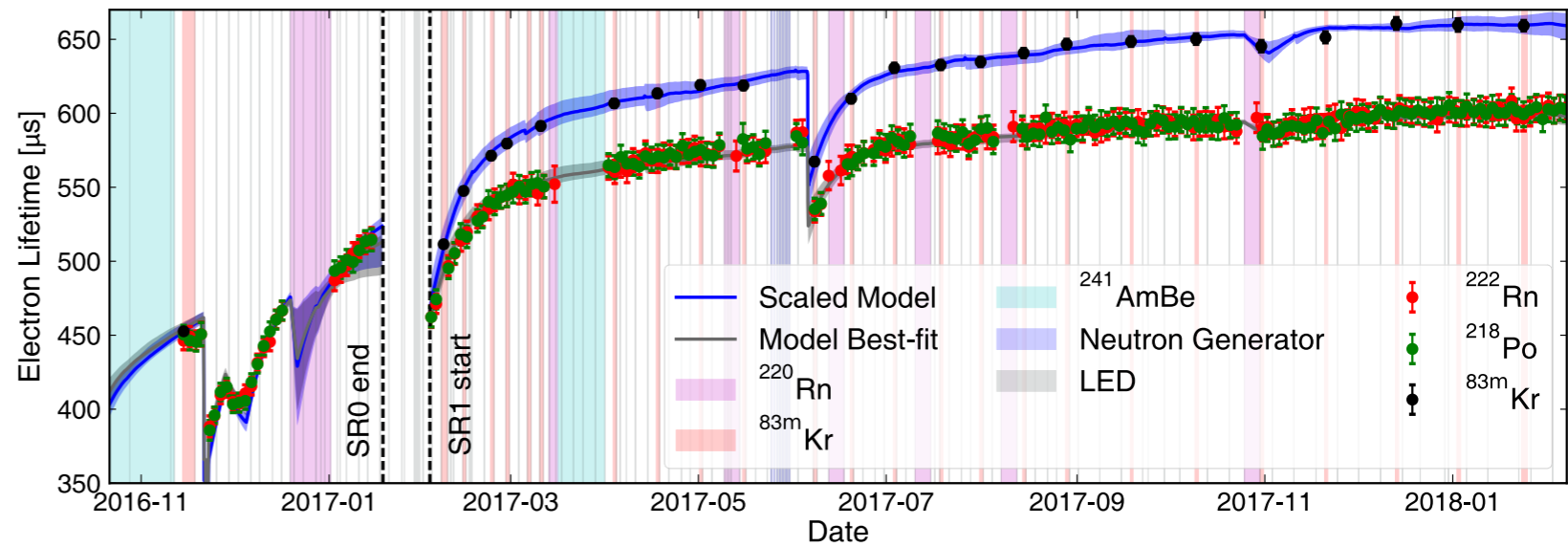
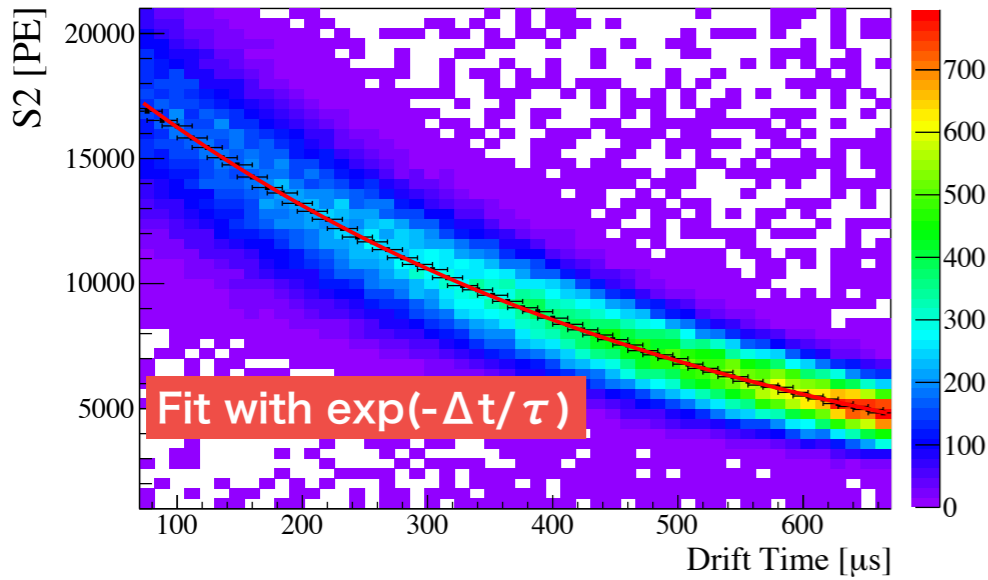
## S1光の吸収

Ozone, 2005



- ▶ Electronegative impurities in the Xe gas and from outgassing reduce both photon (S1) and electron (S2) signals.
- ▶ To detect light S1 signals efficiently, need  $O(1)$  ppb  $H_2O$  concentration. (Tiger Optics HALO+  $H_2O$  monitor used for purity measurement)
- ▶ To drift electrons over 1 meter requires  $< 1$  ppb ( $O_2$  equivalent)
- ▶ Solution: continuous gas circulation at high flow through heated getter material (Zr-V-Fe getter,  $400^\circ C$ ) with hydrogen removal unit
- ▶ SAES PS4-MT50-R ( $O_2$ ,  $H_2O$ , CO,  $CO_2$ ,  $H_2$ ,  $N_2$ ,  $CH_4$ : ~ppb)
- ▶ Total flow rate of 54 slpm driven by up to 3 pumps.



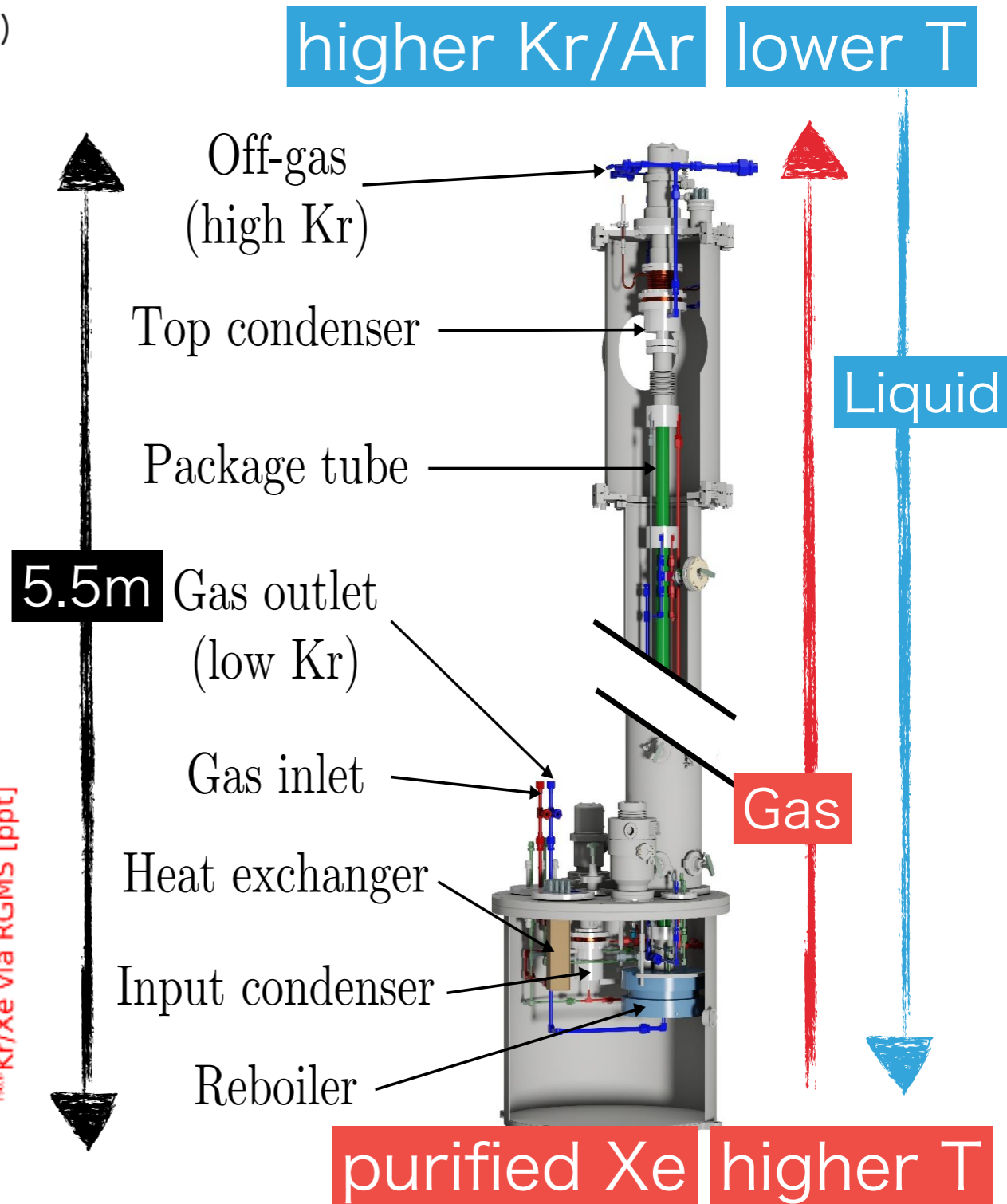
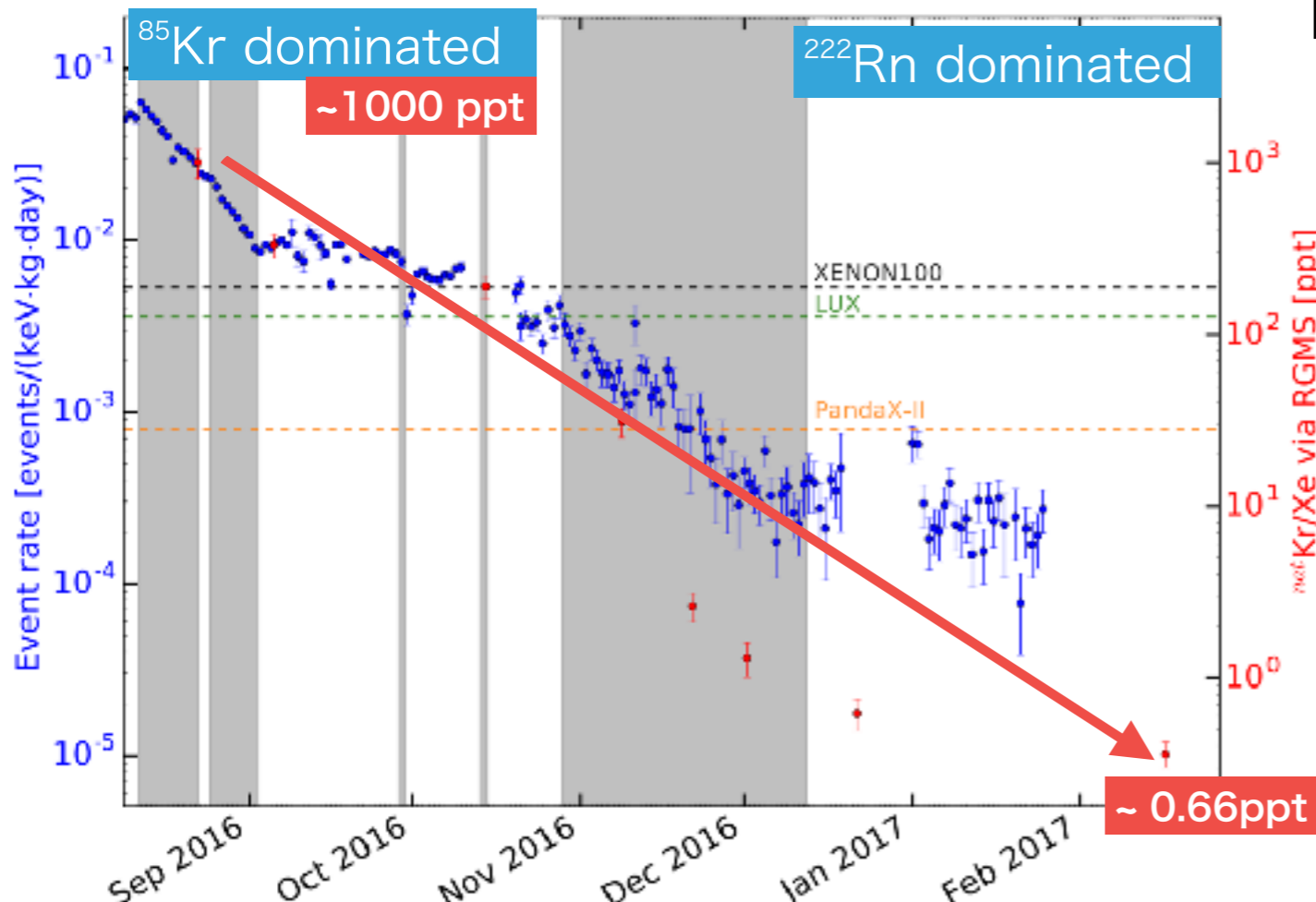


- ▶ Solution: continuous gas circulation at high flow through heated getter material (Zirconium)
- ▶ SAES PS4-MT50-R (O<sub>2</sub>, H<sub>2</sub>O, CO, CO<sub>2</sub>, H<sub>2</sub>, N<sub>2</sub>, CH<sub>4</sub>: ~ppb)
- ▶ Total flow rate of 54 slpm driven by up to 3 pumps.
- ▶ 650us of e-lifetime → oxygen equivalent impurity concentration of ~ 0.5ppb.
- ▶ XENON1T TPC length ~ 650us drift time



# Kr (and Ar) Distillation

- ▶ Commercial Xe: 1 ppm - 10 ppb <sup>nat</sup>Kr,
- ▶ <sup>85</sup>Kr is unstable ( $T_{1/2} = 10.8$  y,  $\beta$ -decay with Q-value = 687 keV)
- ▶ Solution: 5.5 m cryogenic distillation column
- ▶ Utilizes different vapor pressure:
  - Kr: 20900 mbar@178K, Xe: 2010 mbar@178K
- ▶ Feeding flow rate: 8.3 SLPM (3kg/h)
  - Thermodynamically stable up to 18 SLPM (6.5 kg/h)
- ▶ Measured separation:  $6.4 \times 10^5$  @8.3 SLPM, < 48 ppq (RGMS)



# Energy Reconstruction

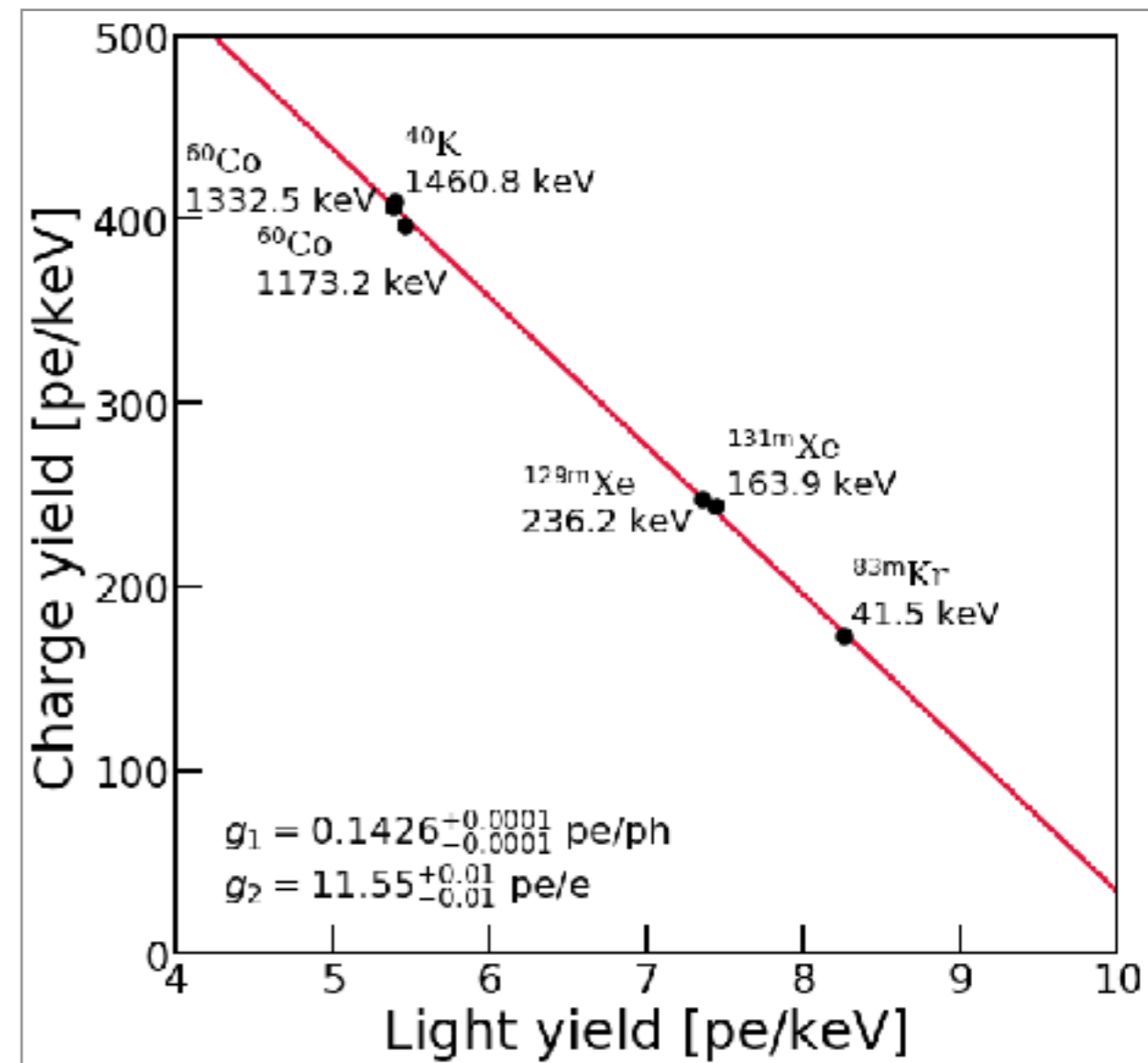
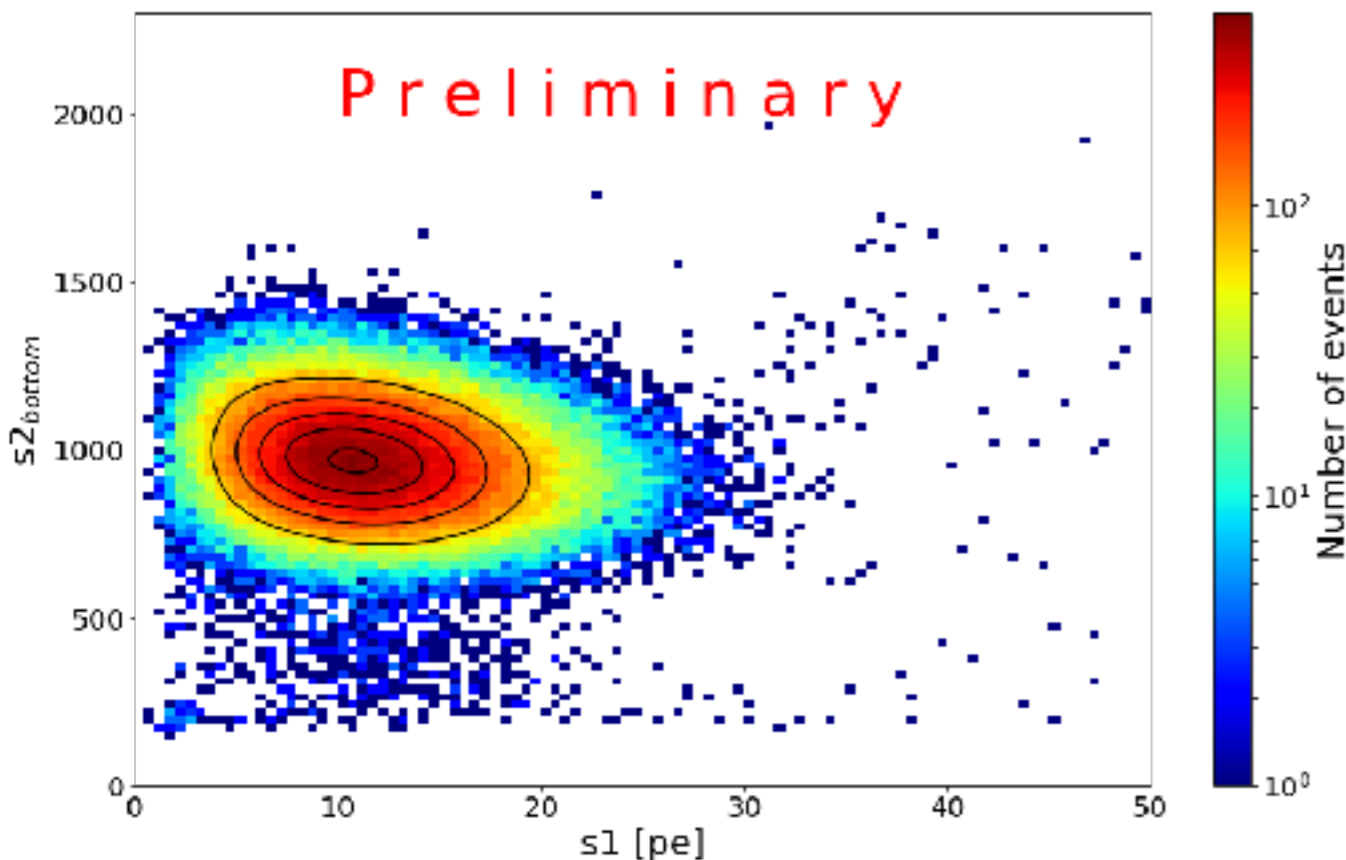
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$$E = (N_{ph} + N_e) \cdot W = \left( \frac{S1}{g1} + \frac{S2}{g2} \right) \cdot W$$

where  $W = 13.7$  eV/quanta

g1 and g2: detector-specific gain constants  
 extract g1/g2 from calibration data, use it to  
 reconstruct energy of each event

S1, S2はPMTで測定される量!

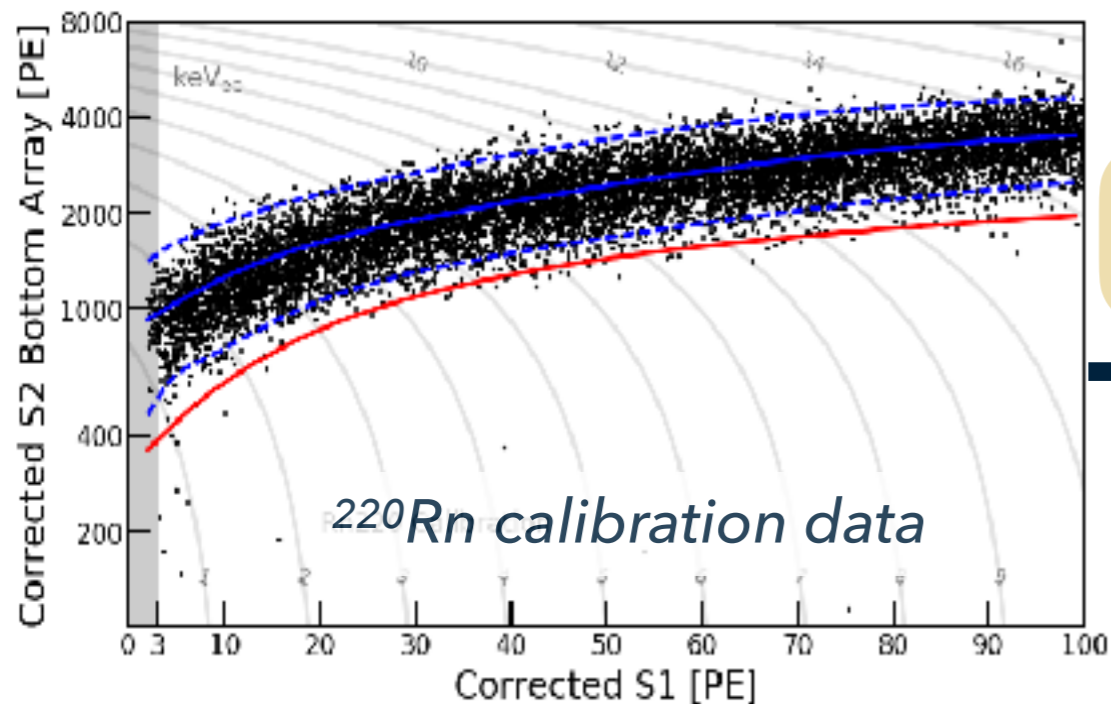
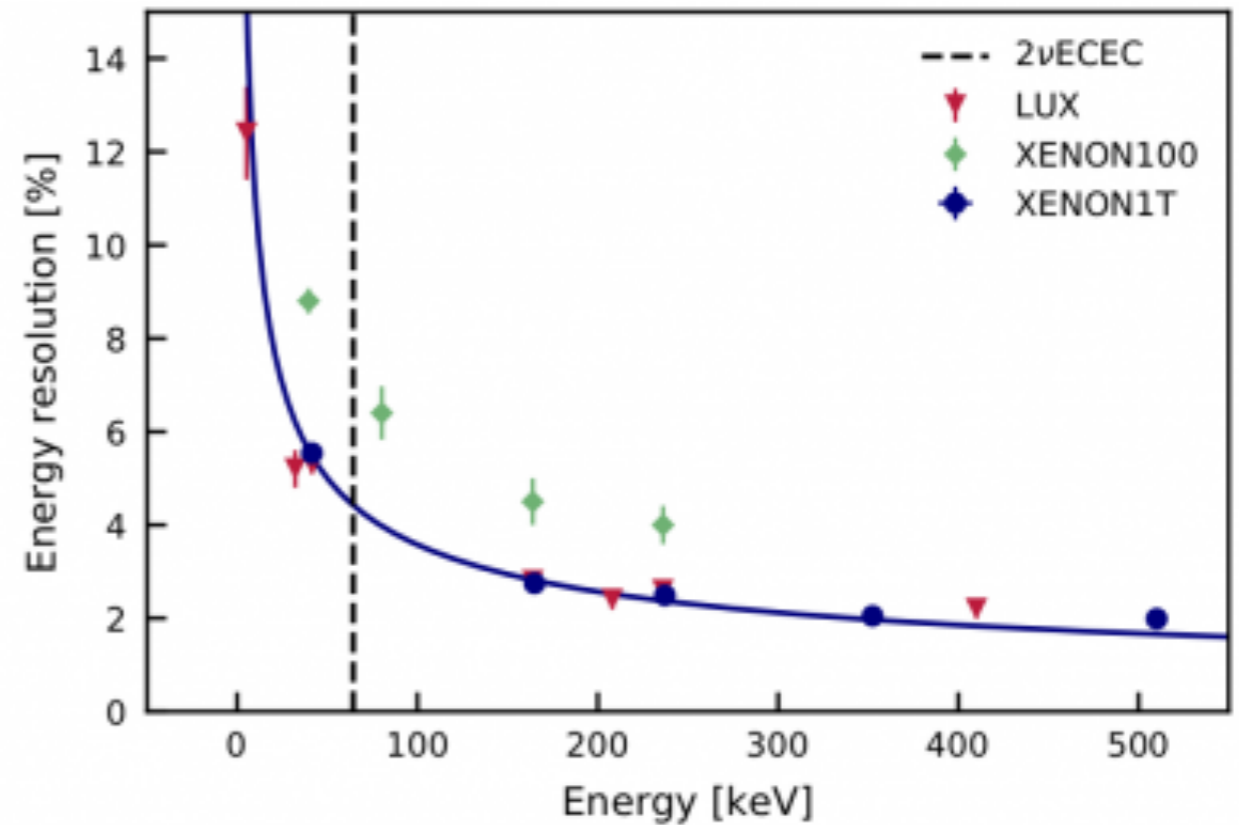


# Energy Reconstruction with LXe TPC

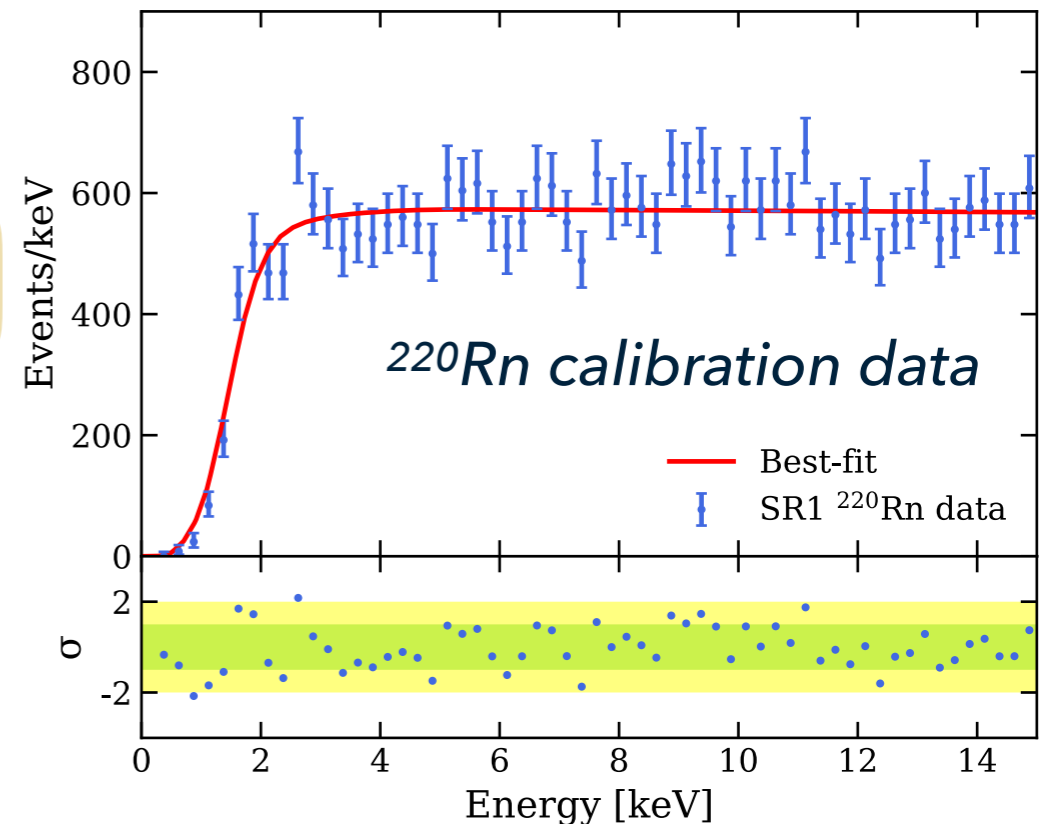
$$E = (N_{ph} + N_e) \cdot W = \left( \frac{S1}{g1} + \frac{S2}{g2} \right) \cdot W$$

where  $W = 13.7$  eV/quanta

g1 and g2: detector-specific gain constants  
 extract g1/g2 from calibration data, use it to  
 reconstruct energy of each event



$$E = W \left( \frac{S1}{g1} + \frac{S2}{g2} \right)$$



2D analysis in s1-s2 space

1D energy spectrum

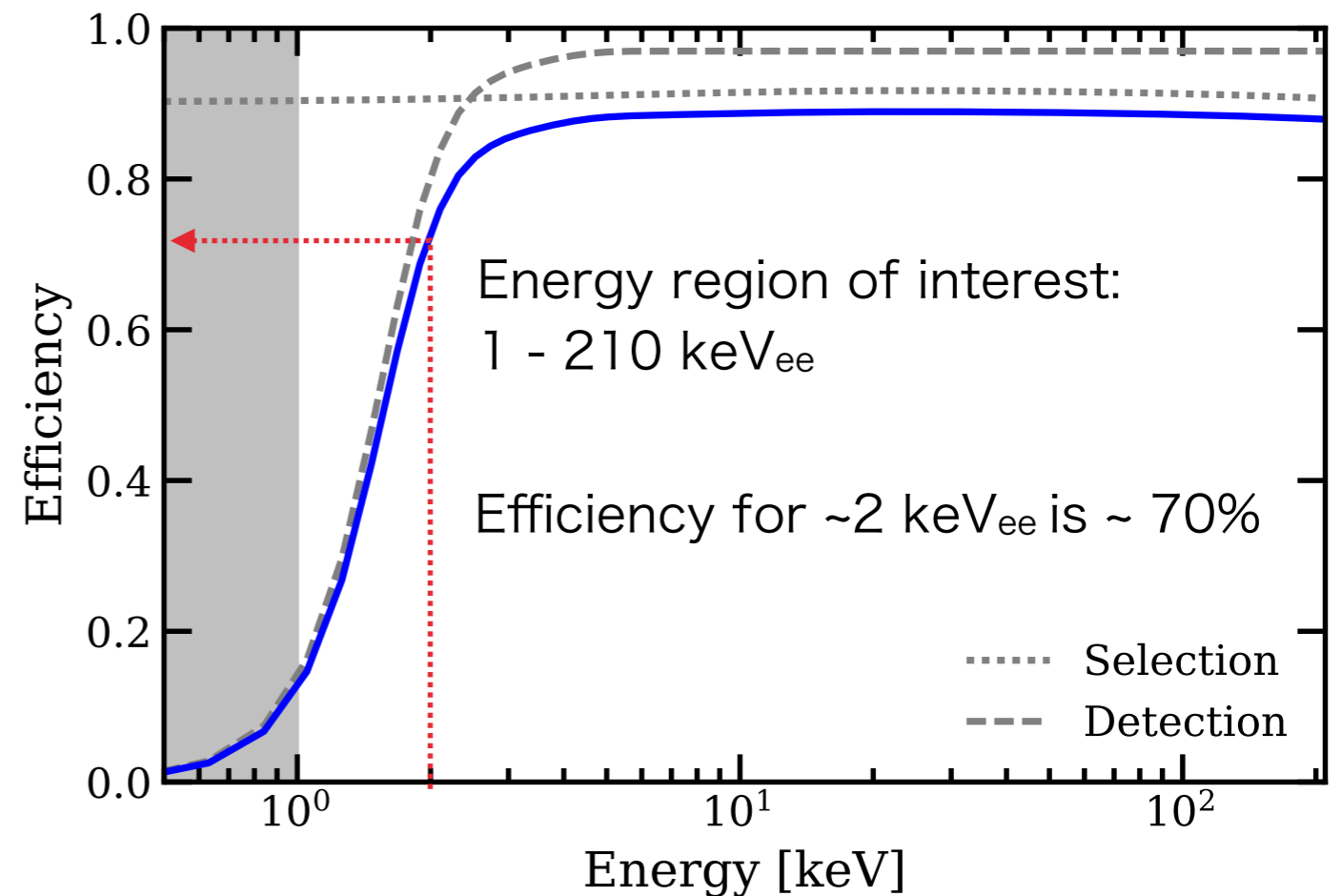
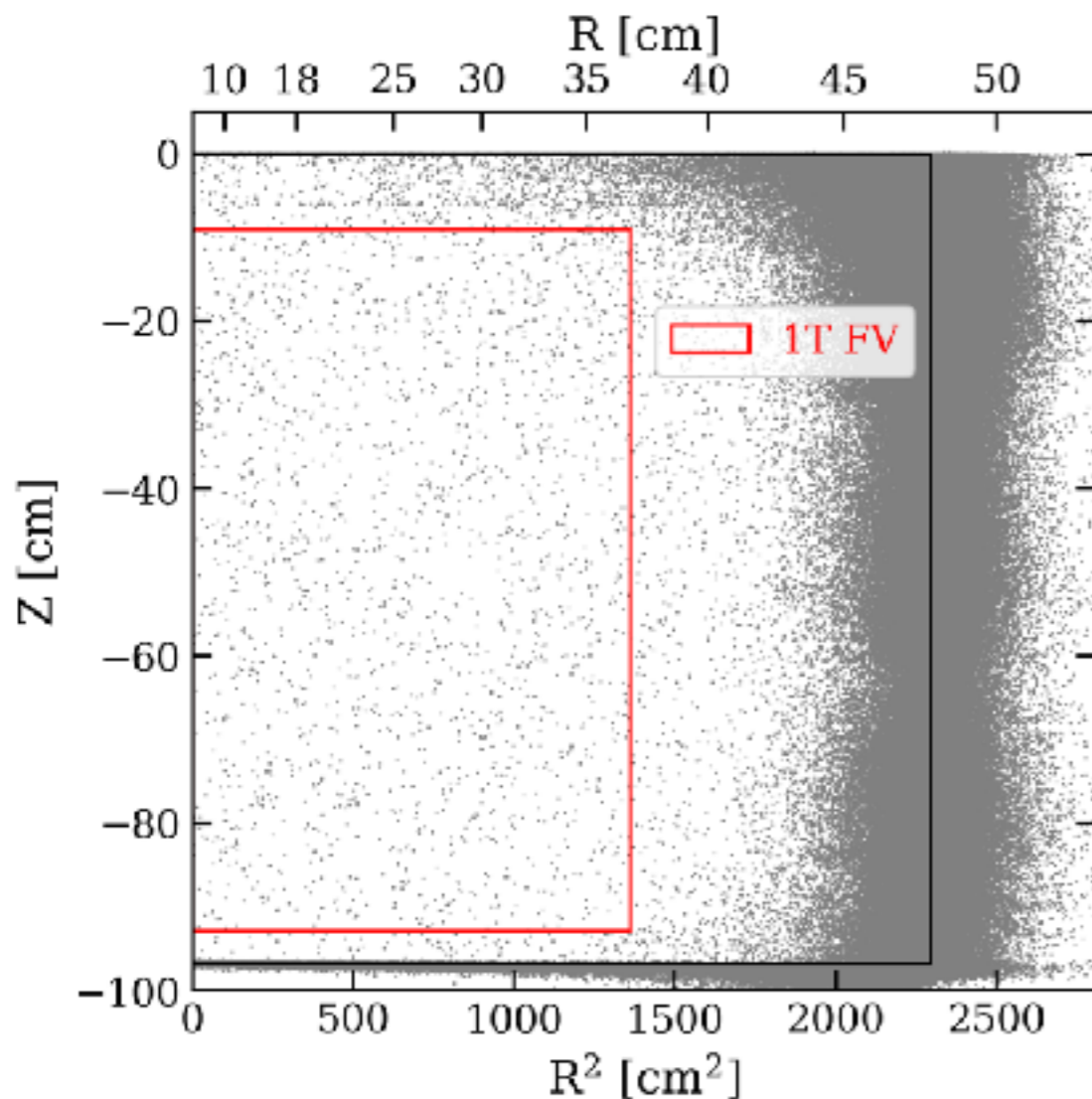
# Signal Models

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- Fiducial volume 1042 kg
- Single-scatter events, standard data quality cuts
- Higher S2 threshold ( $> 500$  pe) to remove instrumental BGs
- Detection efficiency is dominated by 3-fold PMT coincidence for S1 detection



## Three components of solar axion flux

**ABC**

$g_{ae}$

(Atomic recombination and deexcitation, Bremsstrahlung and Compton)

axion-electron interactions dominated by Bremsstrahlung and Compton

**Primakoff:**

$g_{a\gamma}$

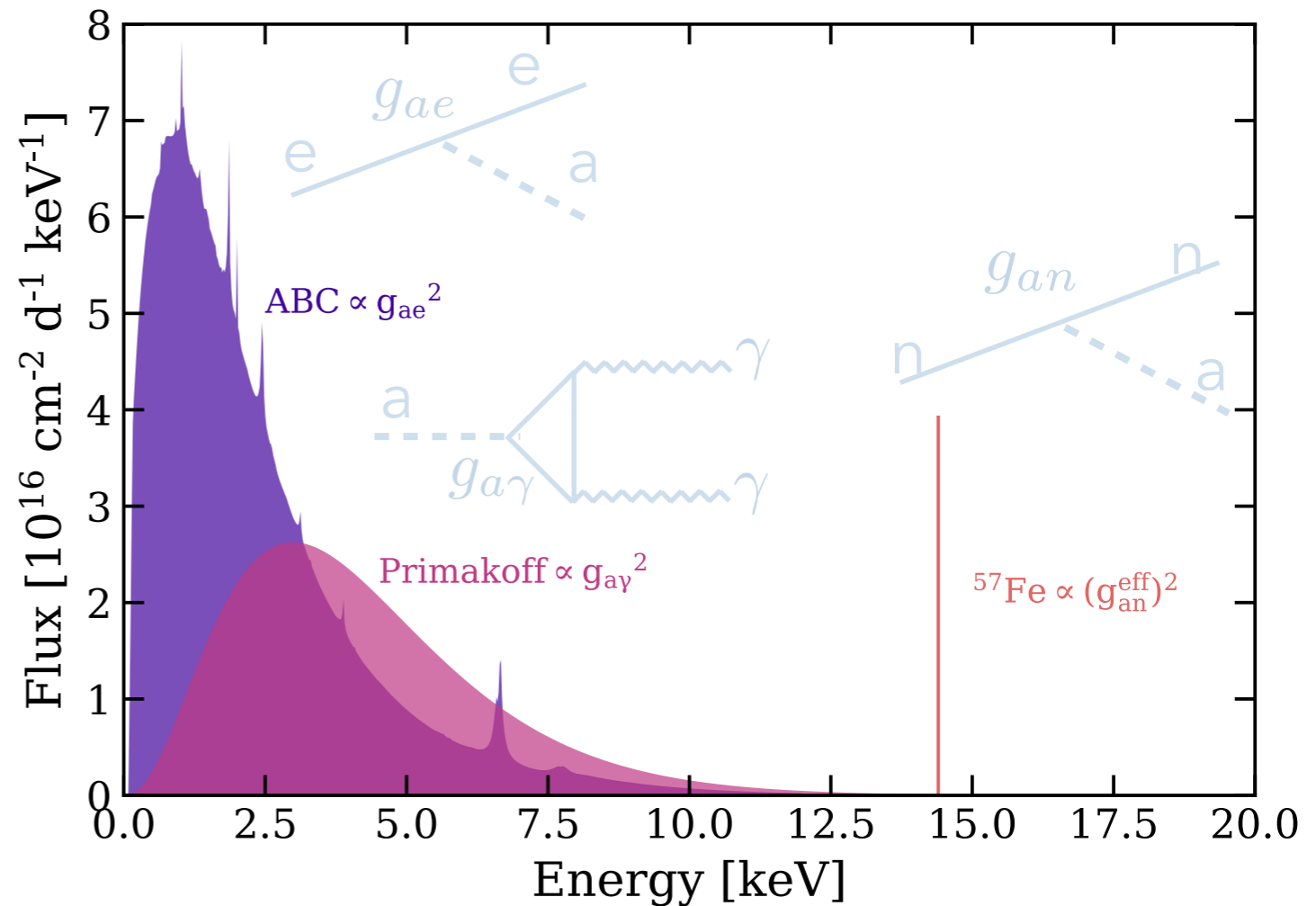
axion-photon coupling  
axions produced from photon conversion induced by the electric field of ions and electrons in the Sun.

**Fe-57 nuclear transition:**

$$g_{an}^{\text{eff}} = -1.19g_{an}^0 + g_{an}^3$$

mono energetic **14.4 keV** M1 transition  
effective axion-nucleon coupling

**Emerge with keV-scale energies**  
In principle, axions from all 3 couplings can be present at the same time.

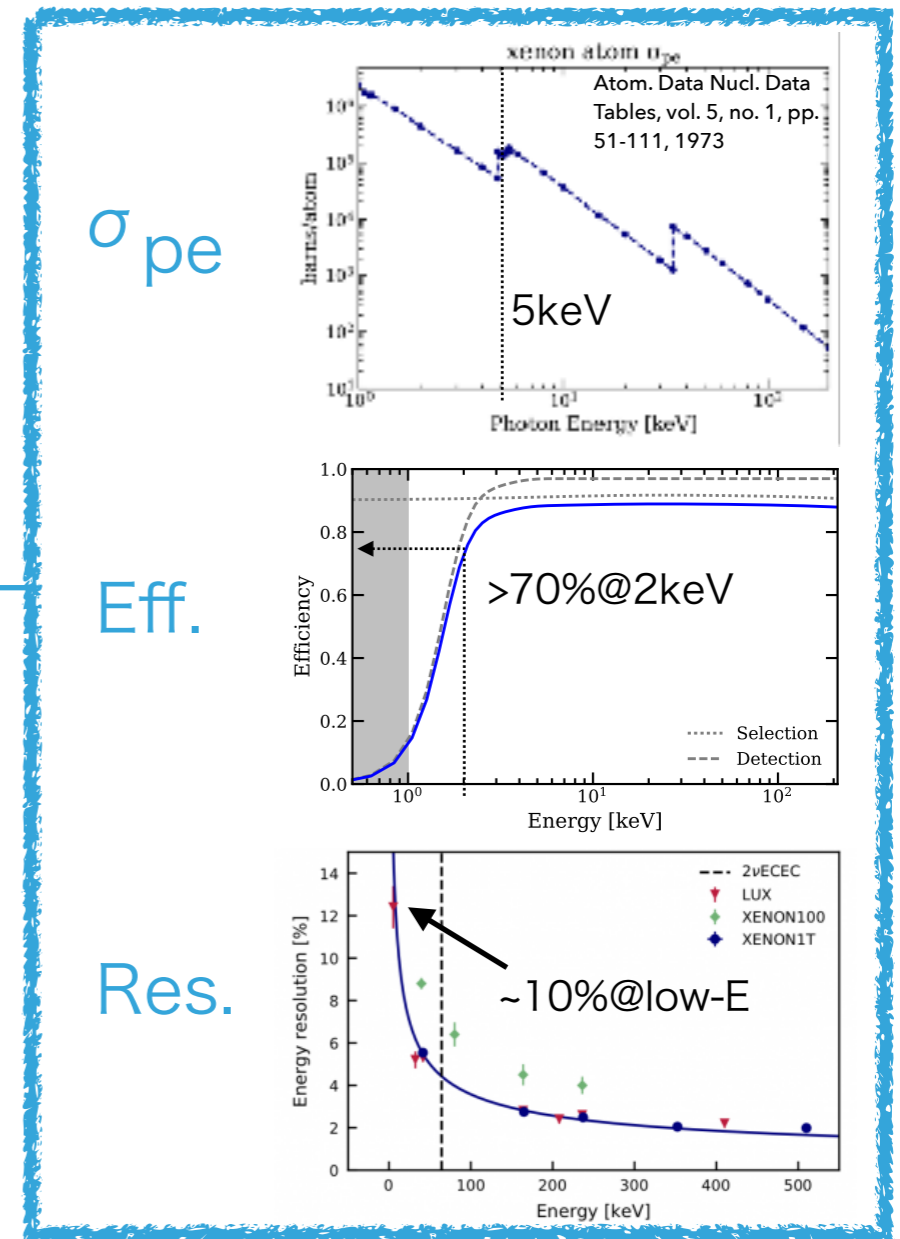
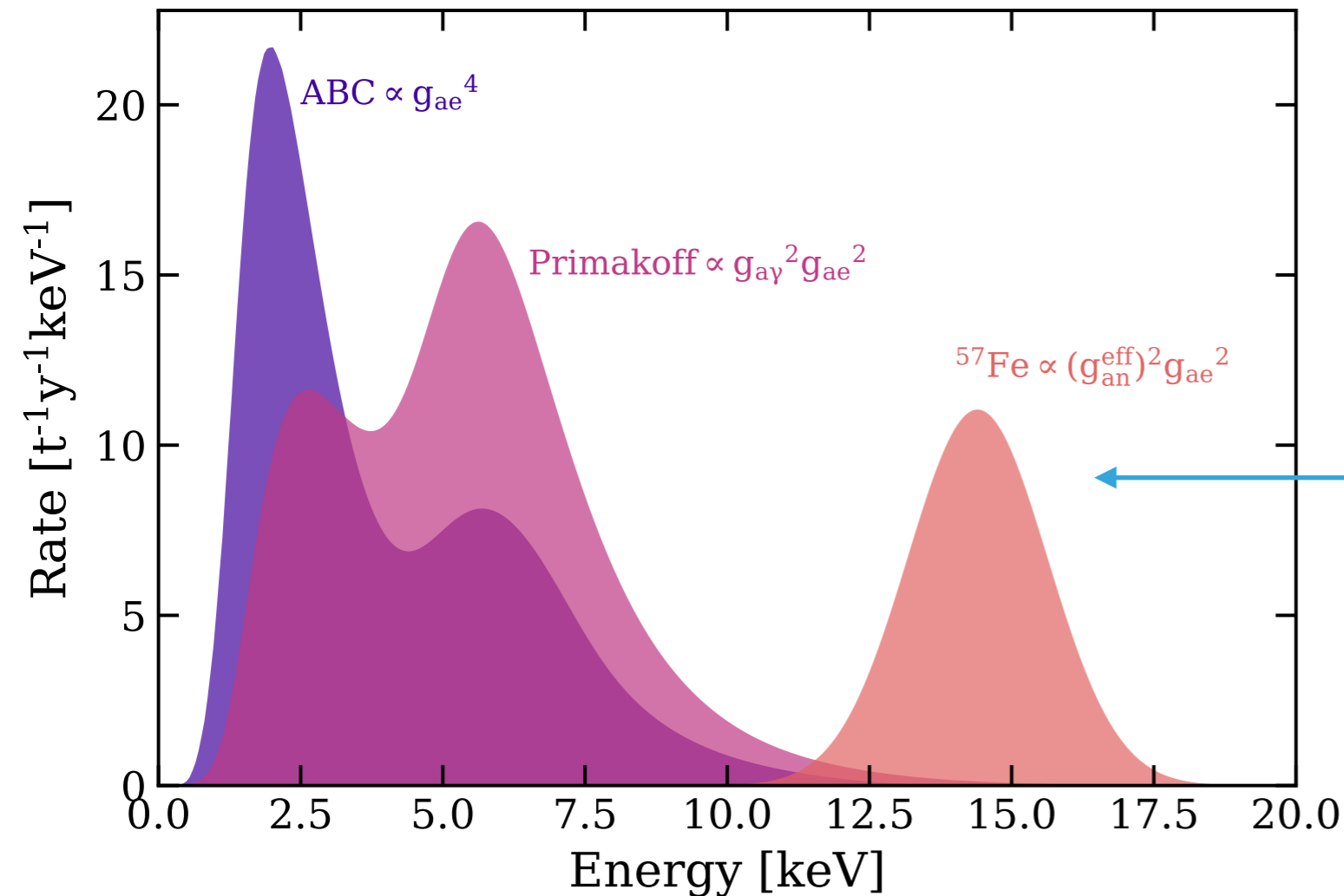


XENON1T sensitive to all 3 channels via **coupling to electrons**  $g_{ae}$   
(electronic recoils via axio-electric effect).

$$\sigma_{ae} = \sigma_{pe} \frac{g_{ae}^2}{\beta} \frac{3E_a^2}{16\pi\alpha m_e^2} \left(1 - \frac{\beta^{2/3}}{3}\right)$$

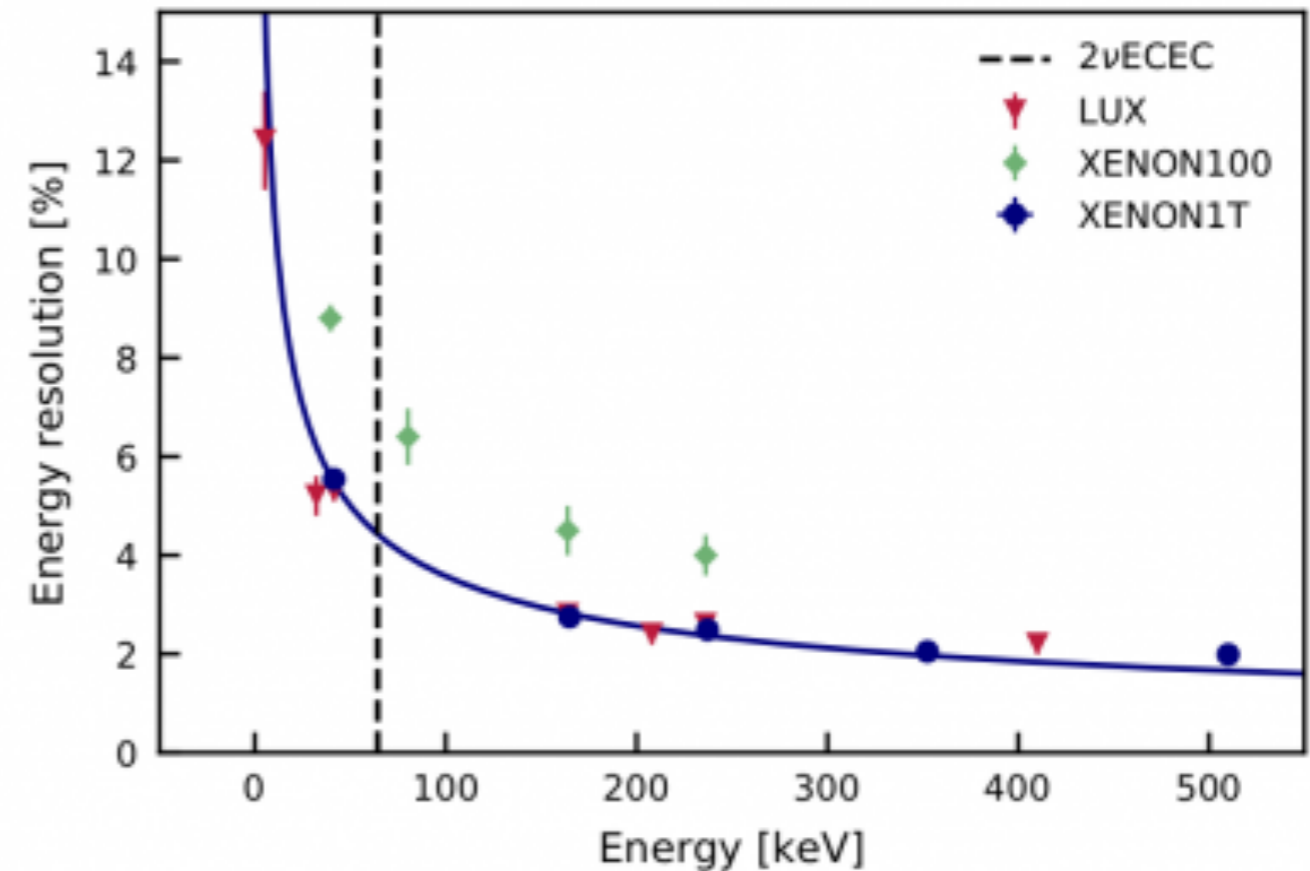
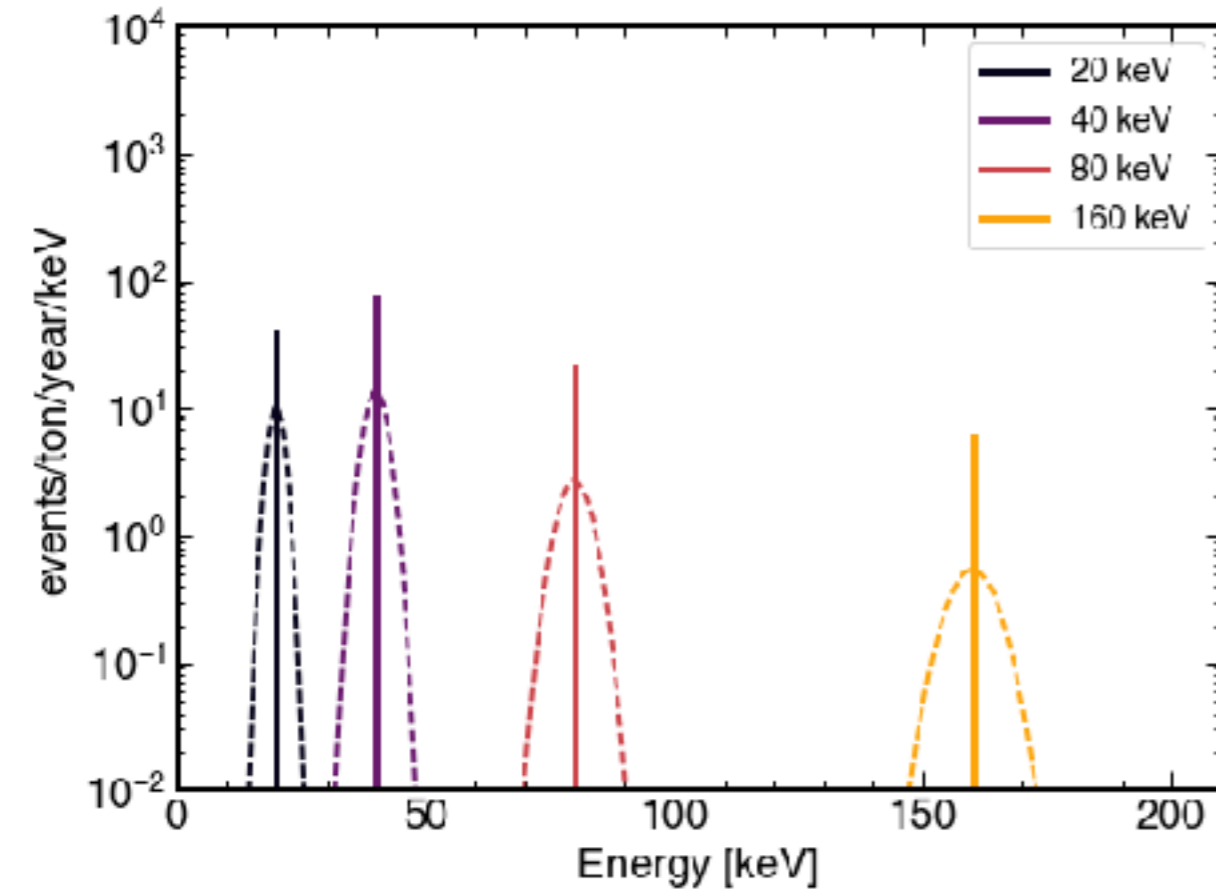
$\beta$  (Ea): velocity (energy) of axion

Expected rate in xenon convolved with detector effects (resolution, efficiency) and  $\sigma_{pe}$



- For Primakoff and  $^{57}\text{Fe}$ , can only deduce product of 2 couplings.
- All the three flux components are considered completely independent of each other  
 → Model-independent search: parameters of interest =  $g_{ae}$  VS.  $g_{ae}g_{a\gamma}$  VS.  $g_{ae}g_{an}^{\text{eff}}$
- XENON1T is more sensitive to DFSZ model, where axions couple to electrons at tree level compared to KSVZ model (couples to electrons at loop level)

Assuming ALPs/dark photons are non-relativistic and make up all of the local dark matter, the expected signal is a mono-energetic peak at the rest mass of the particle



For ALPs

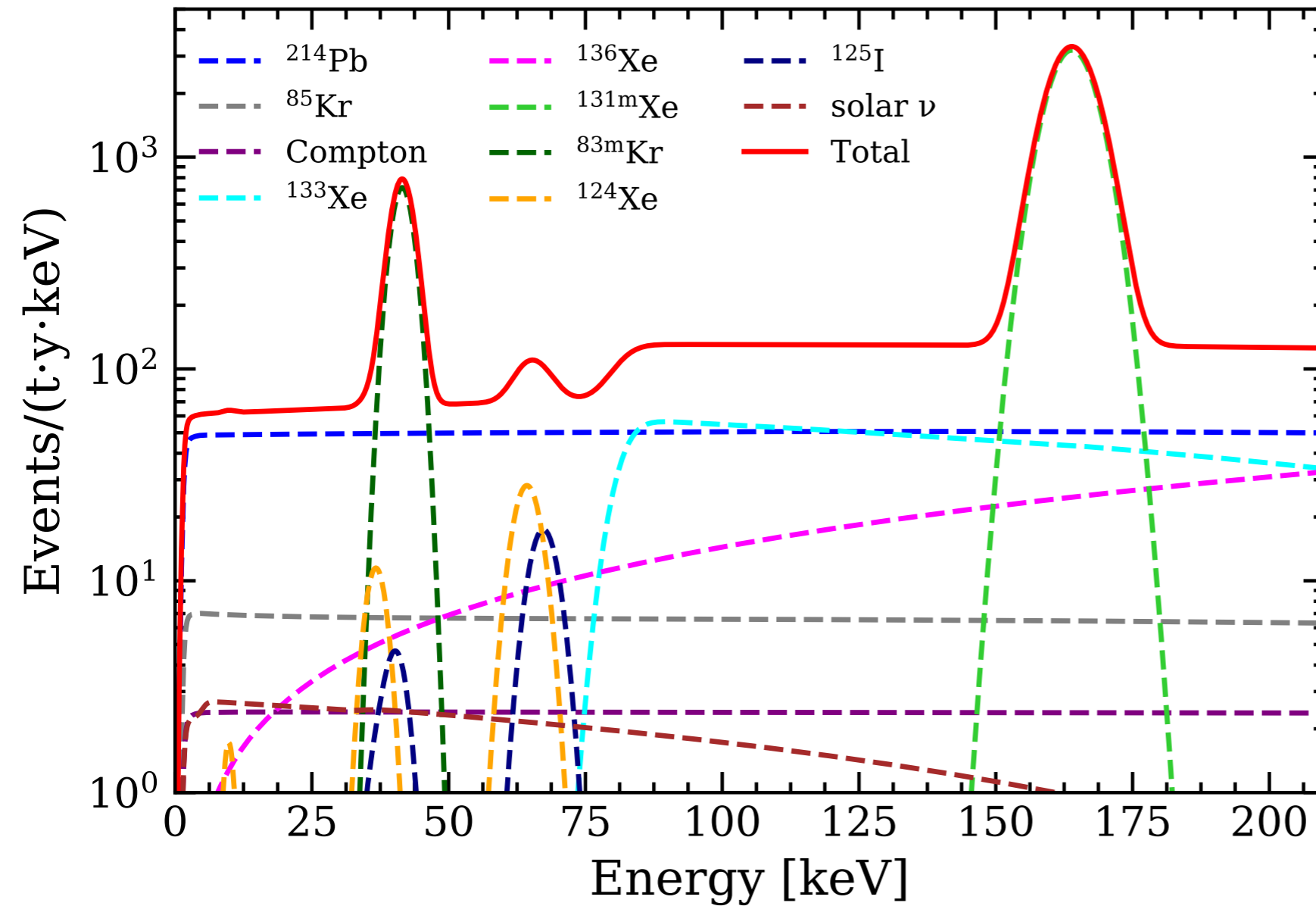
$$R \simeq \frac{1.5 \times 10^{19}}{A} g_{ae}^2 \left( \frac{m_a}{\text{keV}/c^2} \right) \left( \frac{\sigma_{pe}}{b} \right) \text{kg}^{-1} \text{d}^{-1}$$

For dark photons

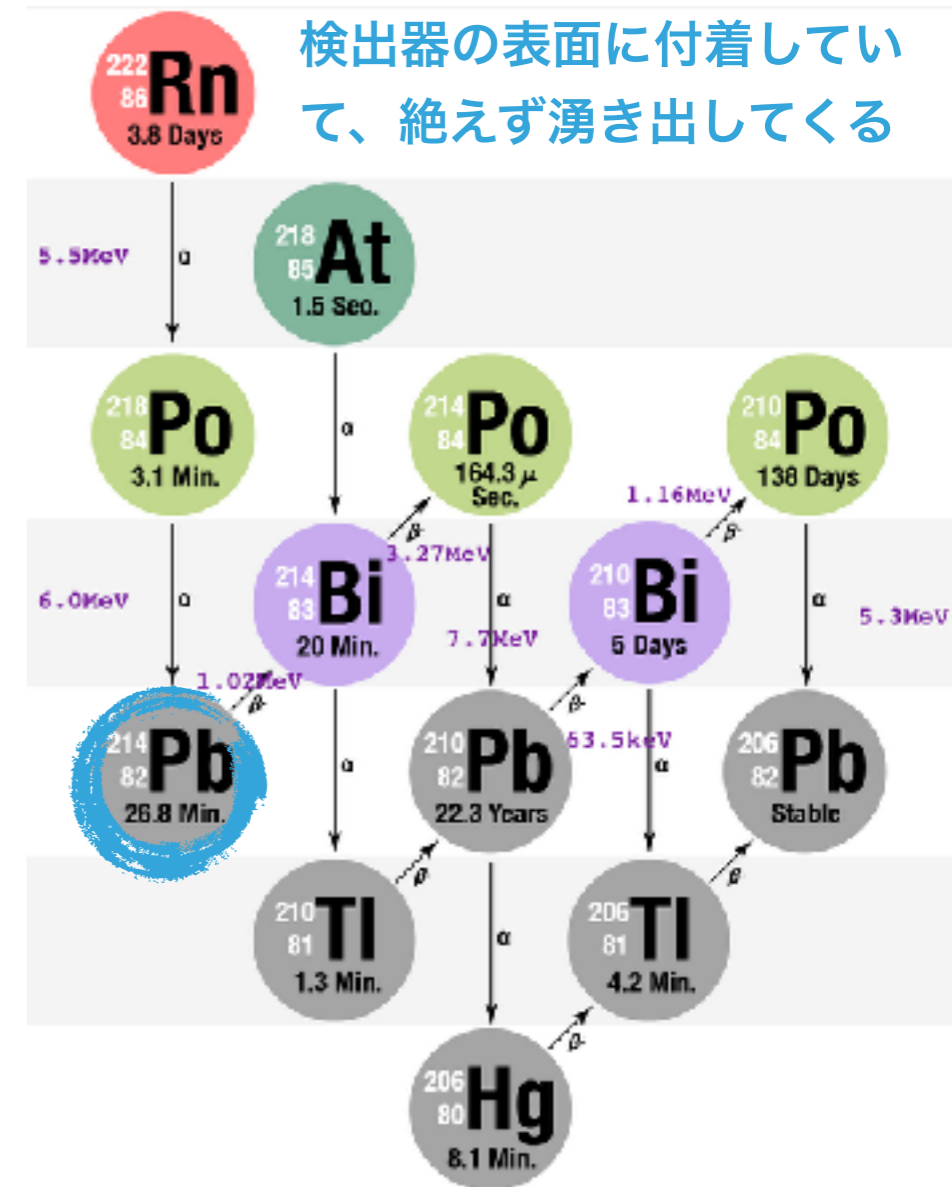
$$R \simeq \frac{4.7 \times 10^{23}}{A} \kappa^2 \left( \frac{\text{keV}/c^2}{m_V} \right) \left( \frac{\sigma_{pe}}{b} \right) \text{kg}^{-1} \text{d}^{-1}$$

# Background Models

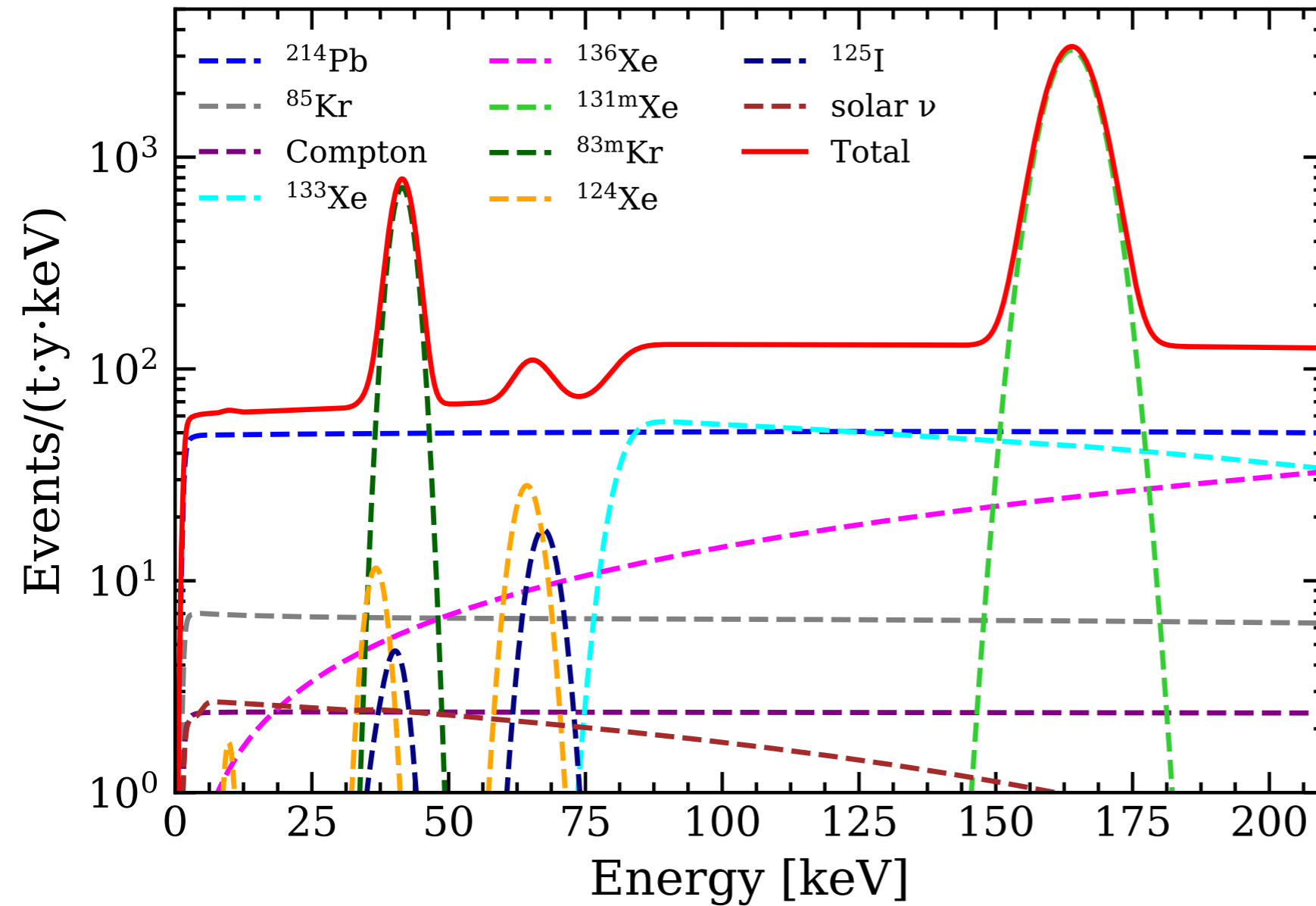
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## $^{214}\text{Pb}$ : main ER BG



Predicted energy spectra based on detailed modeling of each background component  
 Rates constrained by measurements and/or time dependence



**10 Components**

*Internal Backgrounds*

$^{214}\text{Pb}$        $^{85}\text{Kr}$

*Intrinsic Backgrounds*

$^{124}\text{Xe}$        $^{136}\text{Xe}$

solar neutrino

*Activated Backgrounds*

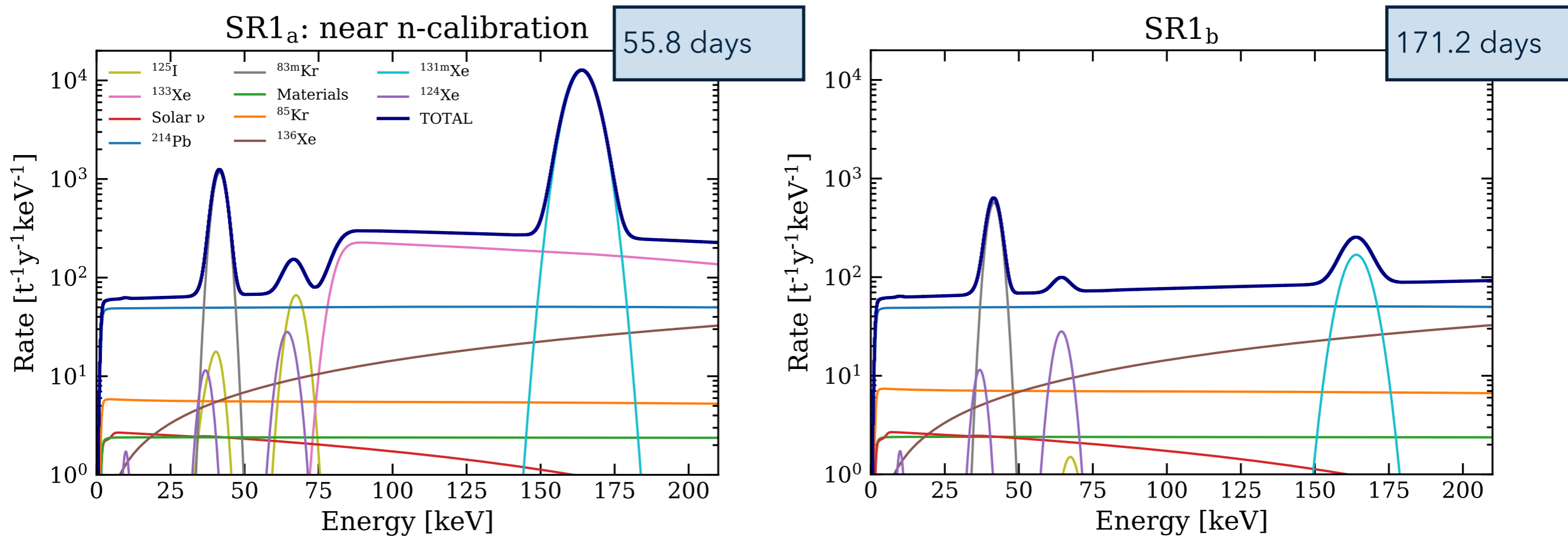
$^{131\text{m}}\text{Xe}$        $^{133}\text{Xe}$

$^{125}\text{I}$

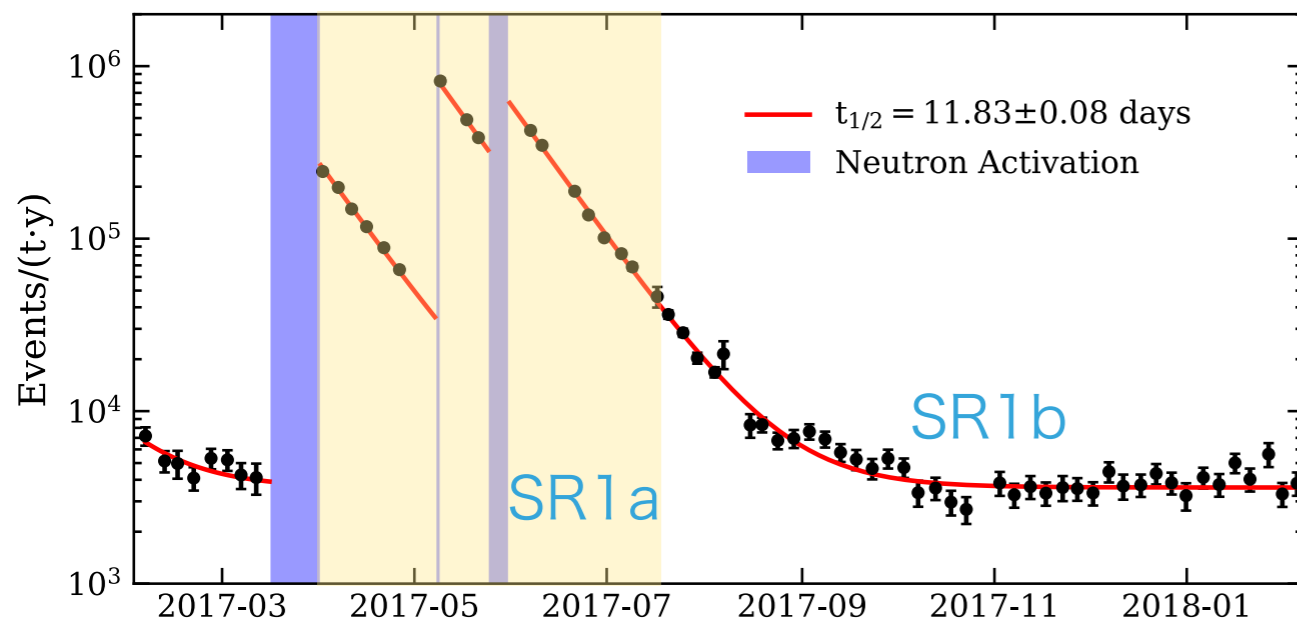
*Contaminant Backgrounds*

$^{83\text{m}}\text{Kr}$       Materials

Predicted energy spectra based on detailed modeling of each background component  
 Rates constrained by measurements and/or time dependence



Time-evolution and model of  $^{131m}\text{Xe}$  (generated by neutron activation)



**Divided into two datasets, fit simultaneously.**

- SR1<sub>a</sub>: <50 days from neutron calibration, includes more activated backgrounds
- SR1<sub>b</sub>: the rest, less activated backgrounds
- background model denoted B<sub>0</sub>



# Unbinned Profile Likelihood Analysis

$$\begin{aligned} \mathcal{L}(\mu_s, \mu_b, \theta) &= \text{Poiss}(N | \mu_{\text{tot}}) \\ &\times \prod_i^N \left( \sum_j \frac{\mu_{b_j}}{\mu_{\text{tot}}} f_{b_j}(E_i, \theta) + \frac{\mu_s}{\mu_{\text{tot}}} f_s(E_i, \theta) \right) \\ &\times \prod_m C_{\mu_{\tau_m}}(\mu_{b_m}) \times \prod_n C_{\theta_n}(\theta_n), \quad (14) \\ \mu_{\text{tot}} &\equiv \sum_j \mu_{b_j} + \mu_s, \end{aligned}$$

Profile over the nuisance parameters

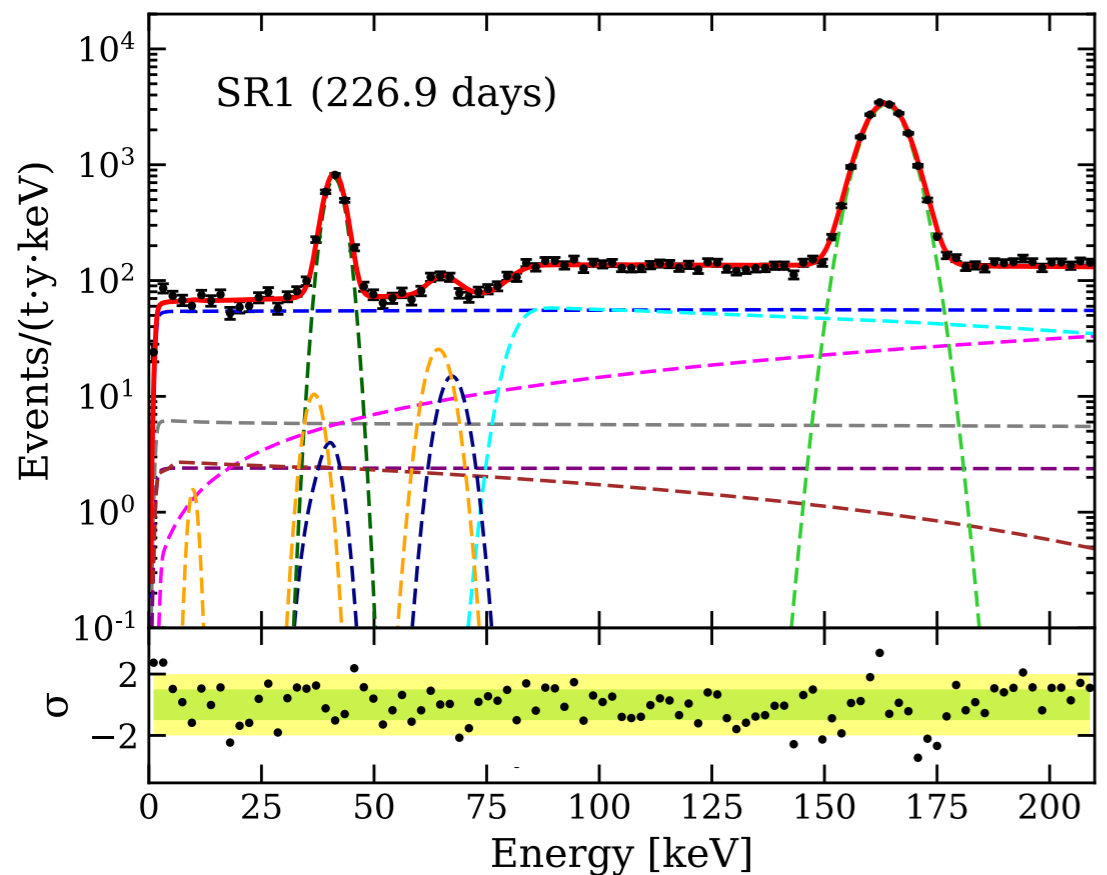
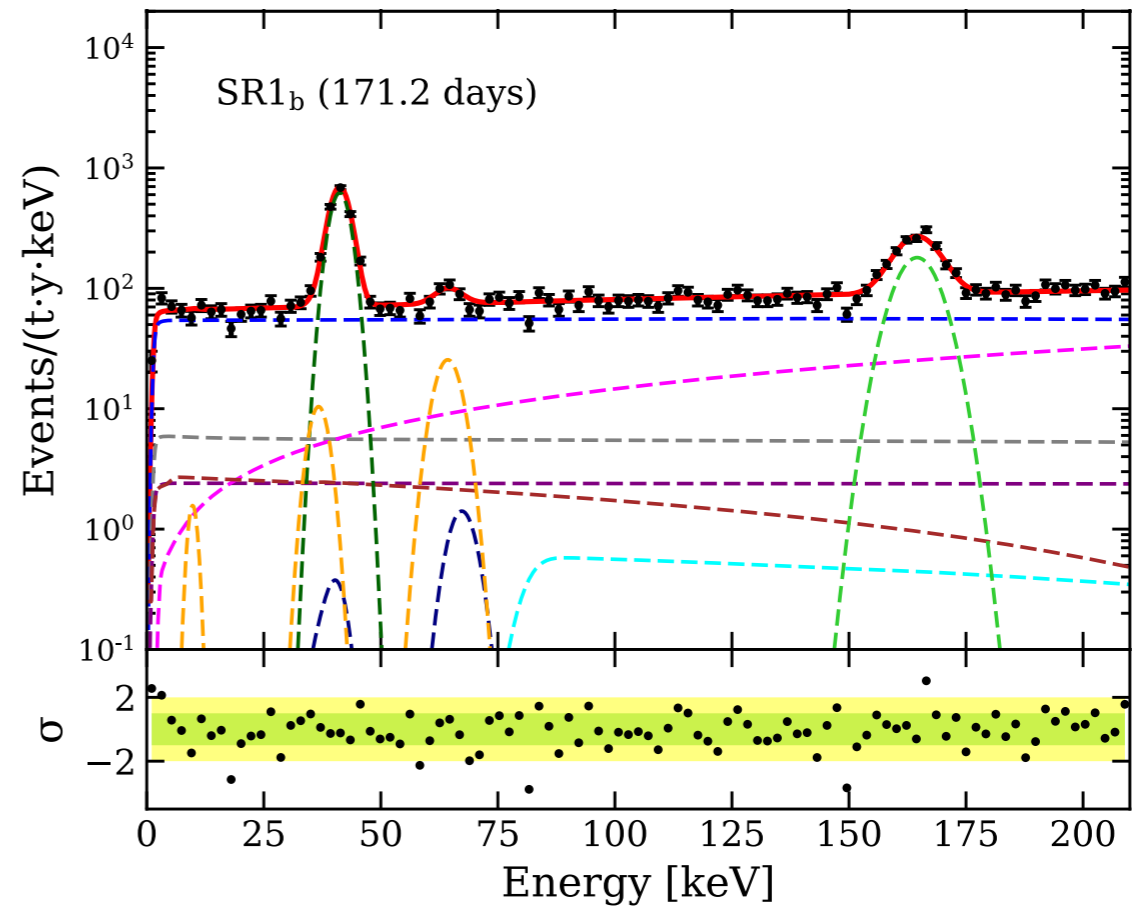
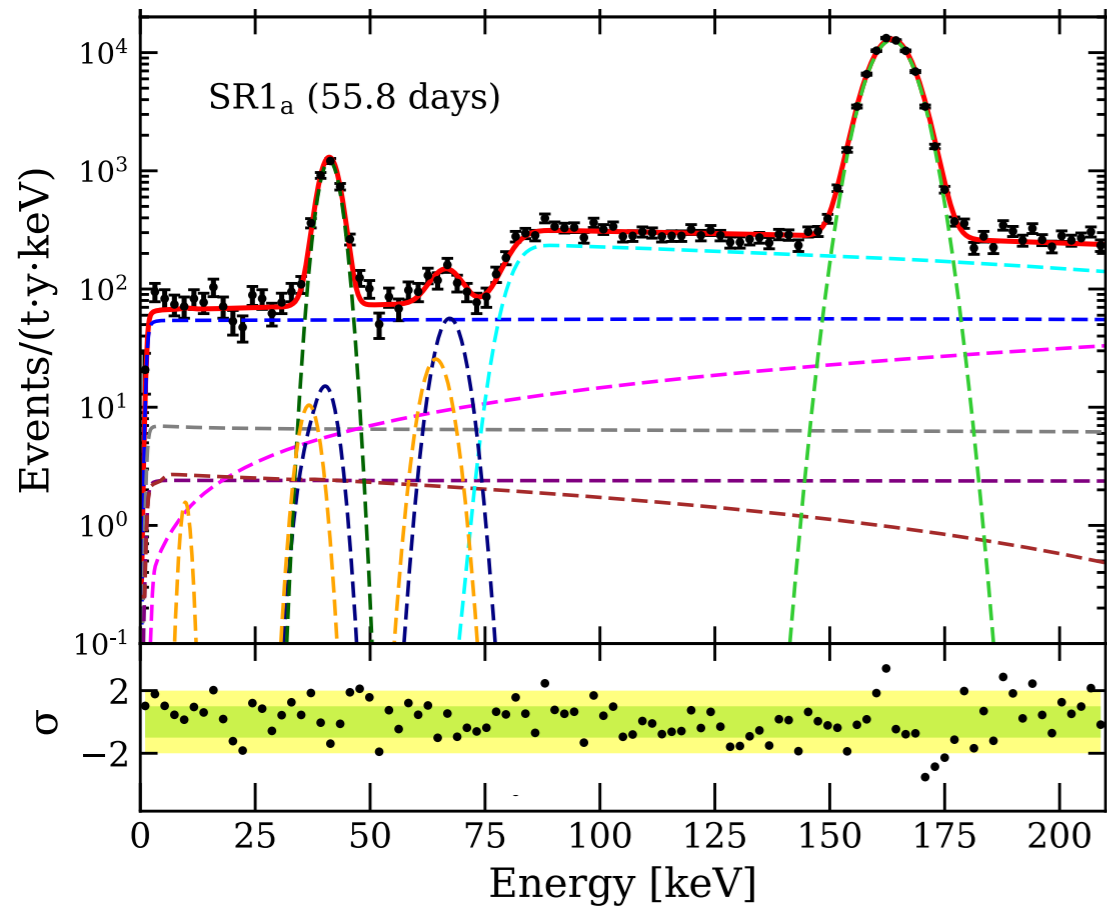
Combining the likelihoods of the 2 partitions

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# Fit to Data

$$\mathcal{L} = \mathcal{L}_a \times \mathcal{L}_b$$

SR1a SR1b

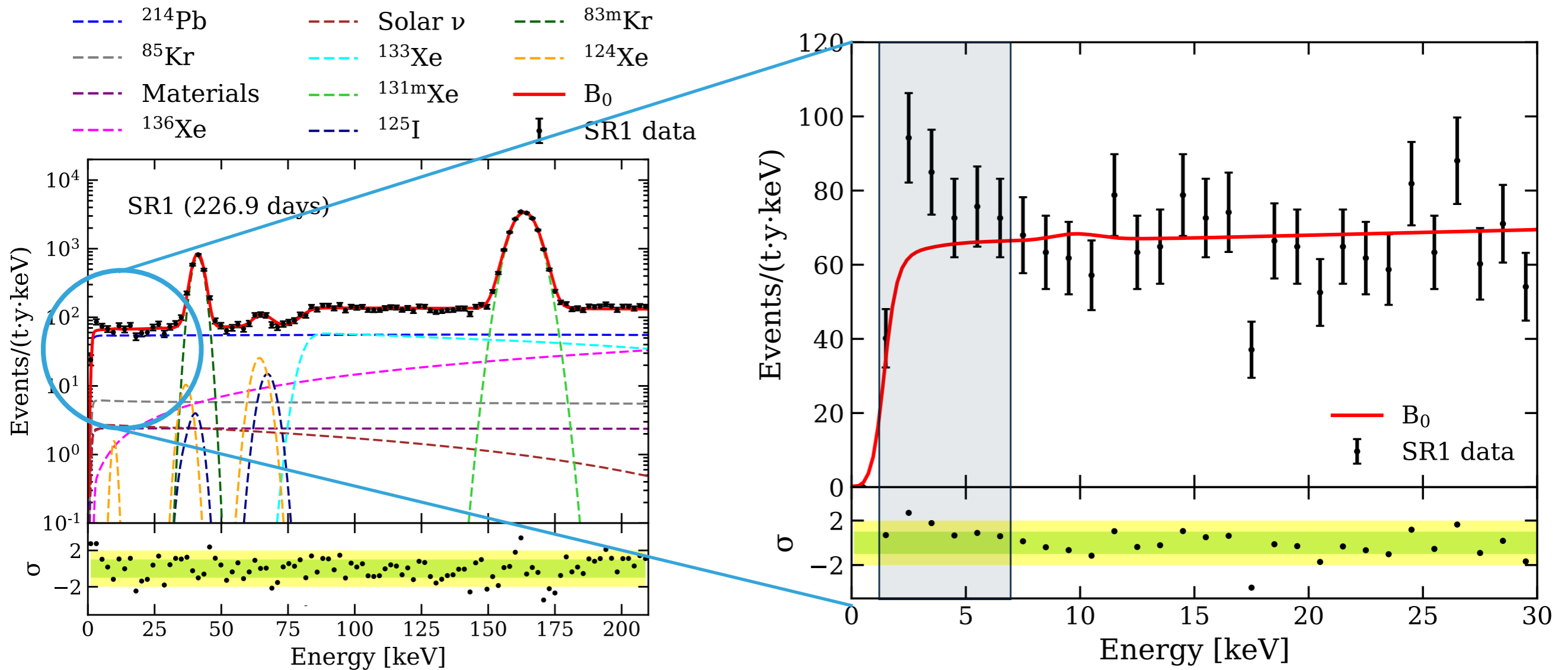


- $^{214}\text{Pb}$
- $^{85}\text{Kr}$
- Materials
- $^{136}\text{Xe}$
- Solar  $\nu$
- $^{133}\text{Xe}$
- $^{131\text{m}}\text{Xe}$
- $^{125}\text{I}$
- $^{83\text{m}}\text{Kr}$
- $^{124}\text{Xe}$
- $\text{B}_0$
- █ SR1 data

Decent matching across the whole energy range in 1-210 keV

(76 +/- 2) events/(t.y.keV) in [1, 30] keV

Lowest background rate ever achieved in this energy range!



## Excess between 1-7 keV

**285** events observed

vs.

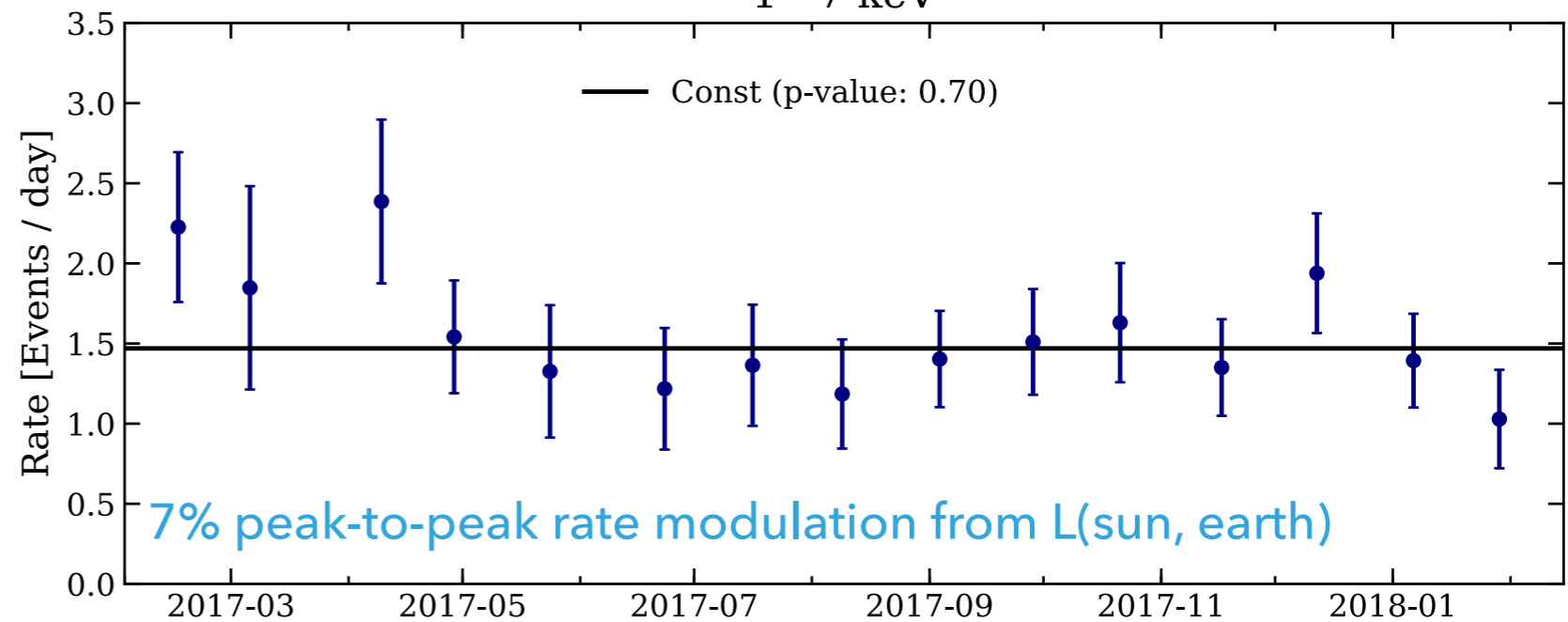
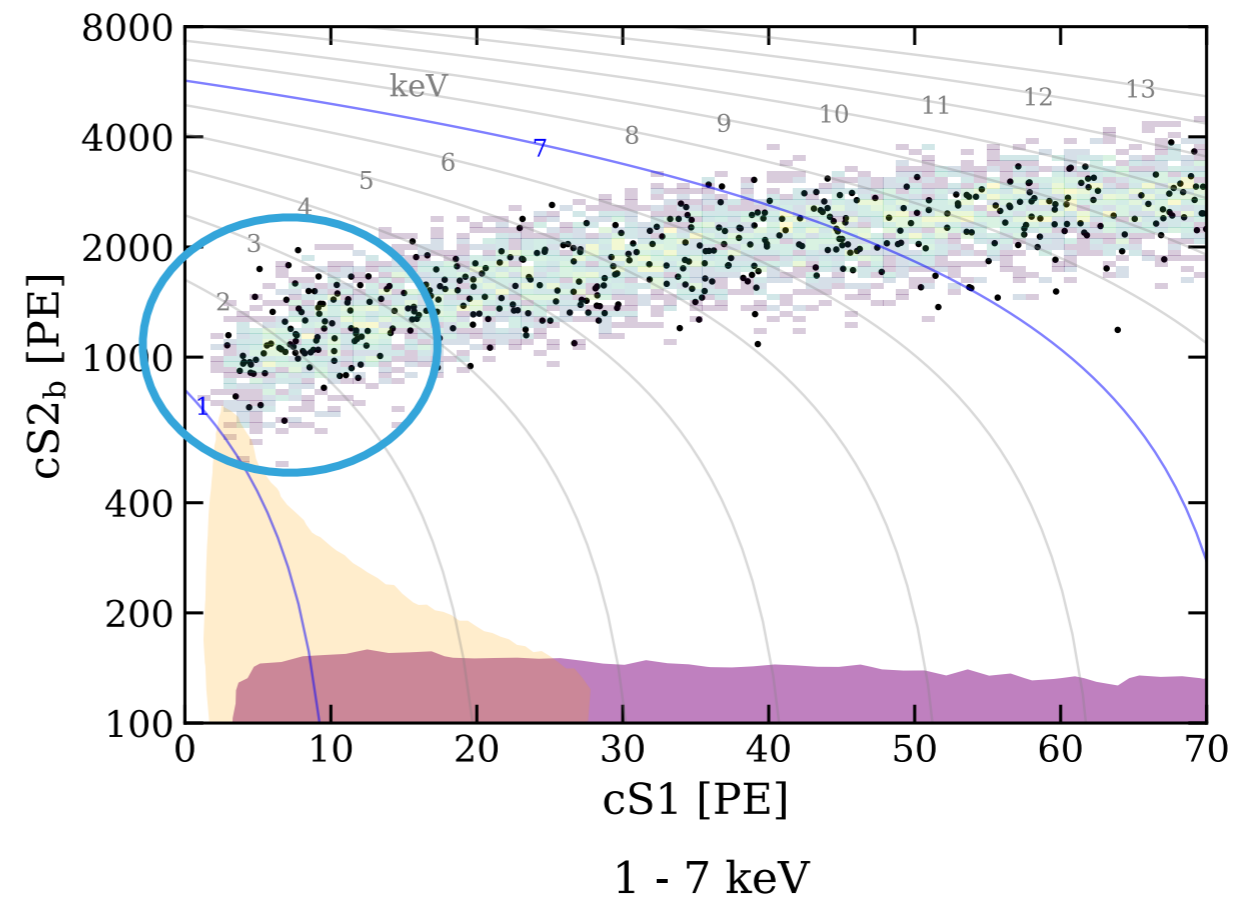
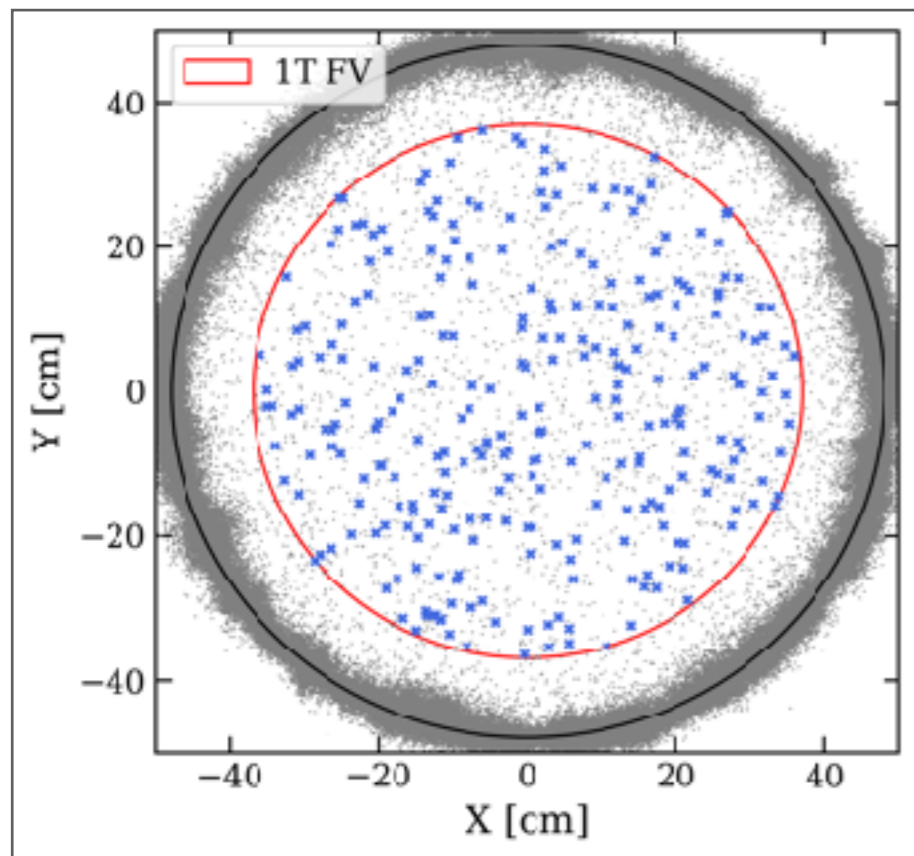
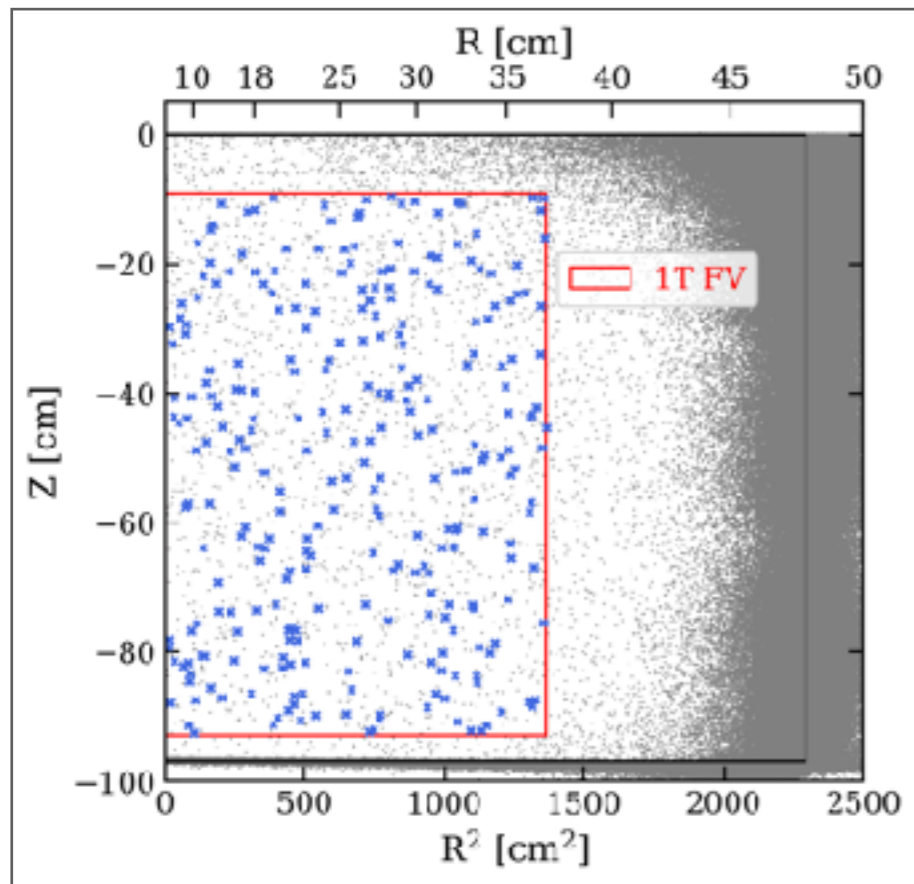
**232** events expected (from BG-only best-fit)

Would be a  **$3.3\sigma$**  Poissonian fluctuation

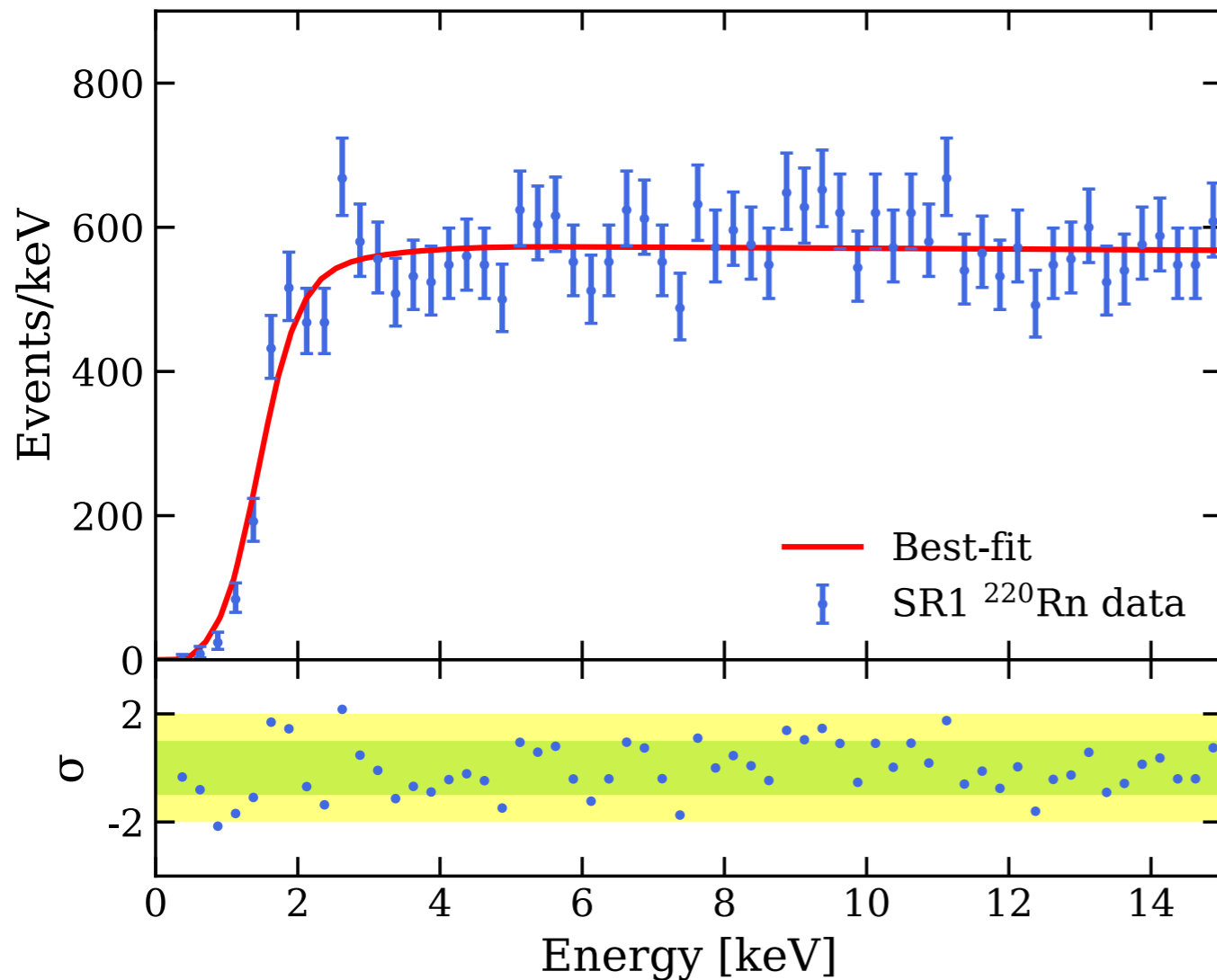
**Are We Missing Something?**

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# Event Location / Time-dependence



- Uniform in 3D-space and cS1/cS2b plane
- consistent with constant time (but with low statistics)

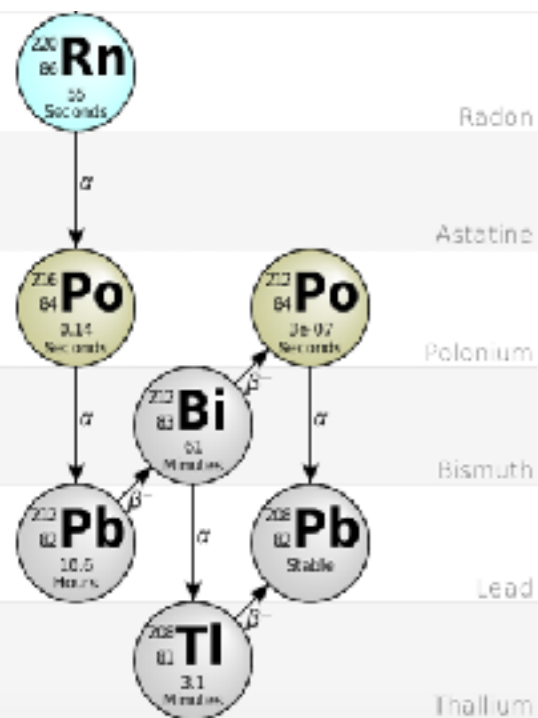


<sup>220</sup>Rn calibration reconstructs as expected

Fit to <sup>220</sup>Rn (<sup>212</sup>Pb) calibration data using same analysis framework

**Validates efficiency and energy reconstruction**

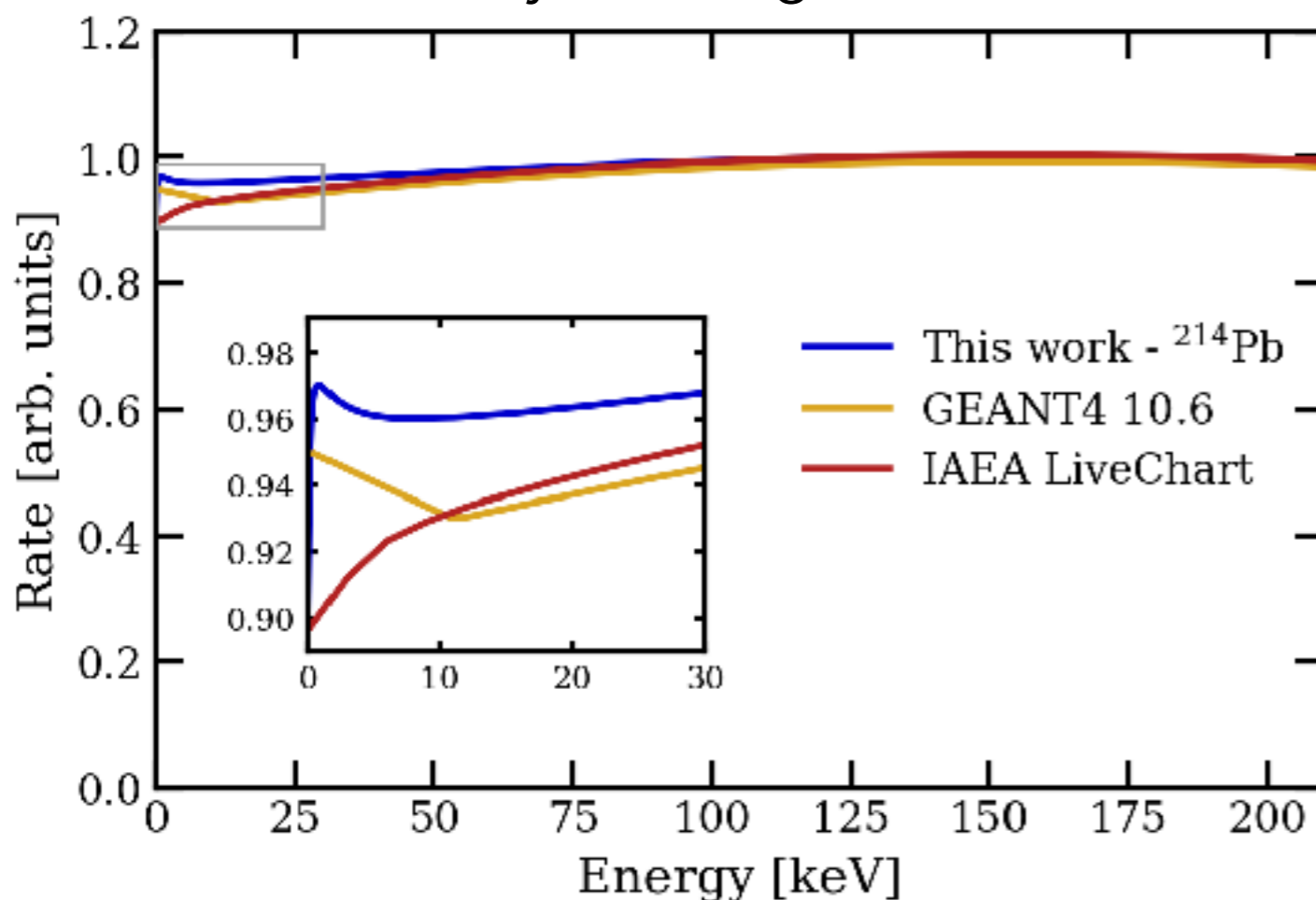
g.o.f p=0.58



$\beta$ -decay of Pb212 is used to calibrate detector's response to ER background

[Again, unbinned fit is performed here.](#)

Calculated by X. Mougeot



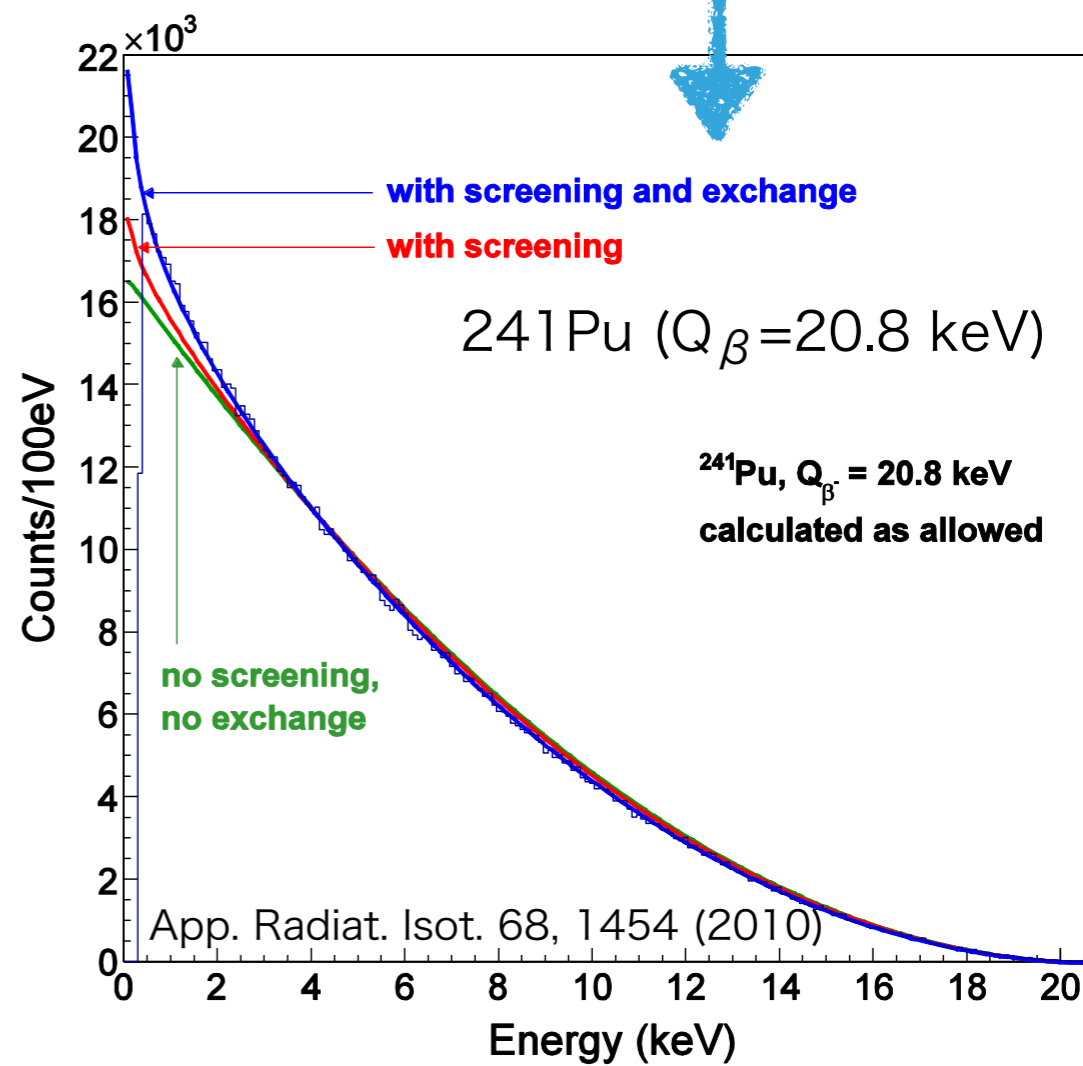
- IAEA model: No screening/exchange effects
- GEANT4 model: only screening effect
- This work: both screening/exchange effects

Atomic screening and exchange effects can increase rate at low energies.

~6% uncertainty on the shape

~50% needed to account for excess

Good agreement between measurements and calculation for <sup>241</sup>Pu ( $Q_{\beta}=20.8$  keV) and <sup>63</sup>Ni ( $Q_{\beta}=67$  keV)



Low-energy (Q value 18.6 keV)

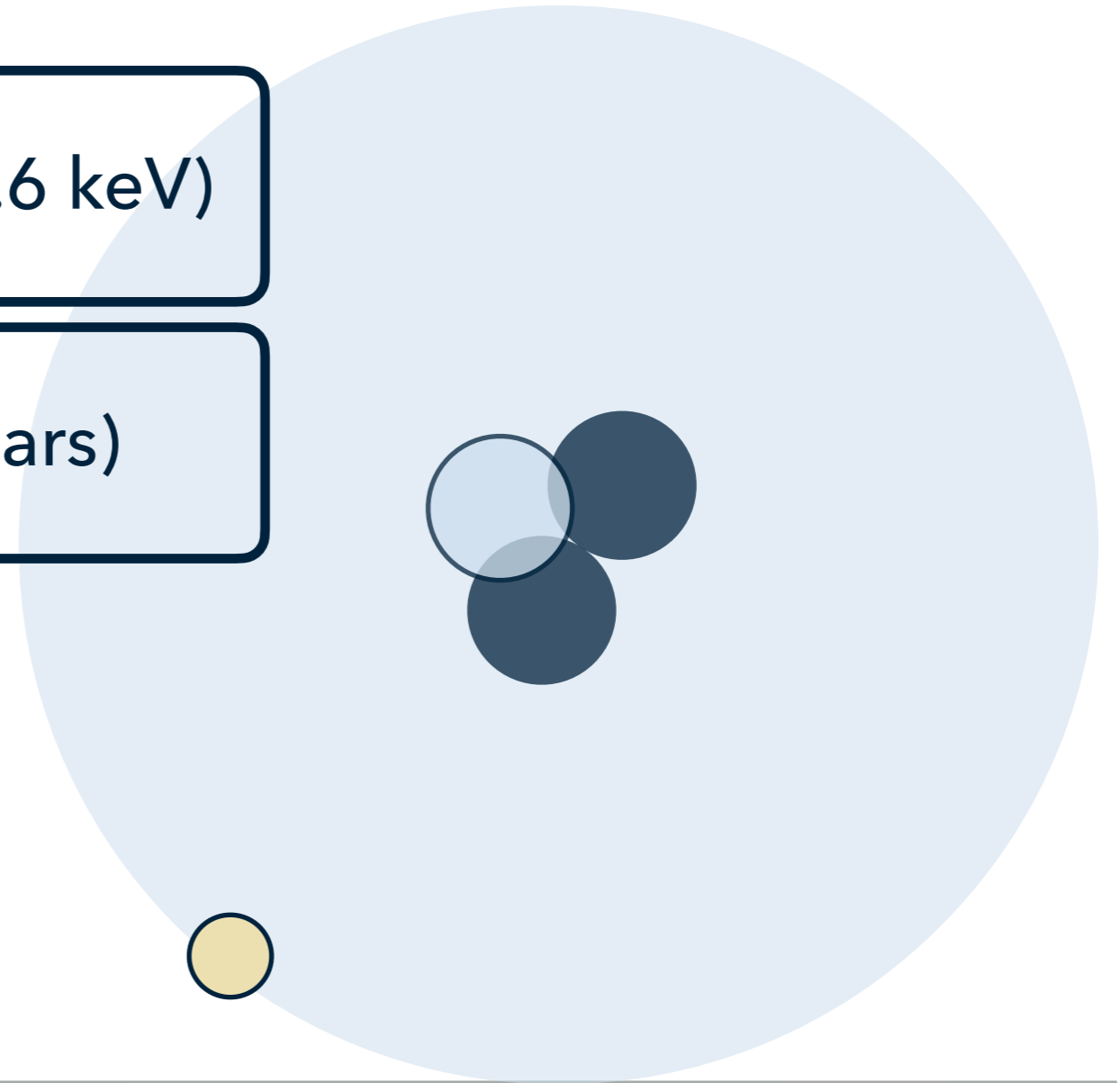
Long half life (12.3 years)

# Tritium?

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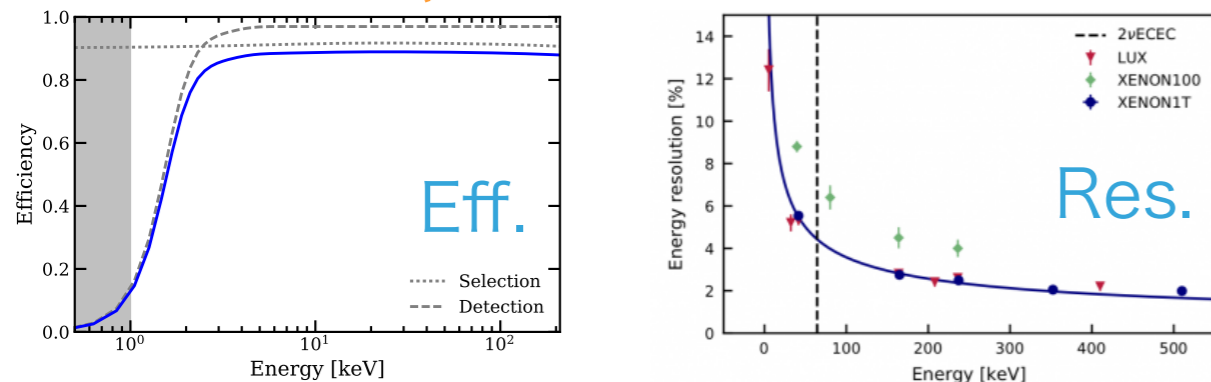
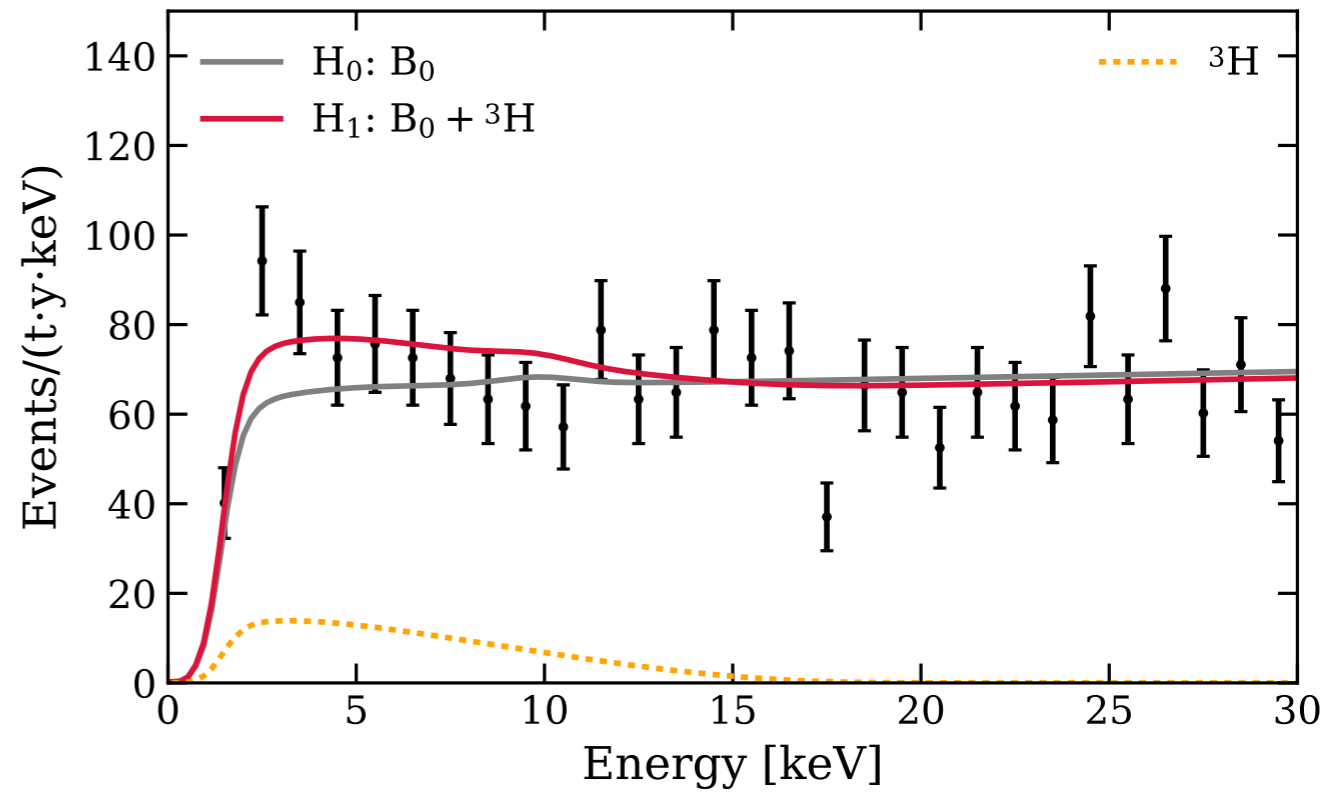
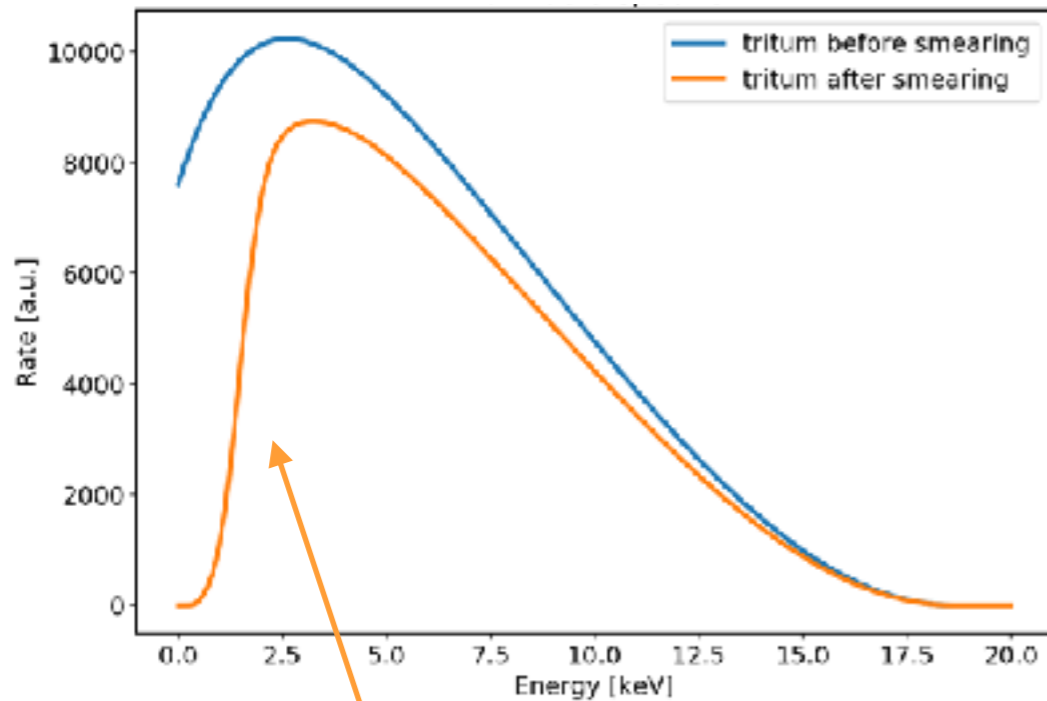
1. Cosmogenic production?

2. Atmospherically abundant?





# Testing Tritium Hypothesis



**Tritium favored over background-only at 3.2σ**

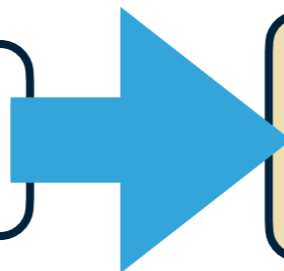
Best-fit tritium rate:

$$159 \pm 51 \text{ events}/(\text{t} \cdot \text{y} \cdot \text{keV})$$

<sup>3</sup>H half-life 12.3 years (too long to observe in SR1)

<sup>3</sup>H:Xe concentration:

$$6.2 \pm 2.0 \times 10^{-25} \text{ mol/mol}$$



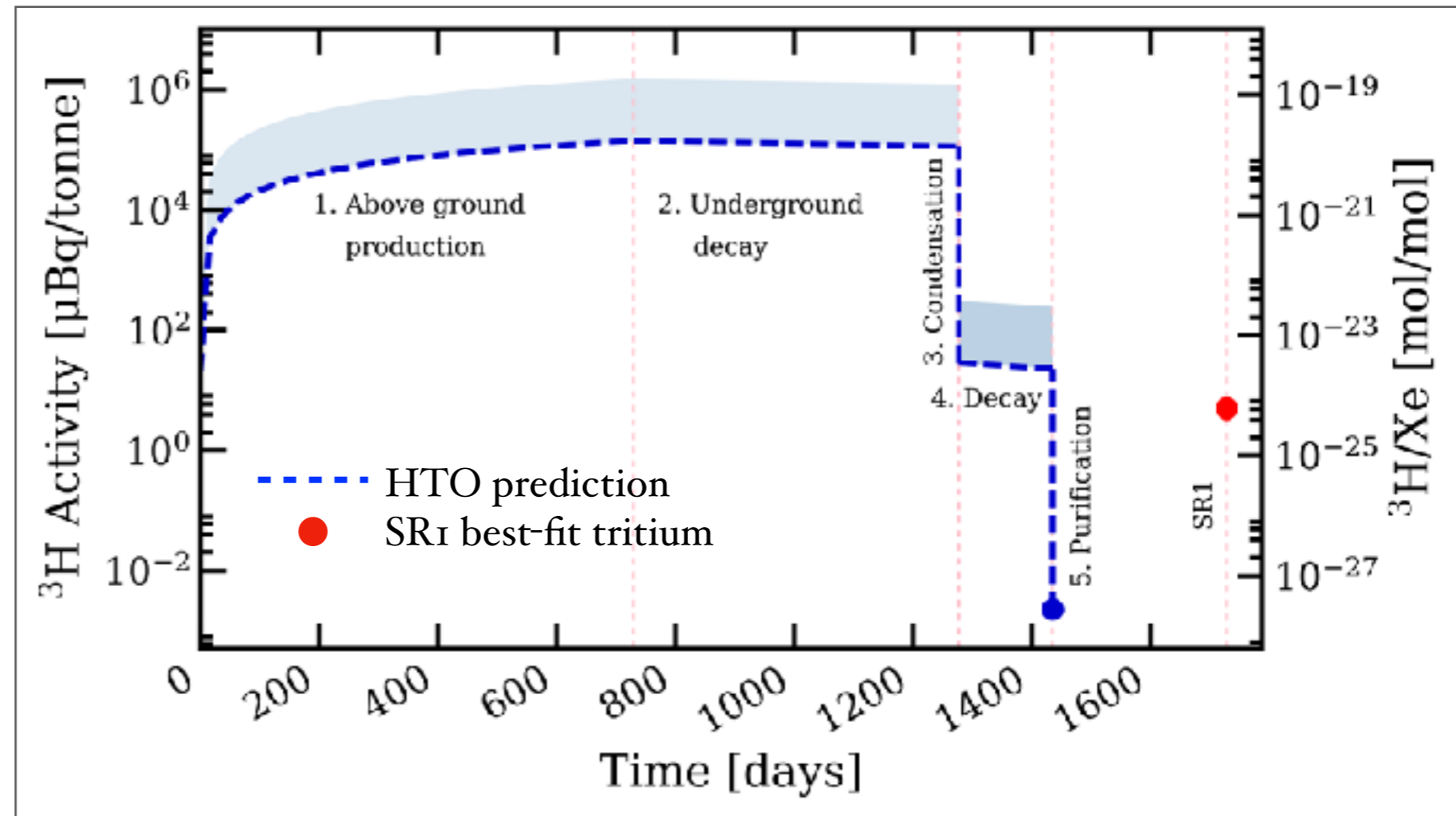
**fewer than 3 tritium atoms per kg of xenon!**

Cosmogenic activation of xenon: ~32 tritium atoms/kg/day (Zhang, 2016)

1 ppm water in bottles implies tritium forms predominately HTO.

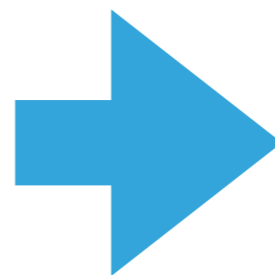
Efficient removal (99.99%) in purification system (SAES getter with hydrogen removal unit)

## Tritiated water (HTO)



(note: tritium from activation while underground is negligible.)

Expected concentration more than 100x smaller than measured



From purification and handling, this component seems unlikely.

## 2. Atmospheric Abundance in Materials

What about T emanating from materials in equilibrium with removal?

\*Hydrology measurements from IAEA nuclear database

HTO:H<sub>2</sub>O concentration\*  $5-10 \times 10^{-18}$  mol/mol

HT:H<sub>2</sub> concentration  Assuming same concentration as for H<sub>2</sub>O

Required (H<sub>2</sub>O + H<sub>2</sub>):Xe  
concentration to explain

**60-120 ppb**

Tritiated molecules can emanate into LXe target from water and hydrogen in detector materials in the form of **HTO** and tritiated hydrogen (**HT**).  
emanation in equilibrium with removal.

**But...**

**H<sub>2</sub>O**  
H<sub>2</sub>O in XENON1T: O(1) ppb,  
otherwise can not detect light

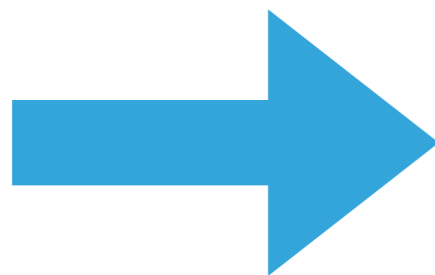
**H<sub>2</sub>**  
O<sub>2</sub> in XENON1T: <1 ppb, otherwise  
can not drift electrons

H<sub>2</sub> ~100 ppb? -> ~100x higher than  
O<sub>2</sub> possible?



## Many unknowns about tritium in a cryogenic LXe environment

- Radiochemistry, particularly isotopic exchange (formation of other molecules?)
- Diffusion properties of tritiated molecules
- Desorption and emanation
- For HT, no direct measure of either abundance or H<sub>2</sub> concentration.



柿内さんのトーク！

**We can neither confirm nor exclude the presence of tritium.**

# Ar37? (時間があれば)

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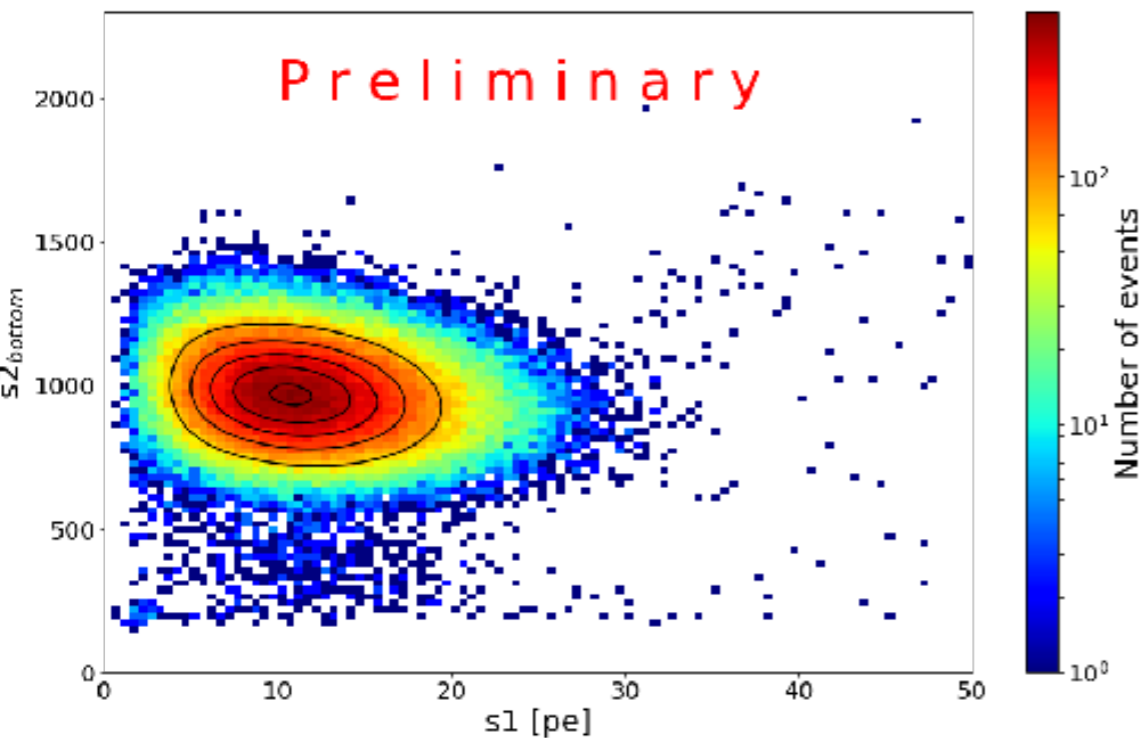
1. In-situ production?

2. Atmospherically abundant?

3. Air leak?

$^{37}\text{Ar}$  K-electron capture to the ground state of  $^{37}\text{Cl}$  ( $^{37}\text{Ar} \rightarrow ^{37}\text{Cl} + \nu e$ )

- Half-lifetime of 35 days & 2.8 keV energy in X-rays & Auger e<sup>-</sup>s
  - Calibration with  $^{37}\text{Ar}$  performed in XENON1T at the end of SR2
- good understanding of the detector at those energies



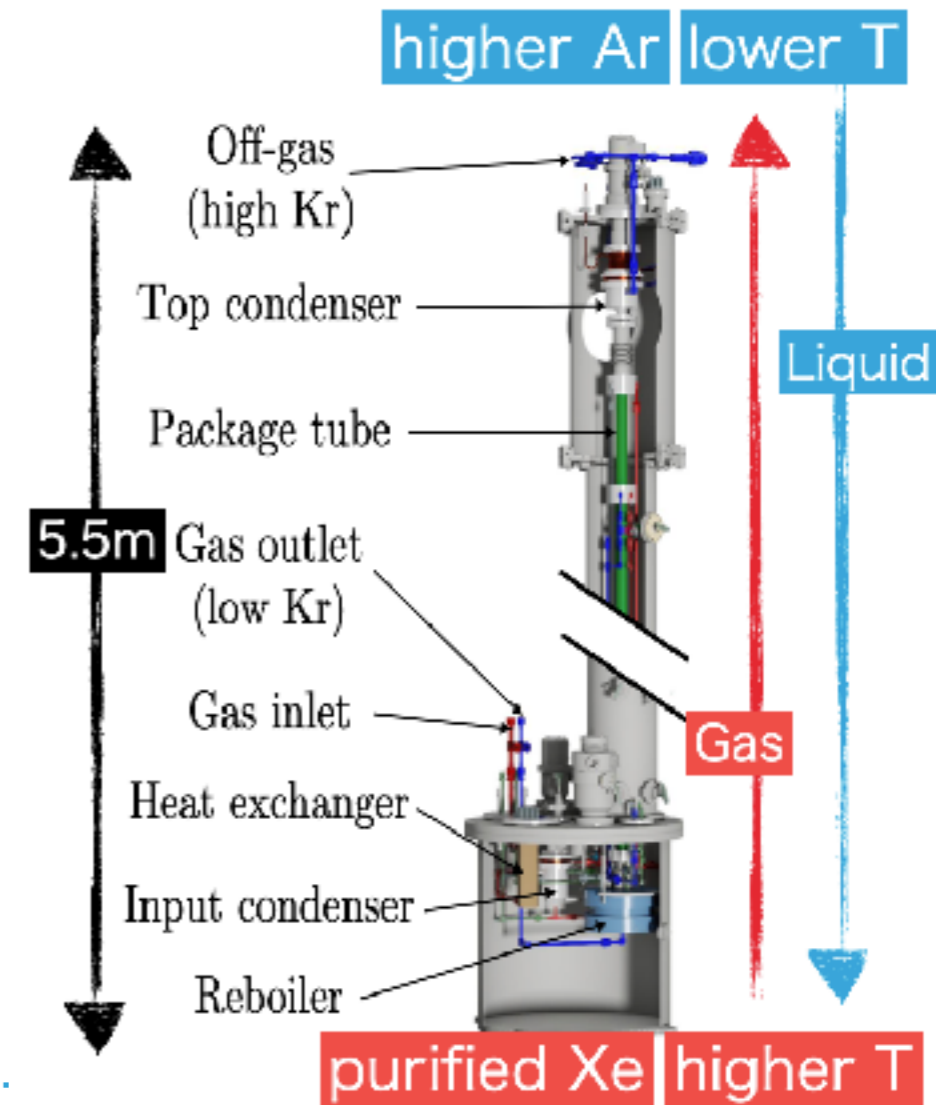
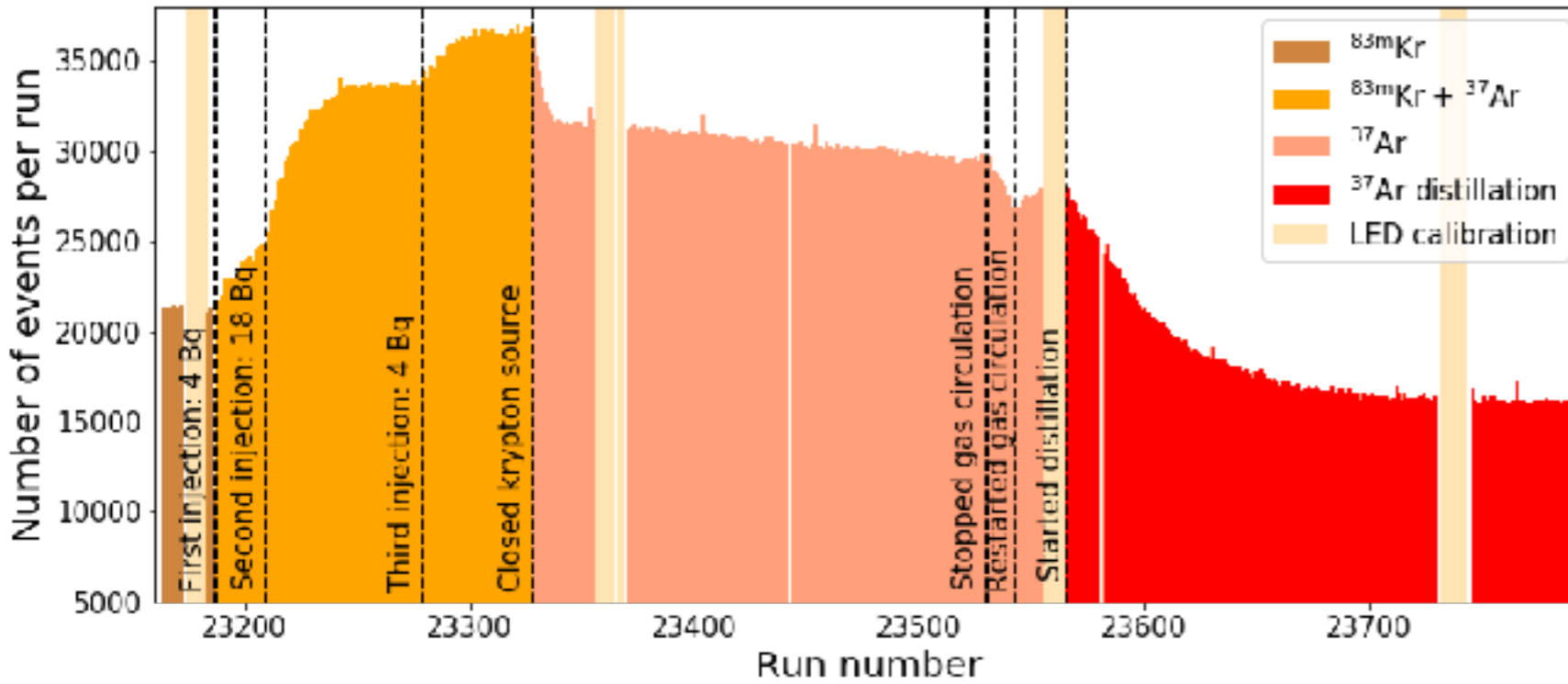
**Table 2.**  $^{37}\text{Ar}$  decay modes and released energy

Decay mode	Energy release, keV	Branching ratio for $^{37}\text{Ar}$ decays through a given mode	Energy release per event of $^{37}\text{Ar}$ decay, keV/decay
<i>K</i> capture	2.8224	$0.9017 \pm 0.0024$	$2.5450 \pm 0.0068$
<i>L</i> capture	0.2702	$0.0890 \pm 0.0027$	$0.0240 \pm 0.0007$
<i>M</i> capture	0.0175	$0.0093^{+0.0006}_{-0.0004}$	0.0002
IB 1s	325 (average)	$\sim 0.0005$	$0.16 \pm 0.02$
IB 2s	325 (average)	$\sim 0.00007$	$0.021 \pm 0.002$
IB <i>p</i>	$\sim 10$ (average)	$\sim 0.00007$	$\sim 0.0007$
Sum			$2.751 \pm 0.021$

*Physics of Atomic Nuclei volume 70, pages 300–310(2007)*

Possible  $^{37}\text{Ar}$  contributions:

1. Its presence in the xenon gas before filling,
2. A possible air leak that could provide a constant source of argon.
3. In-situ production: neutron reactions with  $^{36}\text{Ar}$  or  $^{40}\text{Ca}$



## Possible $^{37}\text{Ar}$ contributions:

1. Its presence in the xenon gas before filling

→ Argon is strongly reduced by 90 days of distillation & decay: combined time constant  $\sim 1.8$  days,  $>10$  times faster than decay

2. A possible air leak that could provide a constant source of argon. →  $\sim 3$  L/day air leak

→ ruled out by  $^{84}\text{Kr}$  measurement (RGMS,  $< 1$  ppt increase/yr =  $\sim 1$  L/year air leak)

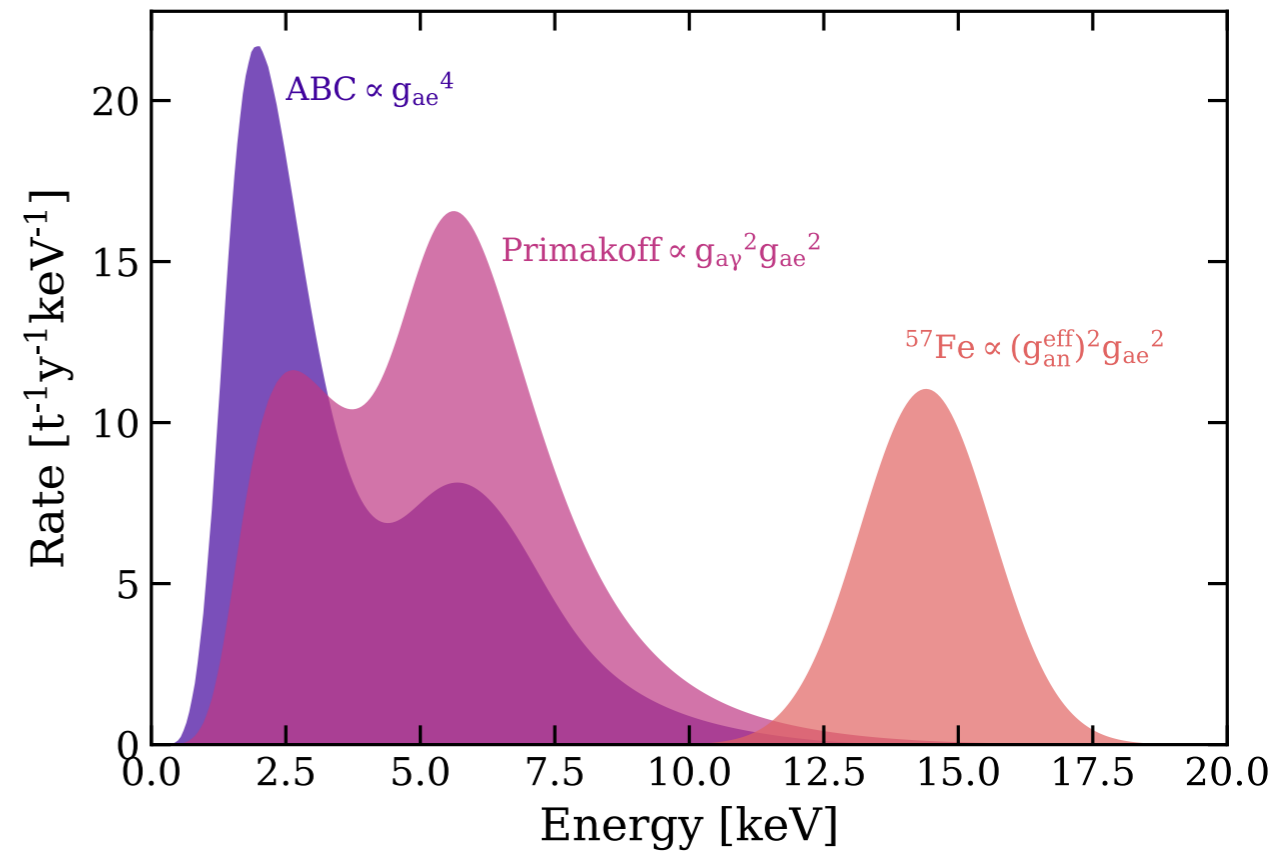
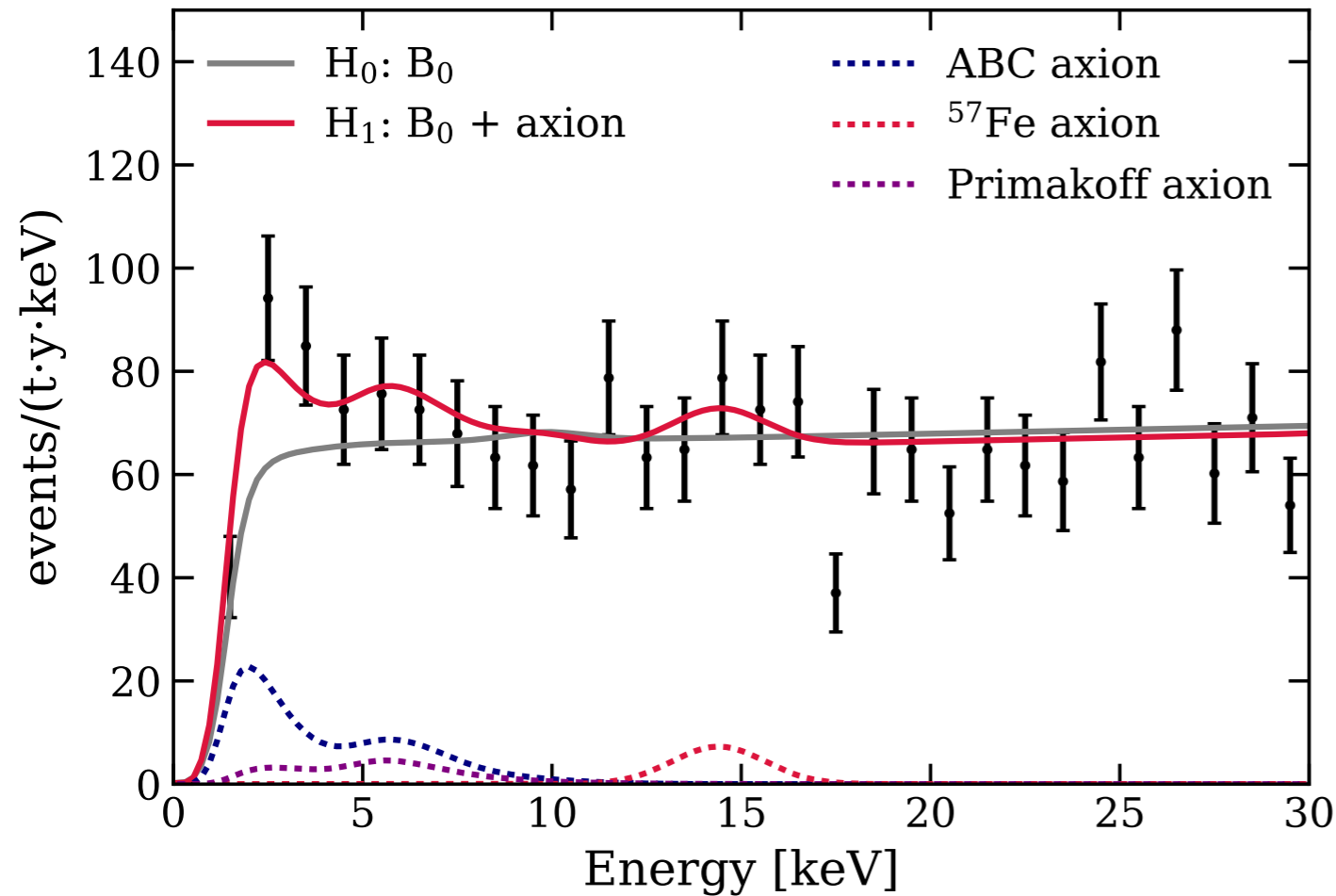
3. In-situ production: neutron reactions with  $^{36}\text{Ar}$  or  $^{40}\text{Ca}$

→ Negligible

# Signal Interpretation?

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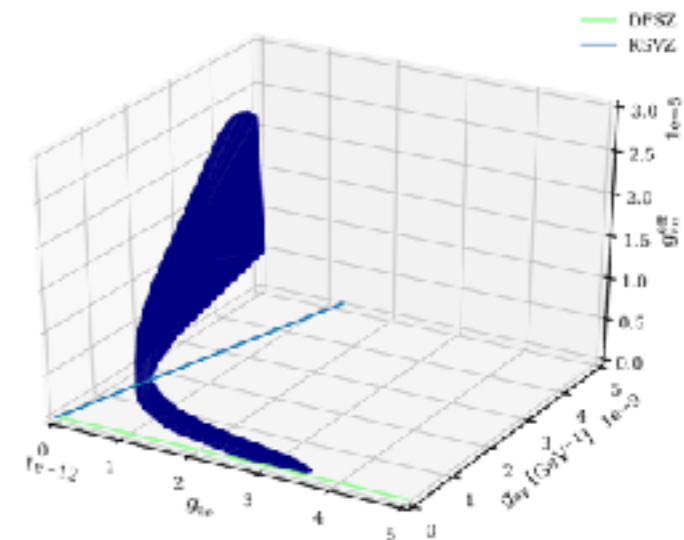


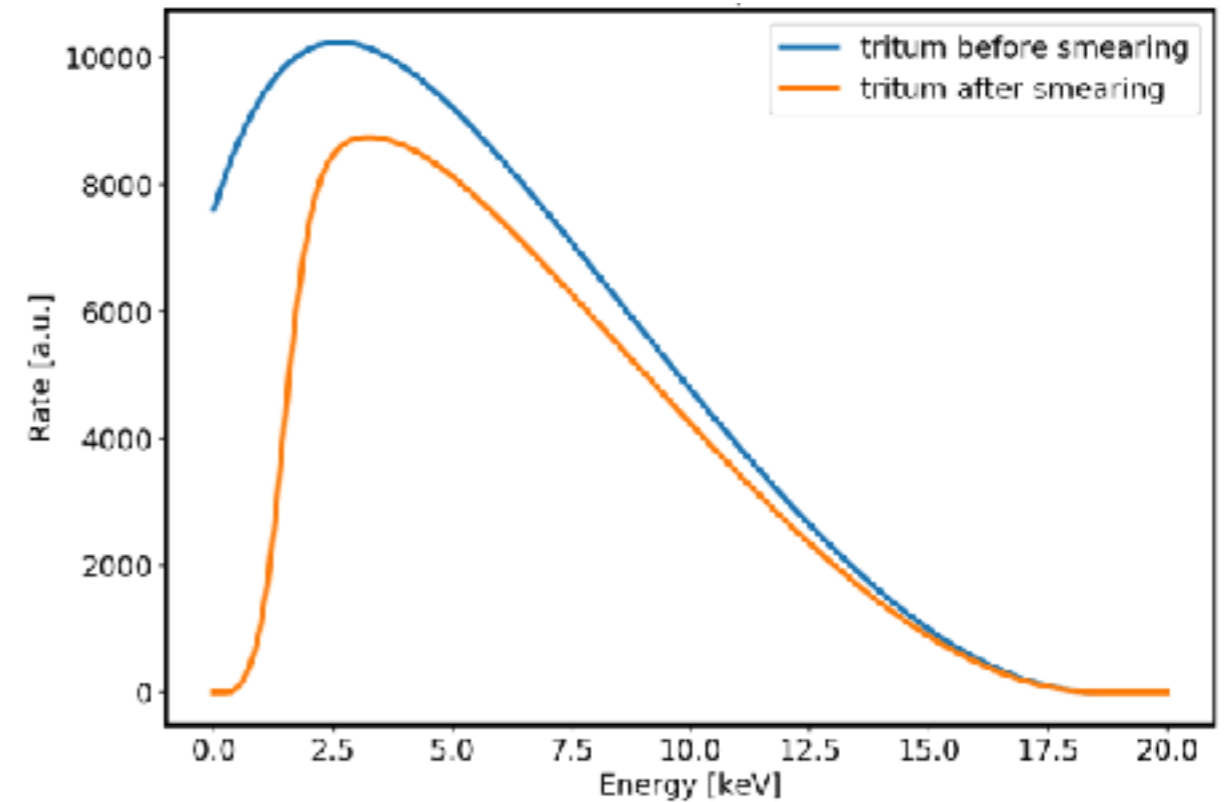
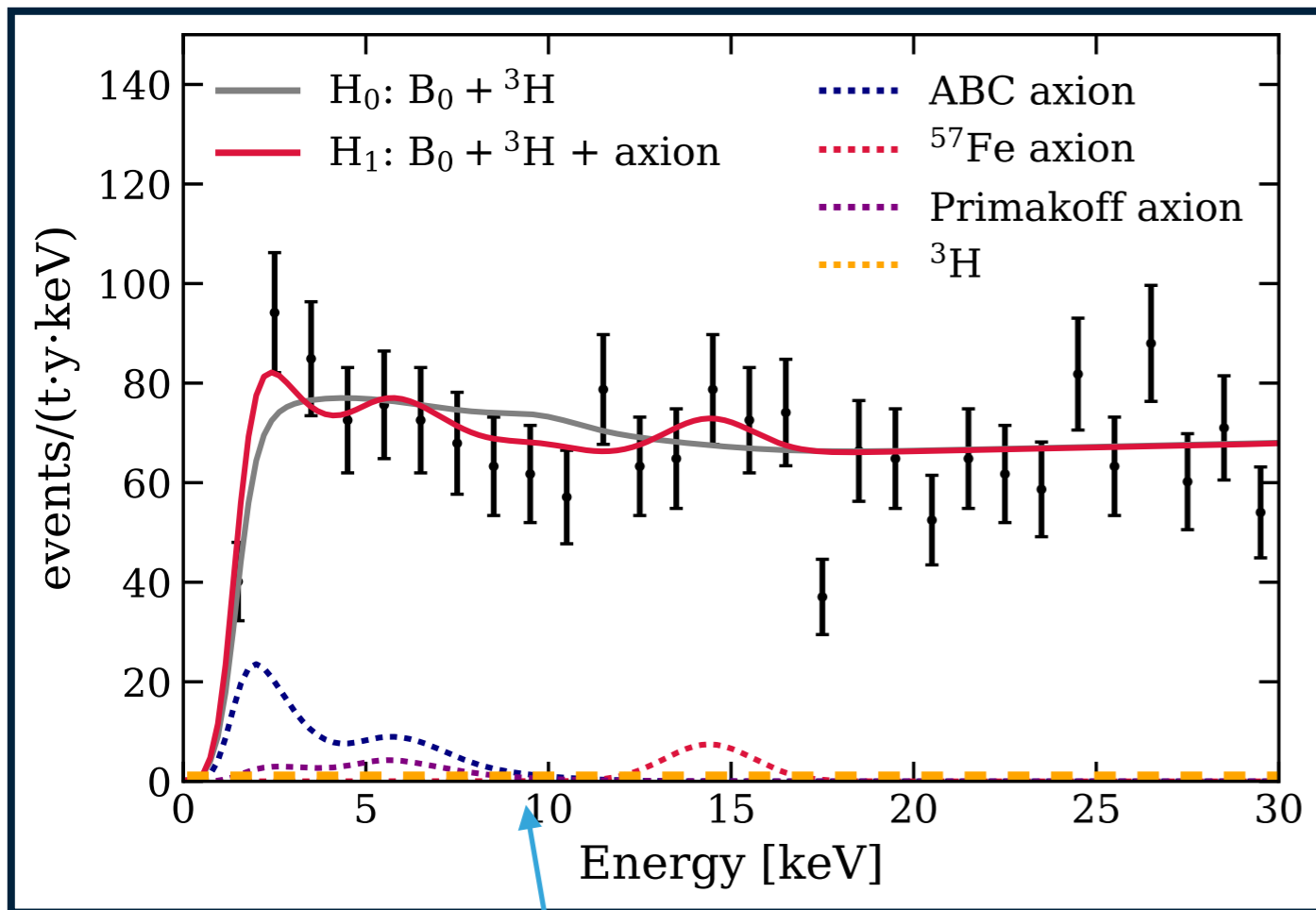


- All the three flux components are considered completely independent of each other
- Parameters of interest in the Profile Likelihood =  $g_{\text{ae}}$  VS.  $g_{\text{ae}}g_{\text{ay}}$  VS.  $g_{\text{ae}}g_{\text{an}}^{\text{eff}}$
- Significance determined using toy-MC methods

$g_{\text{ae}}$  VS.  $g_{\text{ae}}g_{\text{ay}}$  VS.  $g_{\text{ae}}g_{\text{an}}^{\text{eff}}$

**Axion favored over background-only at  $3.5\sigma$**



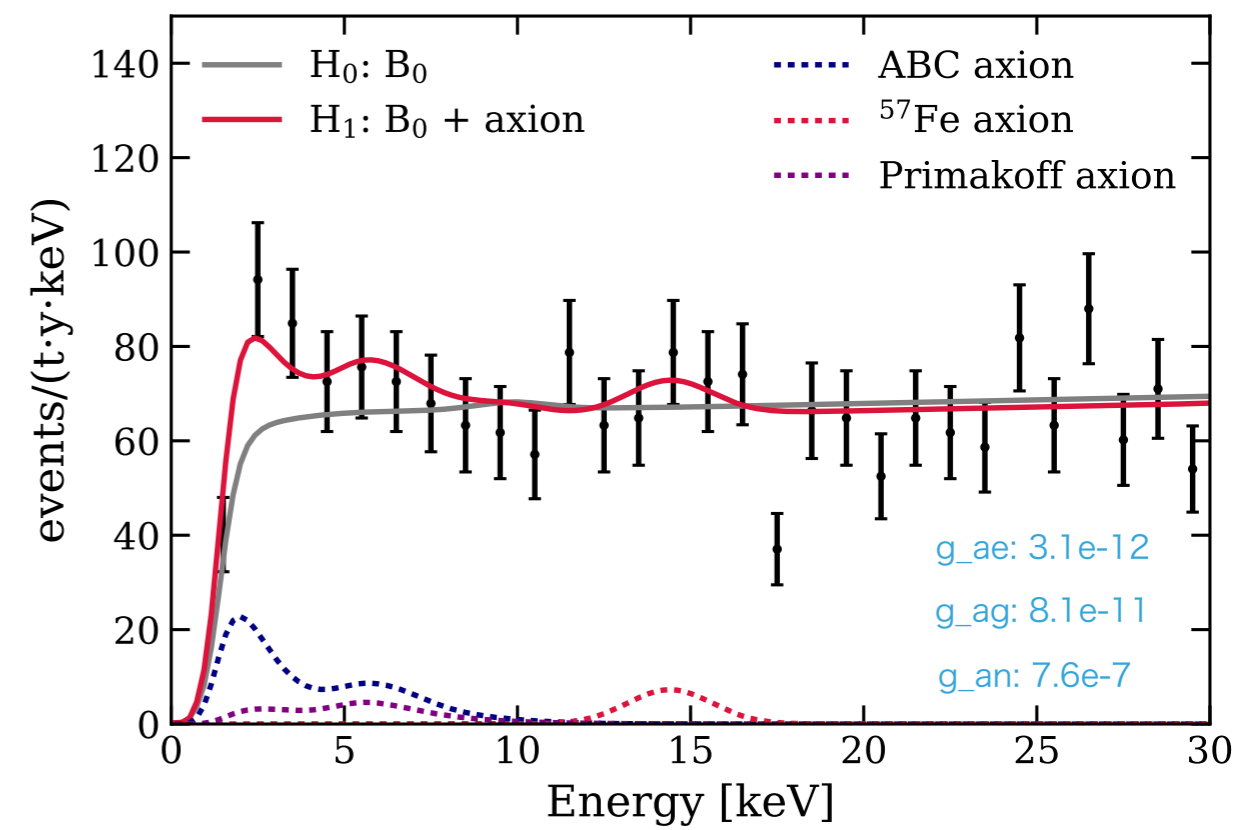
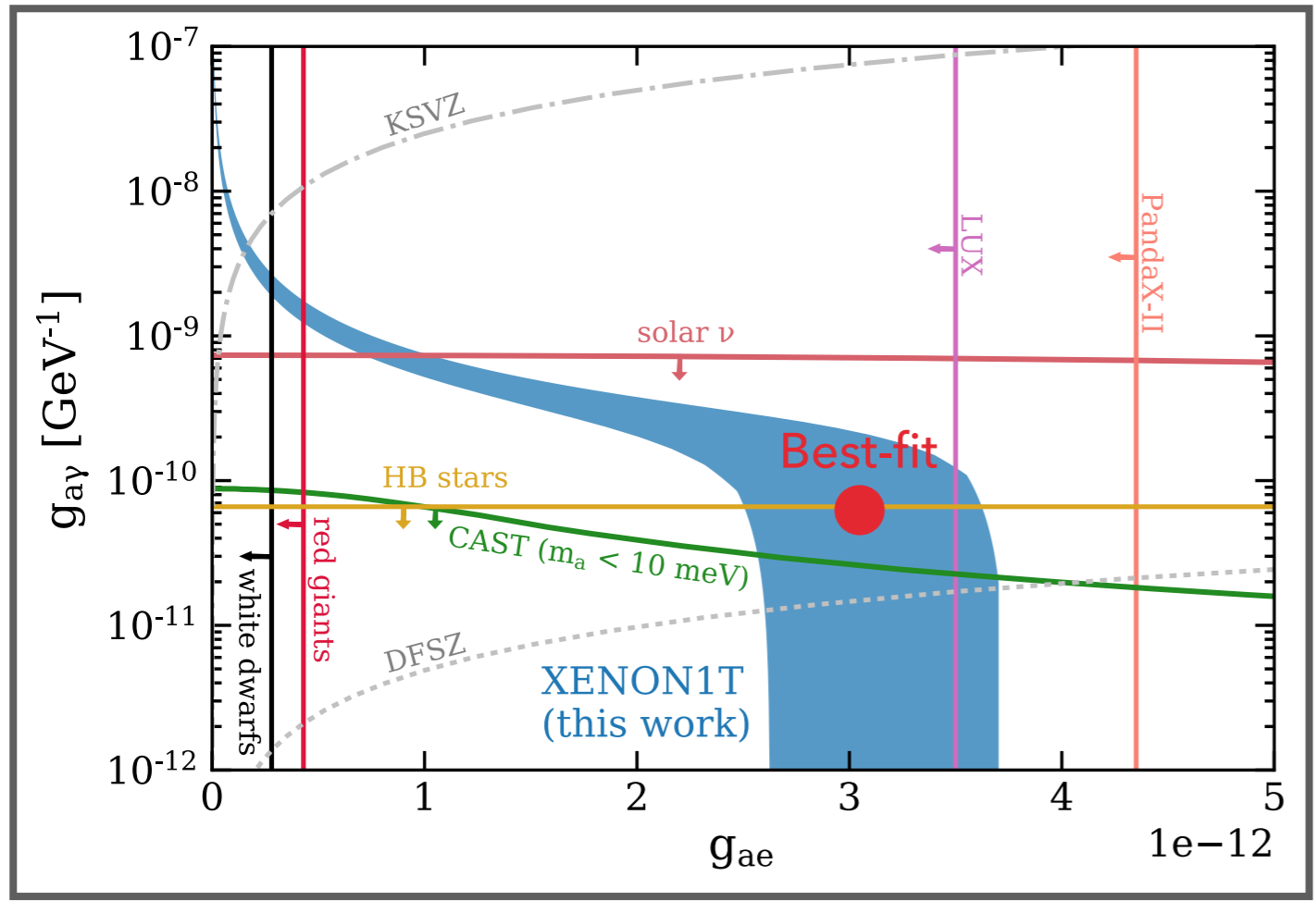


**Axion +  ${}^3\text{H}$  favored over  ${}^3\text{H}$  hypothesis at  $2.1\sigma$**

When both axion and tritium are included in the fit, the **best-fit of tritium is zero** — **in favor of axions.**

Parameters of interest in the profile likelihood

$$g_{ae} \text{ VS. } g_{ae}g_{a\gamma} \text{ VS. } g_{ae}g_{an}^{\text{eff}}$$

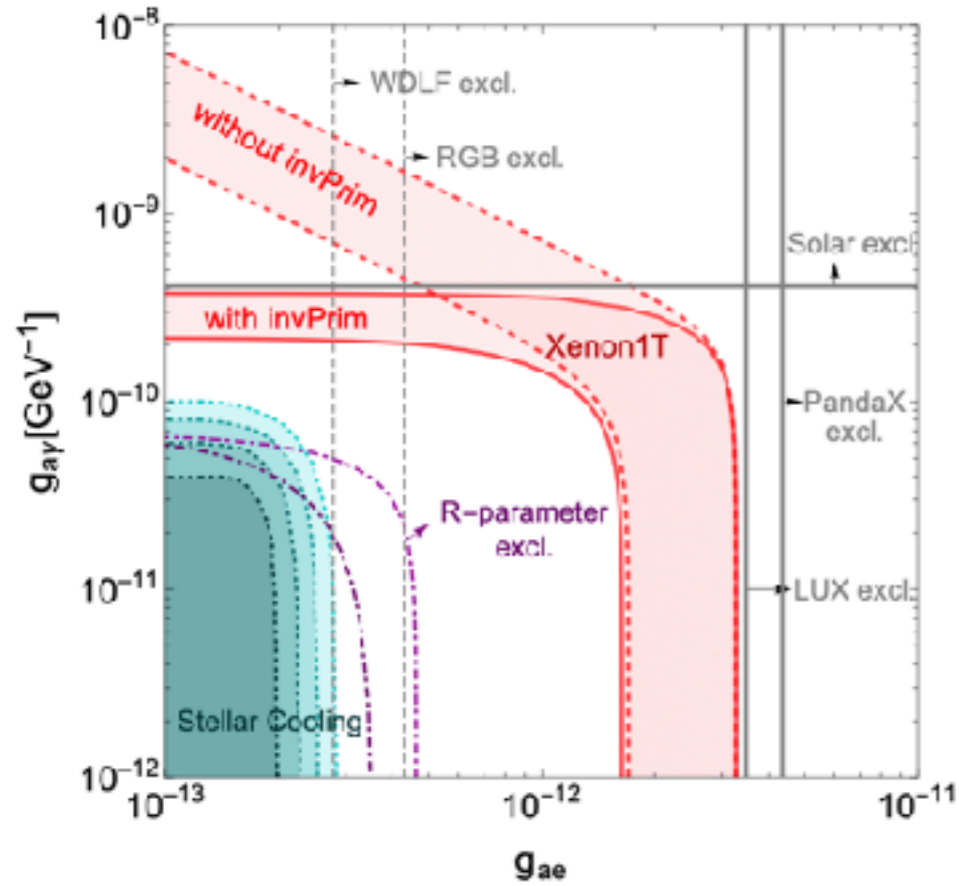


ABC and Primakoff components are both low-energy signals, the favored region is anti-correlated in this space

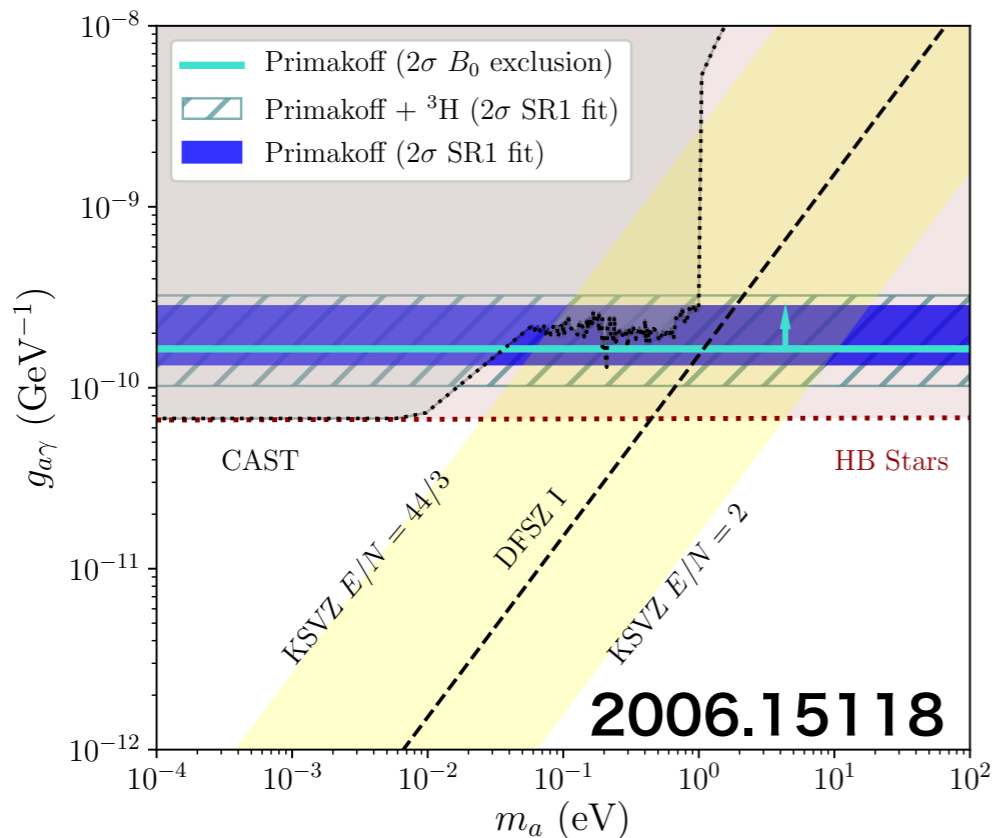
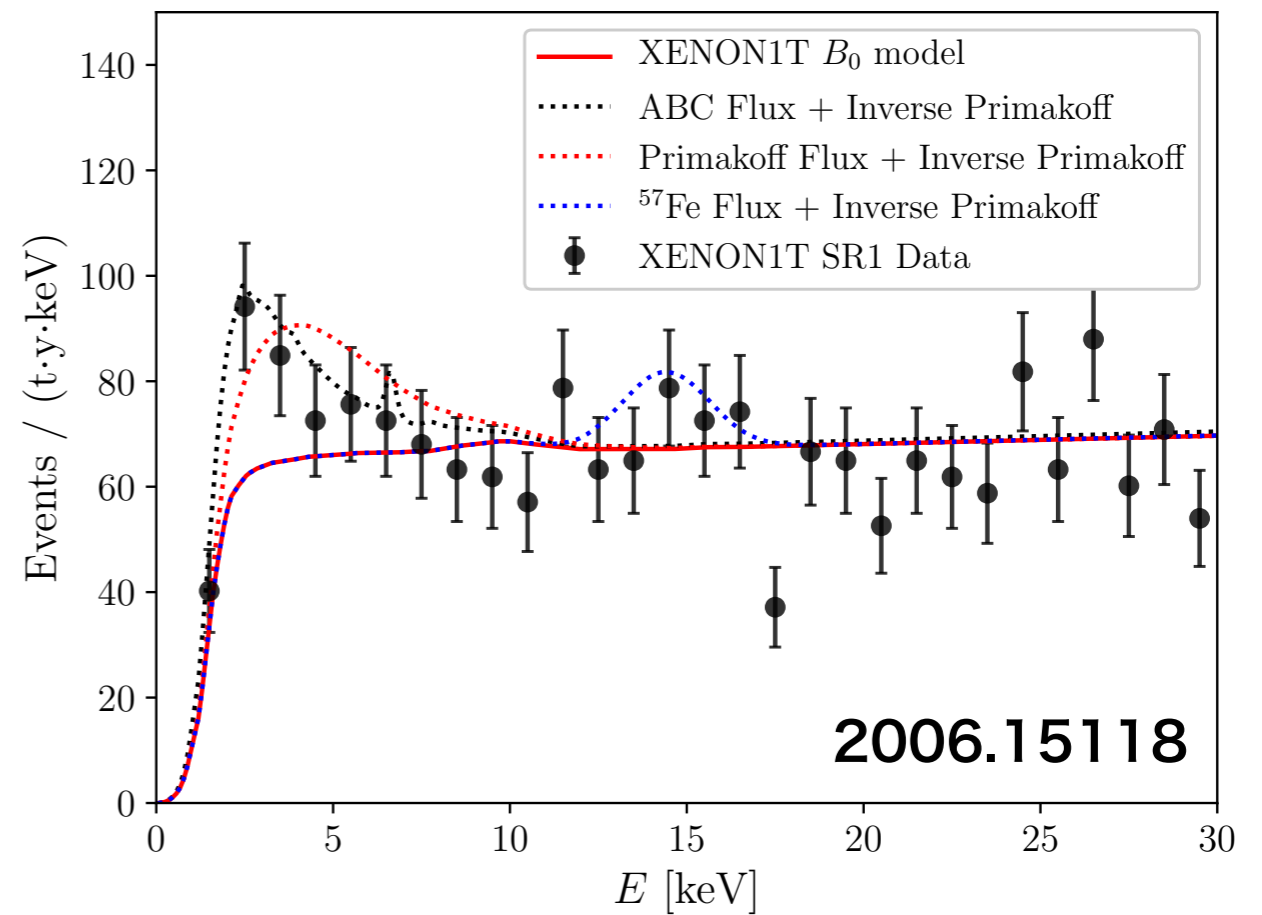
→ suggests either a non-zero ABC component or non-zero Primakoff component.

In tension with astrophysical constraints from stellar cooling bounds from the horizontal branch stars and red giants

2006.14598



$$g_{a\gamma} = 2.7 \cdot 10^{-10} \text{ GeV}^{-1}, g_{ae} = 4.5 \cdot 10^{-12}, g_{an}^{\text{eff}} = 10^{-6}$$

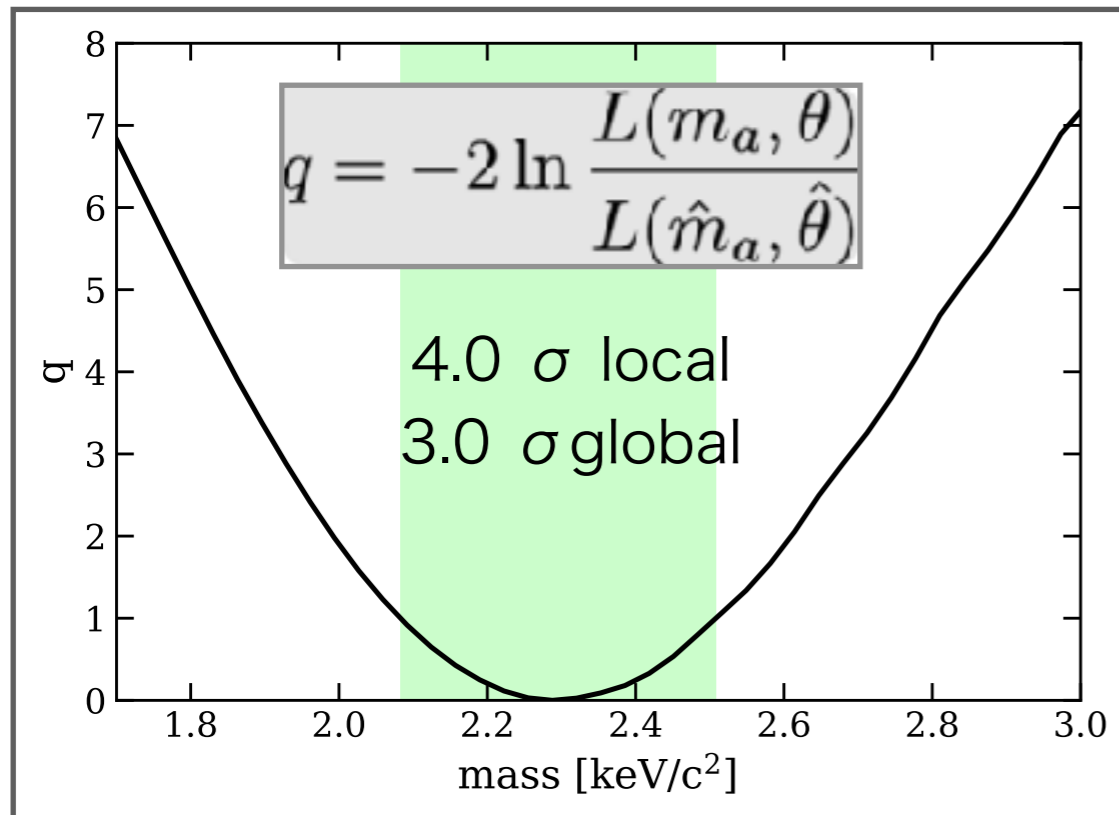


Considering inverse Primakoff process can weaken the tension with stellar cooling constraint

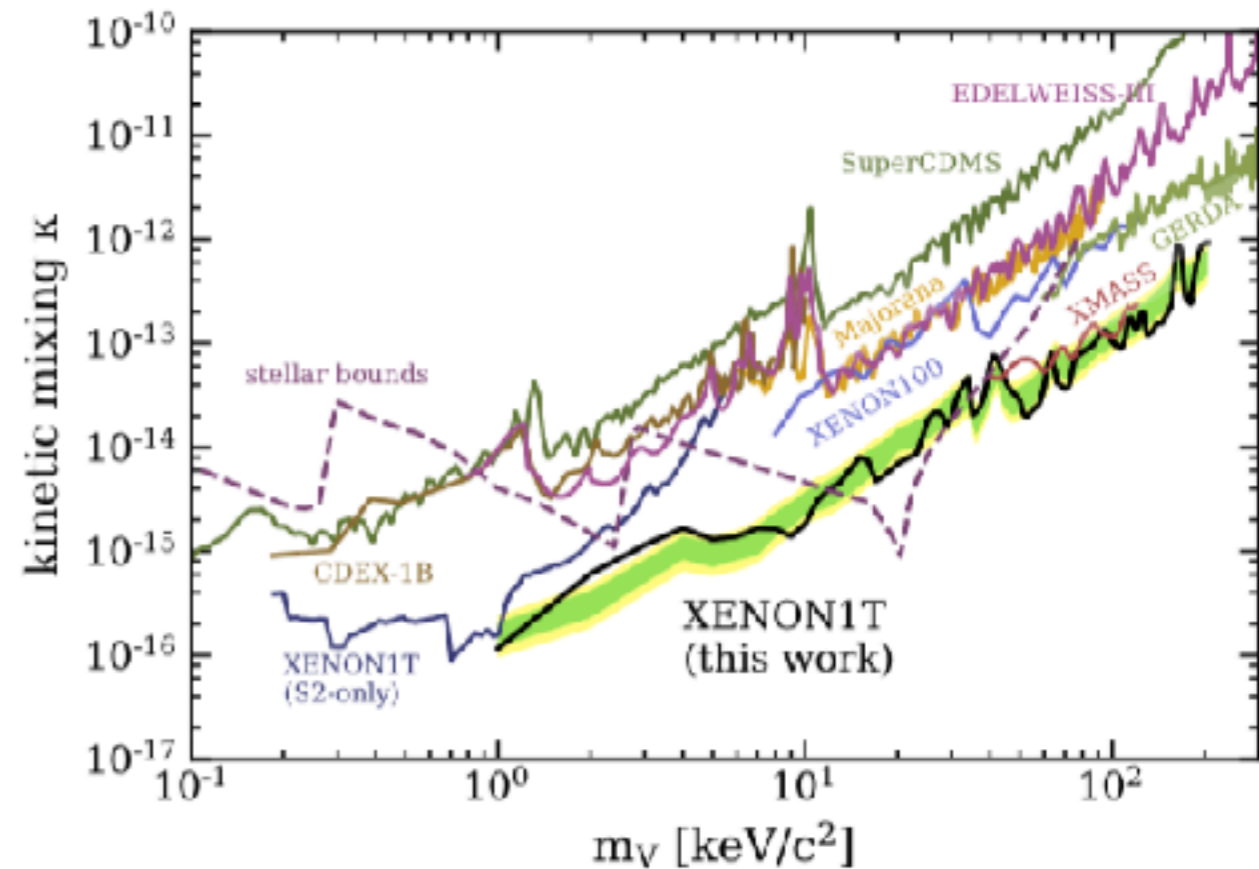
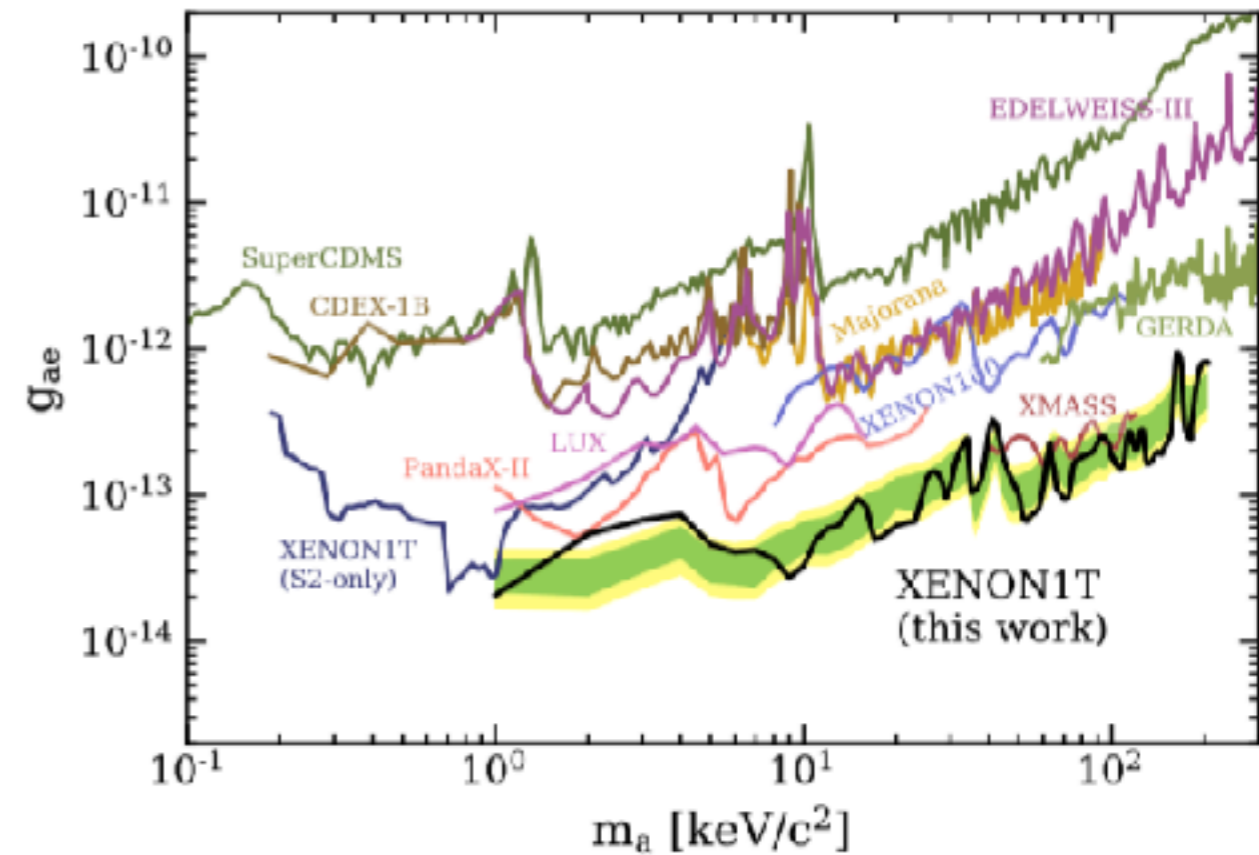
XENON1T has a good sensitivity also for  $g_{a\gamma}$ !

# Axion-like Particle / Dark Photon

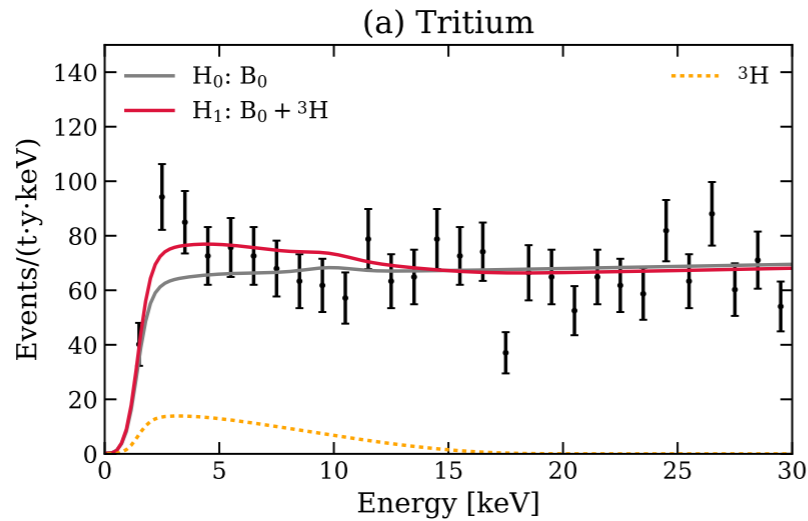
Fitting a mono-energetic peak to the excess:  $2.3 \pm 0.2$  keV



Best fit:  $\sim 60$  events/tonne/year  
4.0  $\sigma$  local significance  
**3.0  $\sigma$  (global, considering look-elsewhere effect).**

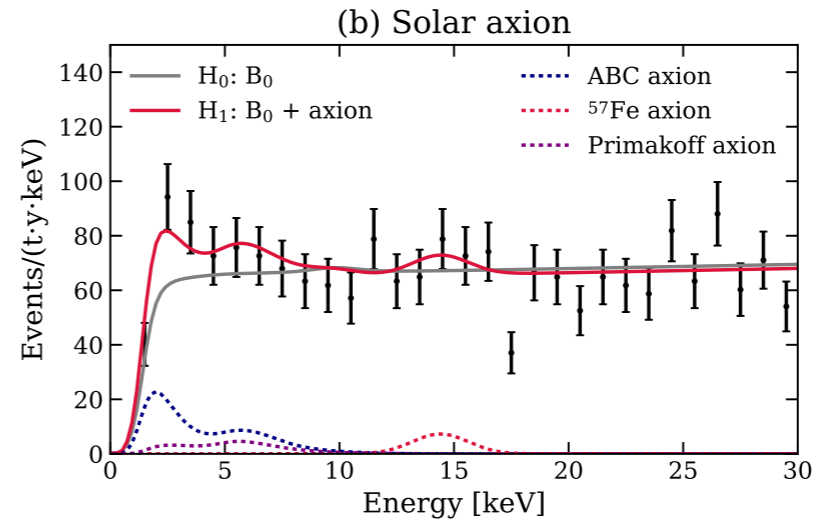


**Tritium**  
 favored over  
 background-only at  
**3.2 $\sigma$**

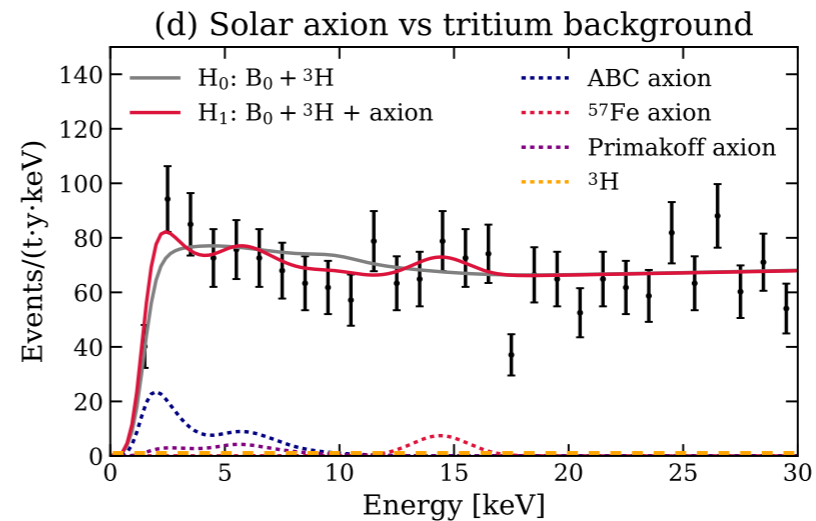


**Solar axion**  
 favored over  
 background-only at  
**3.5 $\sigma$**

$g_{ae}: 3.1e-12$   
 $g_{ag}: 8.1e-11$   
 $g_{an}: 7.6e-7$

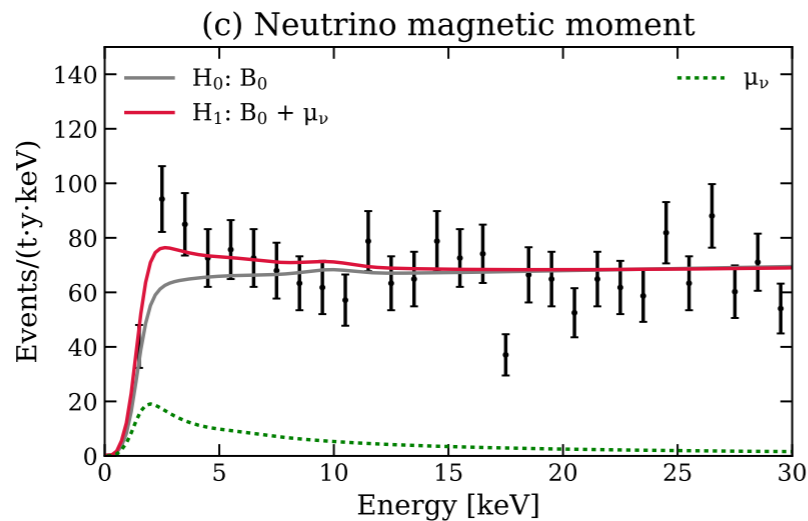


**Axion +  $^3\text{H}$**  favored  
 over  $^3\text{H}$  hypothesis at  
**2.1 $\sigma$**

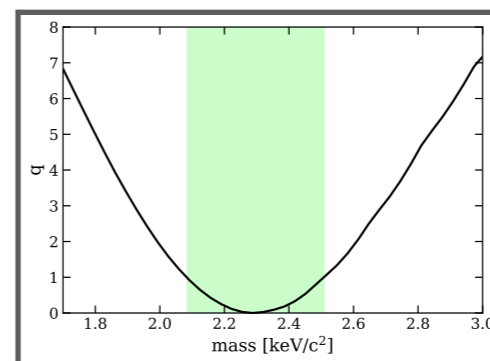


**Neutrino magnetic moment** (see backup slides) favored over background-only at  
**3.2 $\sigma$**

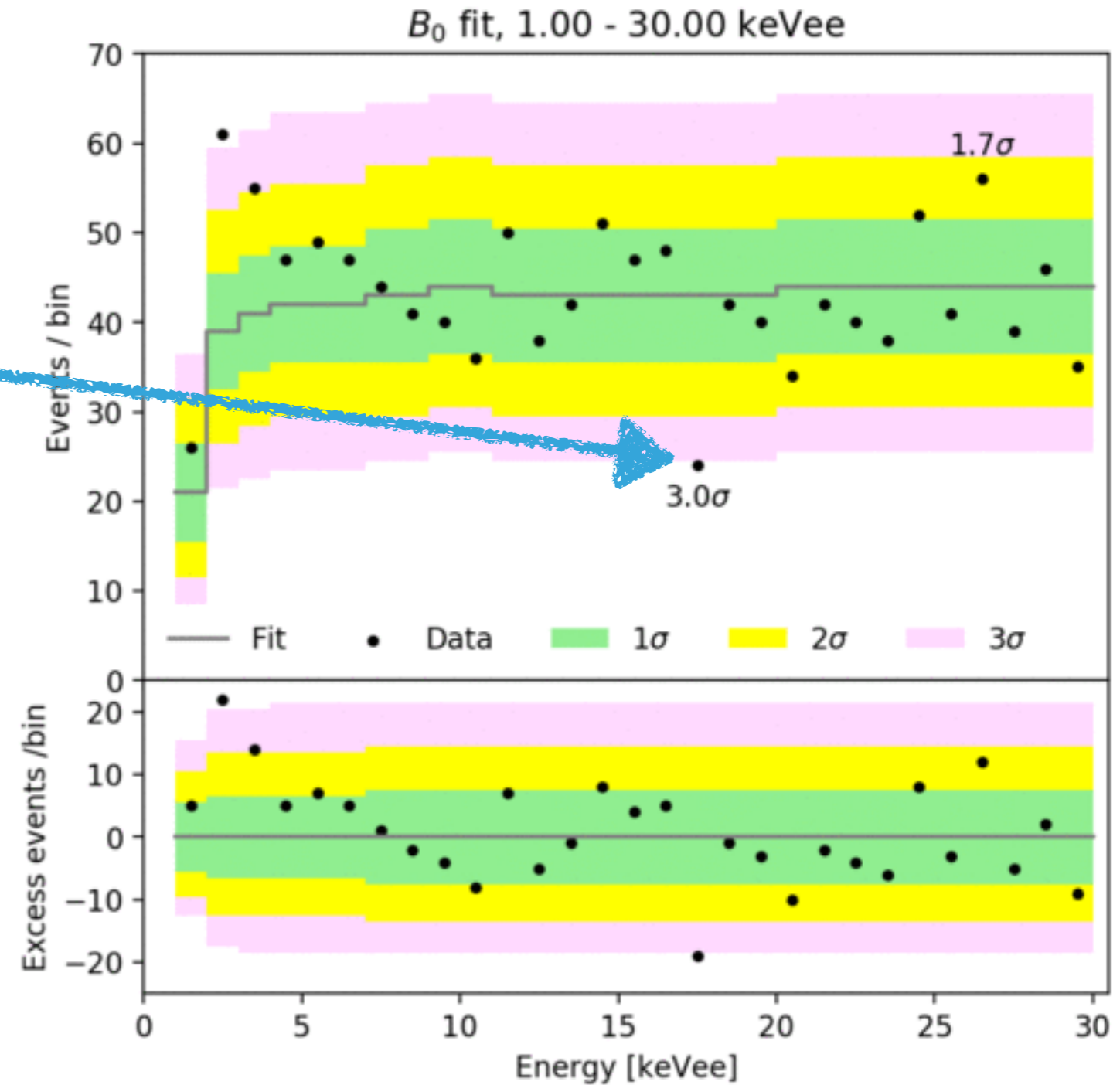
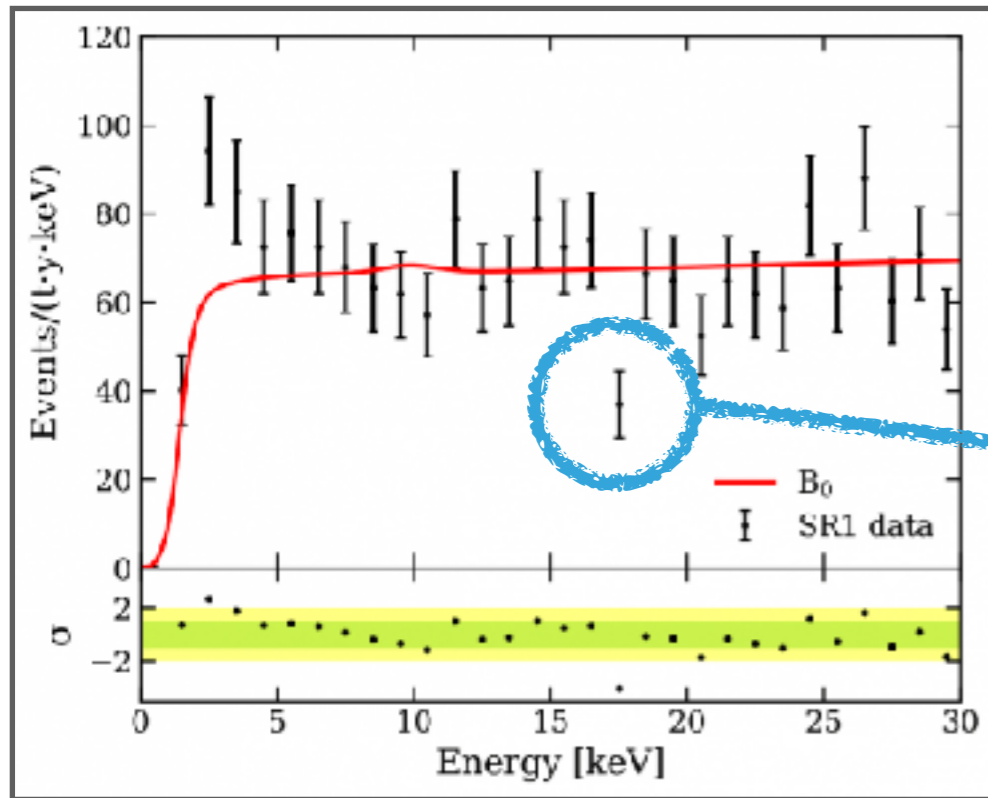
$\mu_\nu: 2.3e-11$



**Monoenergetic peak at 2.3 +/- 0.2 keV**  
 favored over background-only at **3.0 $\sigma$**   
 (global)



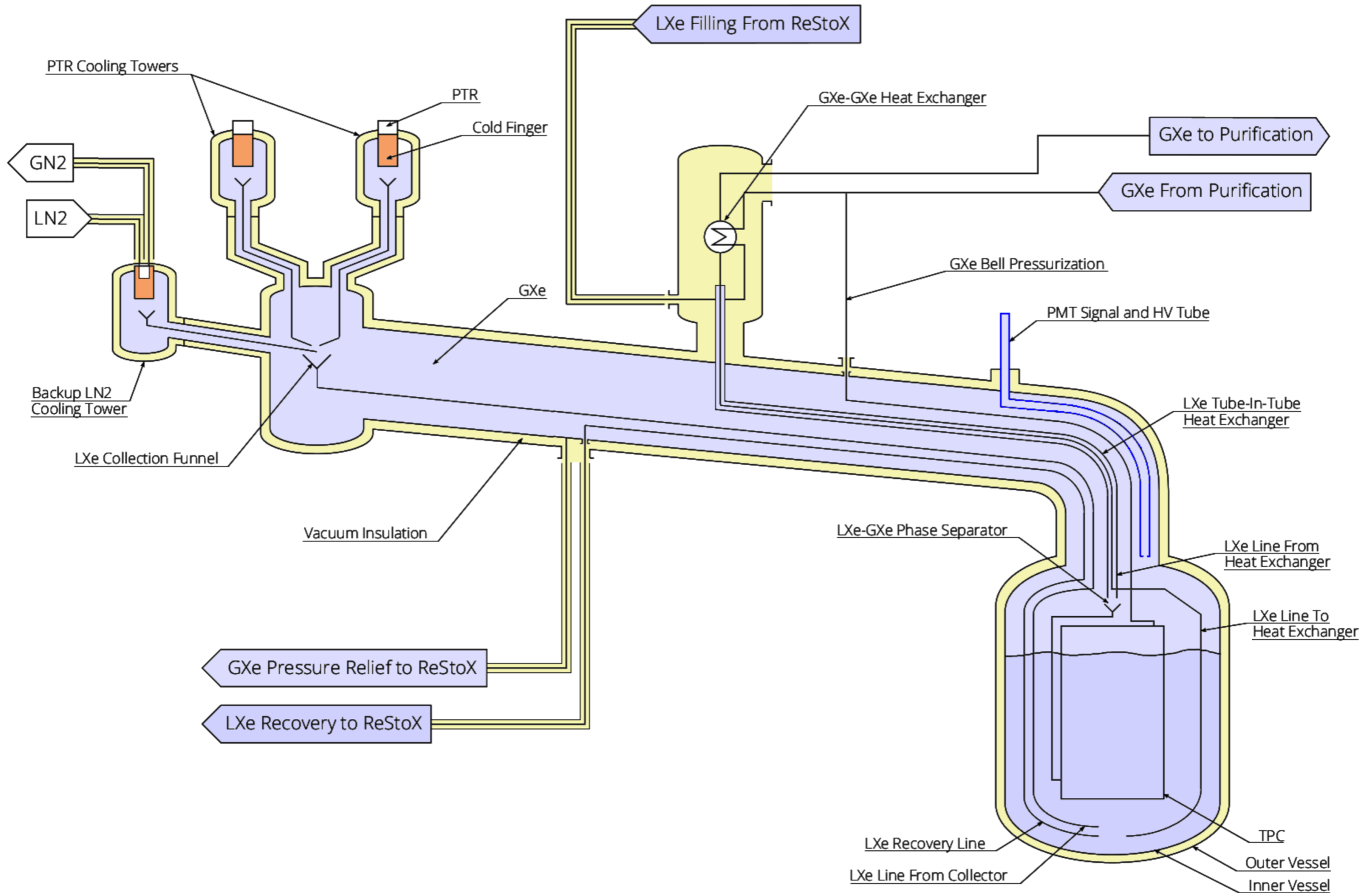
# Back Up



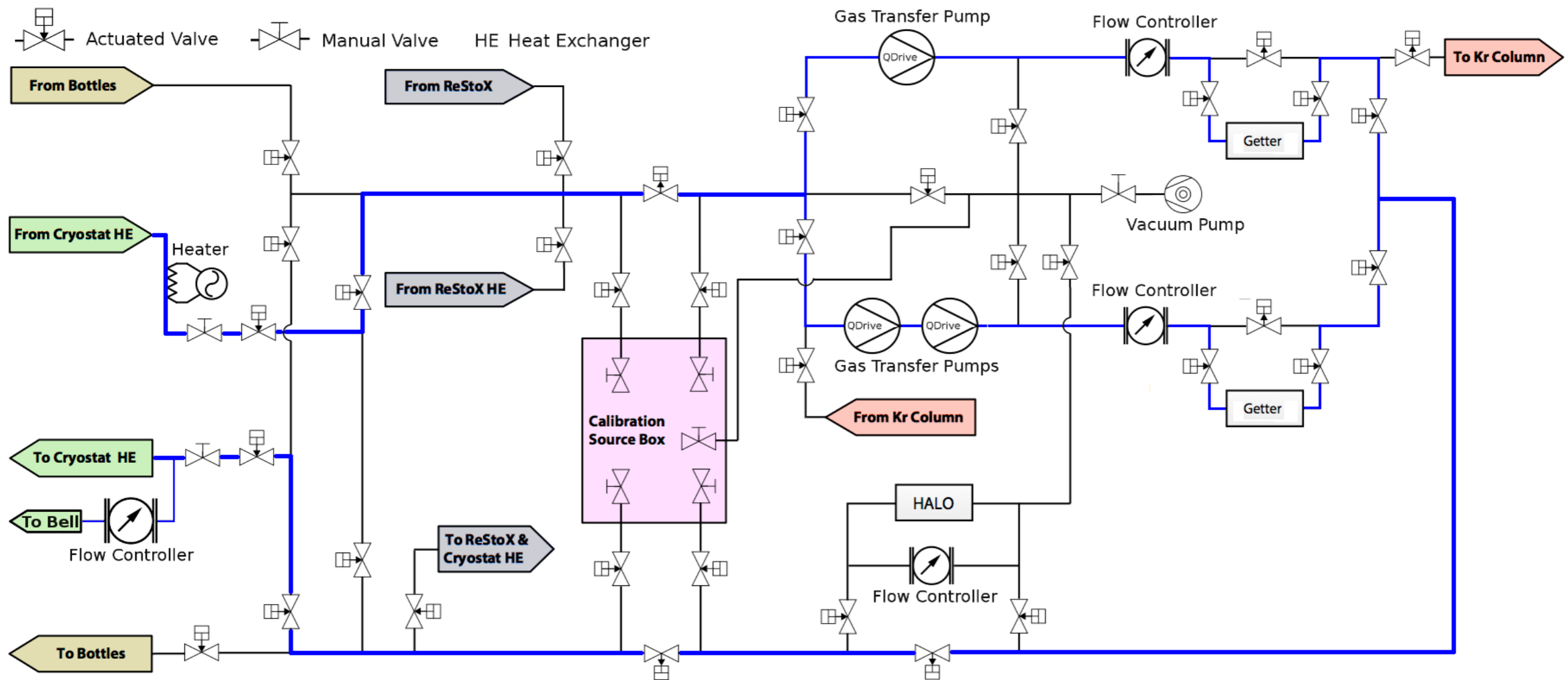
statistical fluctuation? (see 17 keV dip)

Note: we use an unbinned profile likelihood analysis





**Fig. 7** The cryogenic system of XENONnT: cooling is provided by means of three redundant cold heads (two pulse-tube refrigerators (PTR), 1 LN<sub>2</sub>), installed on individual cooling towers located outside of the water shield. The liquefied xenon runs back to the main cryostat in a 6 m long vacuum-insulated cryogenic pipe, through which all connections to the TPC are made with the exception of the cathode bias voltage which is not shown in the figure. The connections to the systems for xenon purification and storage (ReStoX) are also shown. Figure not to scale.



**Fig. 8** Piping and instrumentation diagram (P&ID) of the XENON1T purification system. The system also serves as the main interface to the other components of the gas-handling system (see figure 9) and allows the insertion of short-lived isotopes for calibration. Some instrumentation such as temperature and pressure sensors, as well as several access ports are omitted for clarity. The path of the xenon gas during standard purification is indicated in blue.

Two possible  $^{37}\text{Ar}$  contributions:

1. Its presence in the xenon gas before filling,
2. A possible air leak that could provide a constant source of argon.

1. Its presence in the xenon gas before filling,

- Removal time in distillation is  $\sim 1.8$  day, directly demonstrated at 1T using a dedicated  $^{37}\text{Ar}$  source
- We had  $\sim 90$  days of the online  $^{85}\text{Kr}$  distillation before SR1 (22 orders of magnitude reduction)
- The isotopic abundance of  $^{37}\text{Ar}$  is  $\sim 10^{-20}$

Even if there were 1 ppm of  $^{\text{nat}}\text{Ar}$  in the xenon originally, the  $^{37}\text{Ar}$  concentration would have been reduced to a negligible level ( $\sim 10^{-48}$  mol/mol)

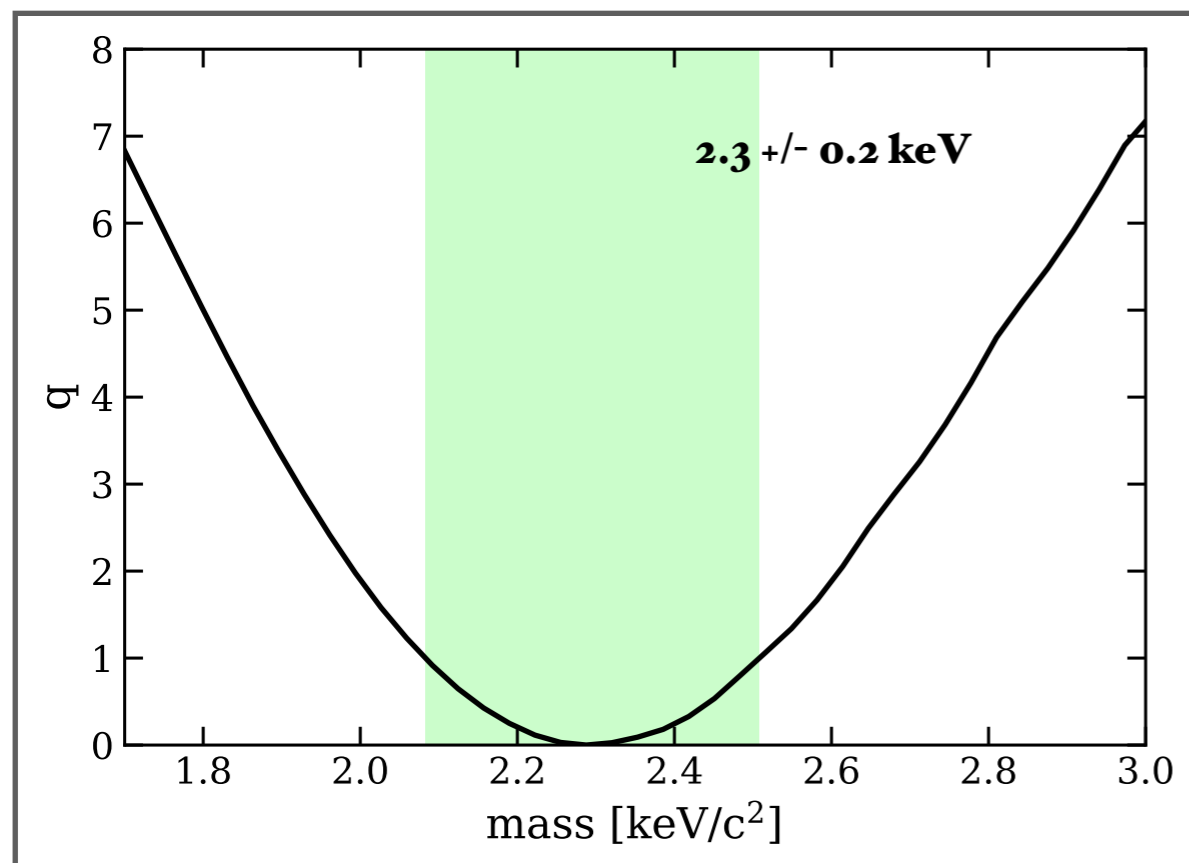
## 2. A possible air leak that could provide a constant source of argon.

Require:  $\sim 10^{-4}$  kg of argon per day, corresponding to a total air leak of  $\sim 3$  L/day.

→ Ruled about by the  $^{\text{nat}}\text{Kr}$  concentration, which increased by  $< 1$  ppt/year during SR1 as informed by RGMS measurements

1-ppt/year increase in  $^{\text{nat}}\text{Kr}$  would correspond to an air leak of  $\sim 1$  L/year in XENON1T.

Also TPC would not work in such a leaky condition because of O<sub>2</sub>...



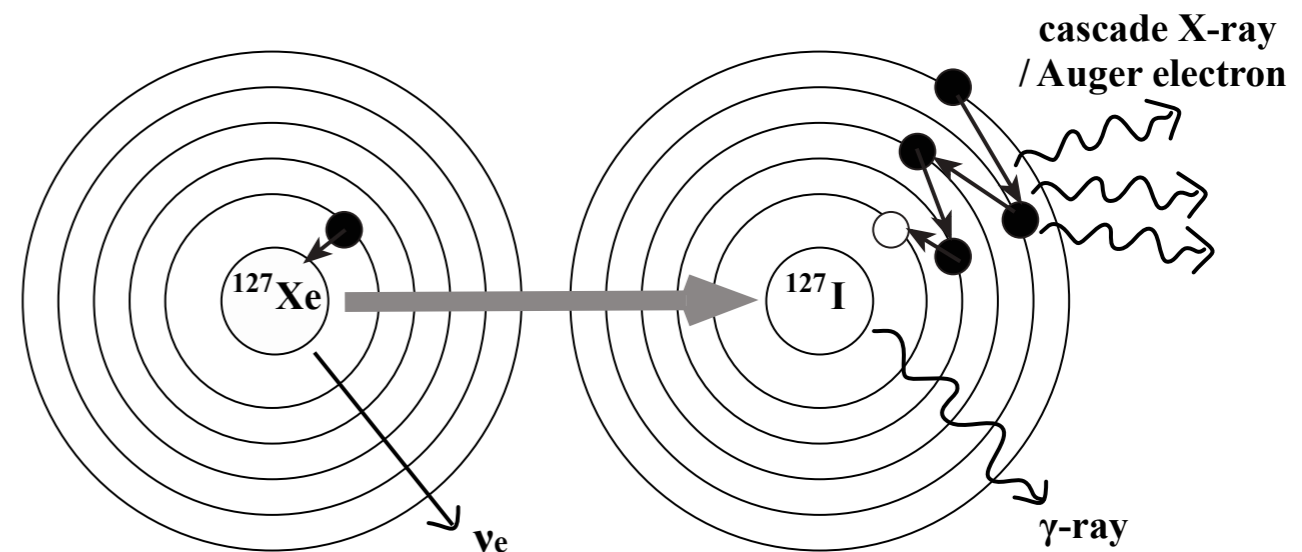
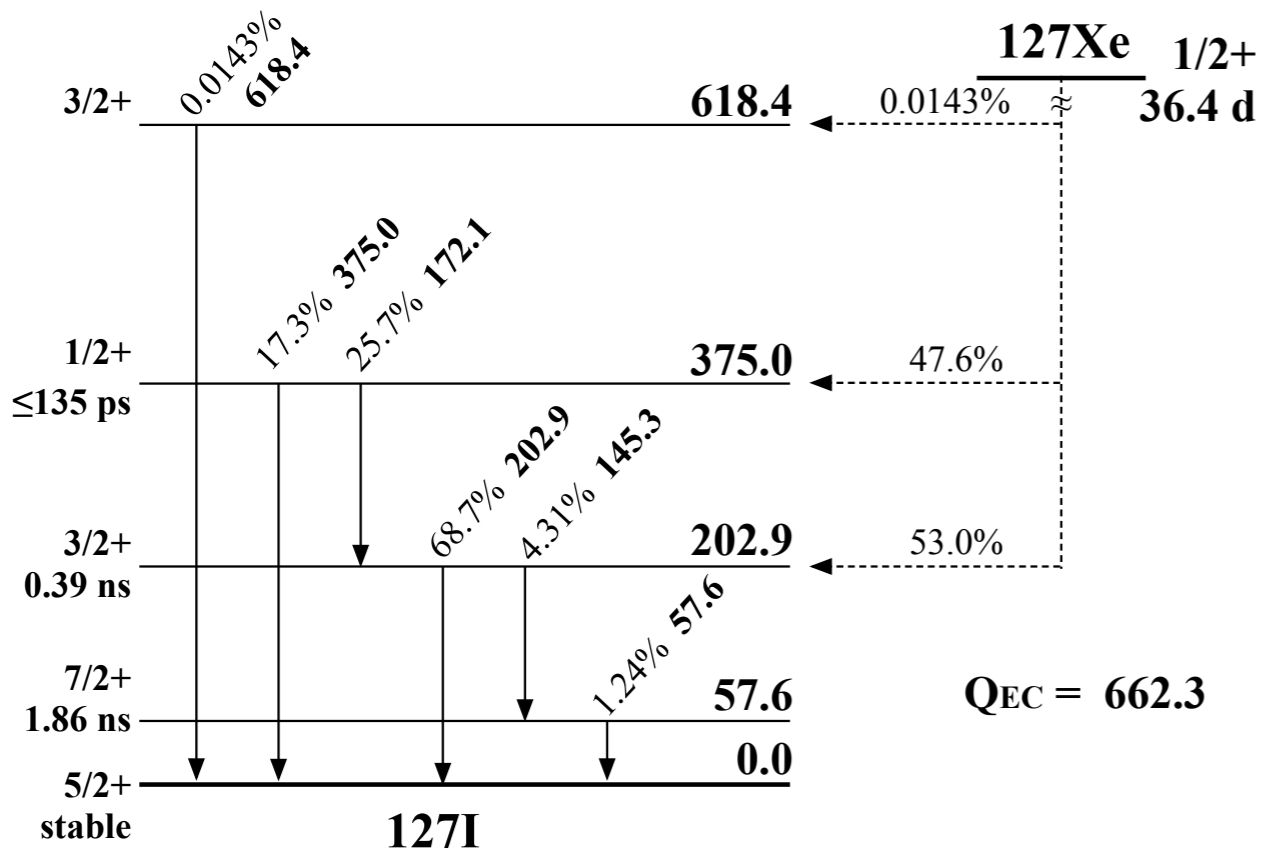
Best-fit mass is  $2.3 \pm 0.2$  keV, so far from 2.8 keV

$^{127}\text{Xe}$  can be produced from cosmogenic activation of Xe at sea level;

Given the short half-life of 36.4 days and the fact that the [xenon gas was underground for O\(1\) years before the operation of XENON1T](#)

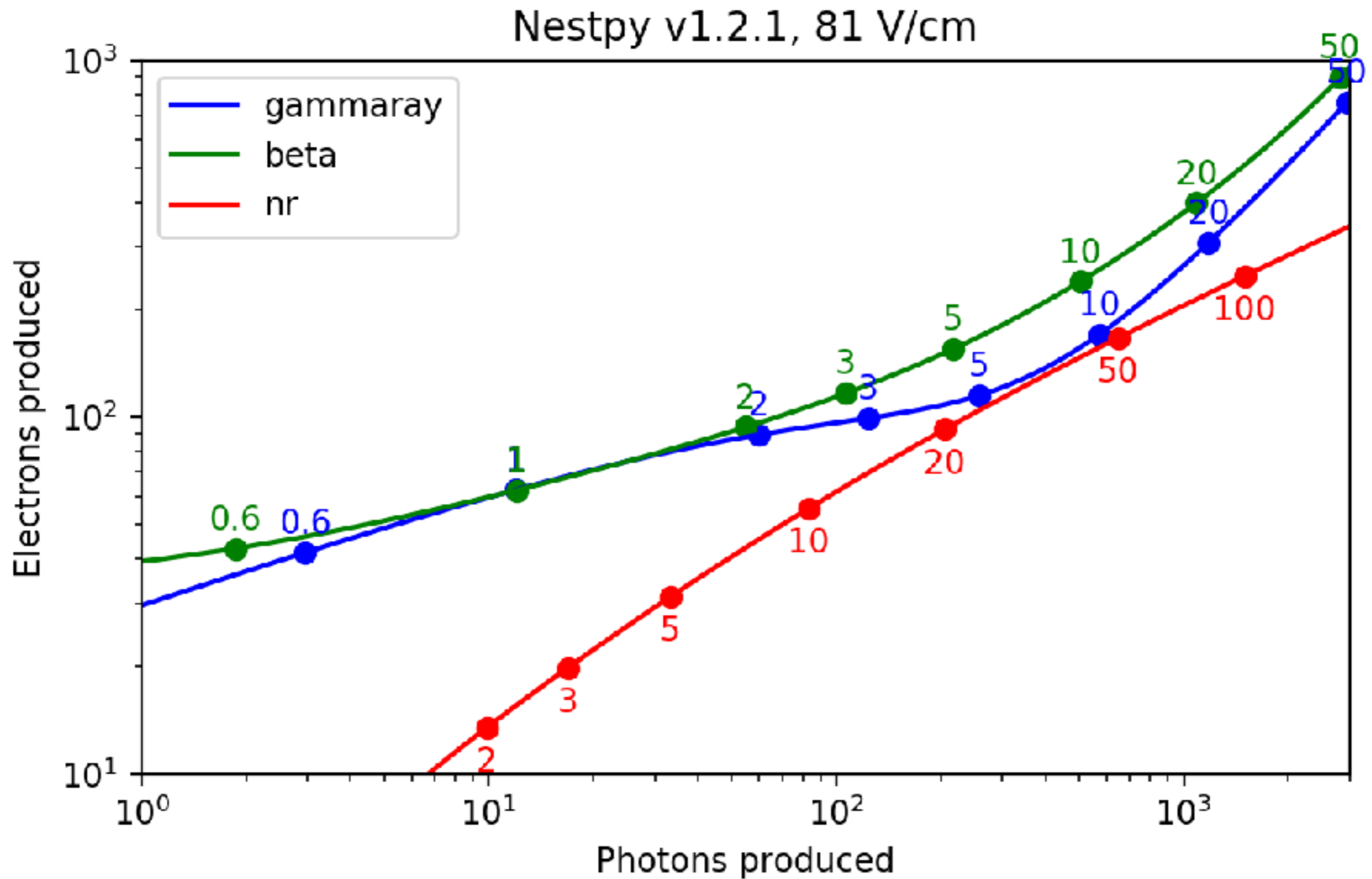
—> already decayed away

Phys. Rev. D 96, 112011 (2017)



Also we did not see high-energy  $\gamma$ s that accompany X-rays

(We are using inner volume for this search, so there is a O(1)cm between FV and the detector wall)



## Exchange effect

$\beta$  electron is created in an atomic orbital of the daughter atom and the atomic electron which was present in the same orbital in the parent atom is ejected to the continuum.

This process leads to the same final state as the direct decay, i.e. one electron in the continuum, and is possible because the nuclear charge changes in the decay.



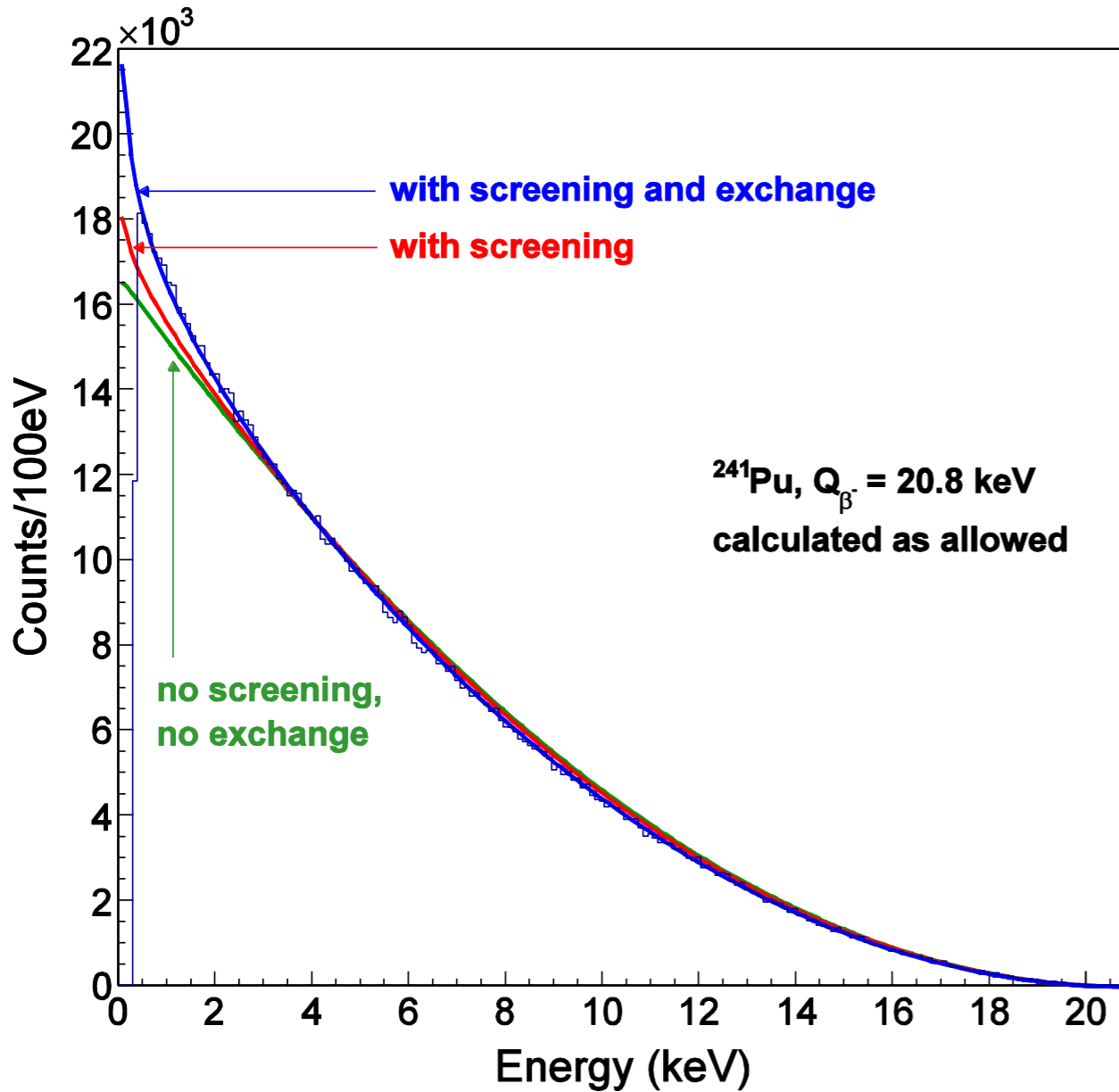
## Screening effect

The electrons in bound states in the atom produce screening of the nuclear charge for the emitted beta particle. This change in electromagnetic field modifies the beta spectrum.

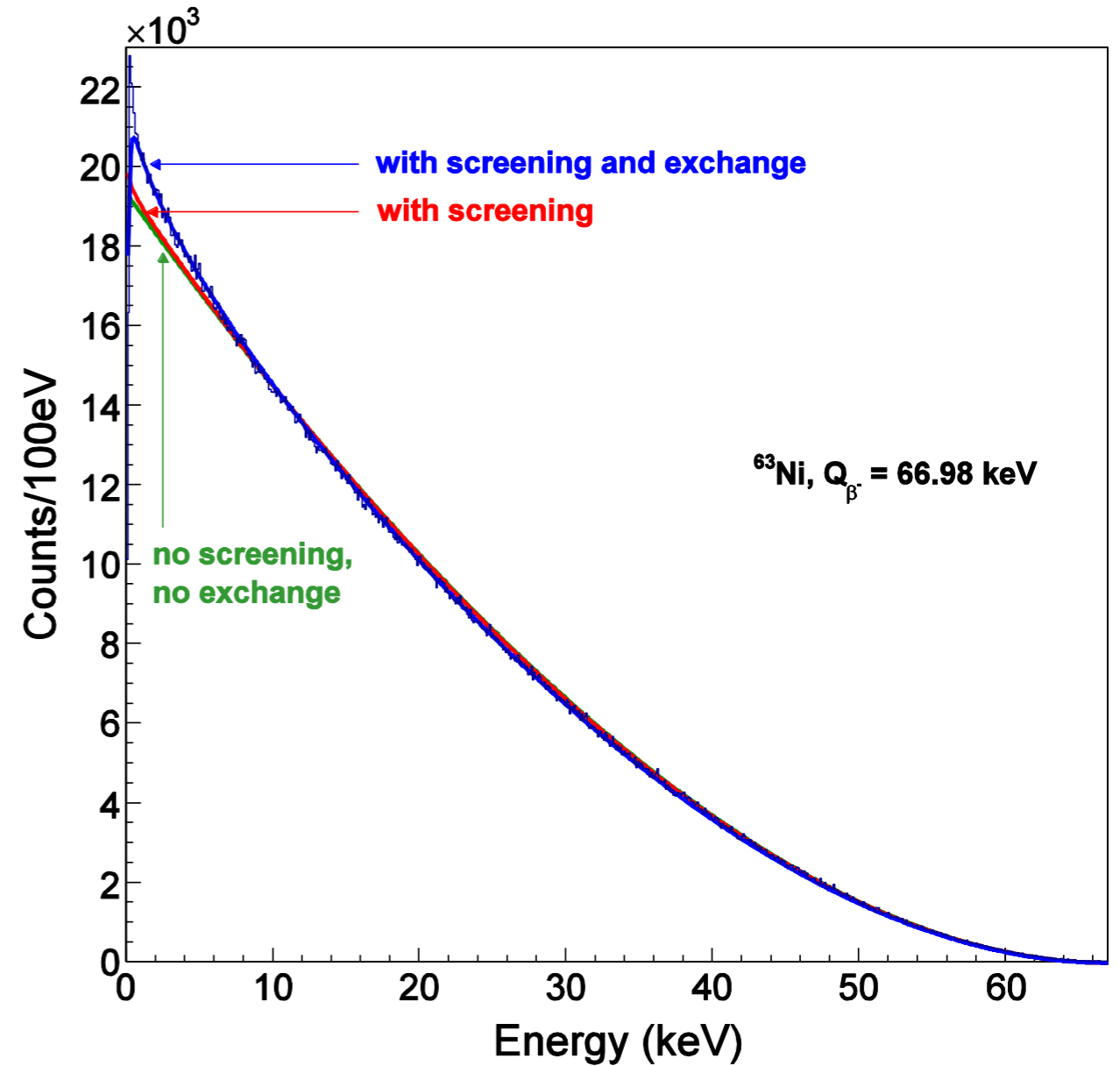
(The atomic screening effect corresponds to the influence of the electron cloud surrounding the daughter nucleus on the  $\beta$  particle wave function)

# $^{241}\text{Pu}$ / $^{63}\text{Ni}$ beta-spectrum

M. Loidl et al., App. Radiat. Isot. 68, 1454 (2010)



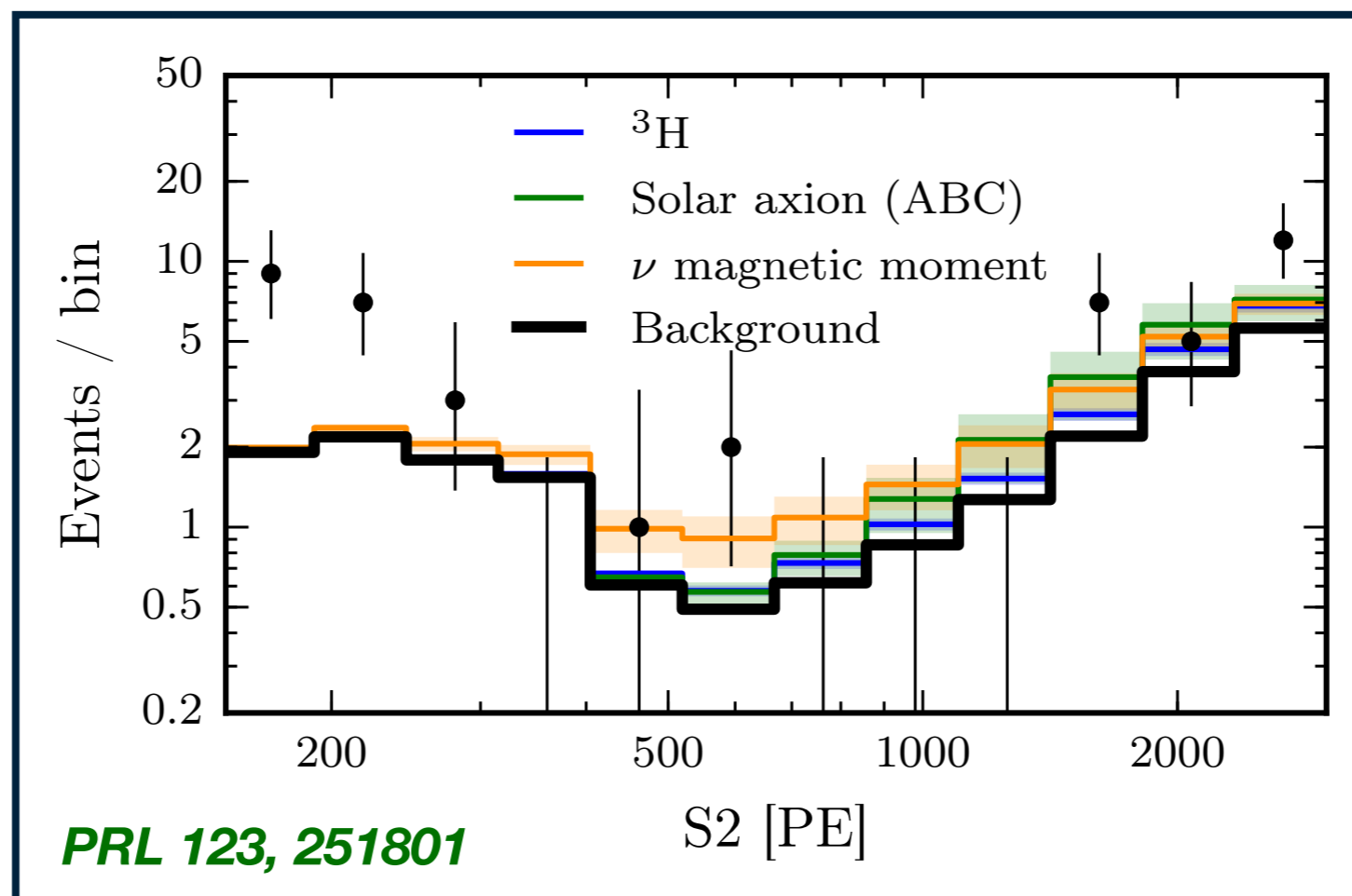
C. Le-Bret, PhD thesis, Université Paris 11 (2012)





# S2-ONLY ANALYSIS

S2-only = No requirement on S1s, allowing for a ~200 eV threshold



S2-only analysis  
allows for a lower  
energy threshold of  
200 eV

$$\mu_\nu < 3.1 \times 10^{-11} \mu_B$$
$$g_{ae} < 4.8 \times 10^{-12}$$
$$R_{\text{H3}} < 2256 \text{ events/t/y}$$

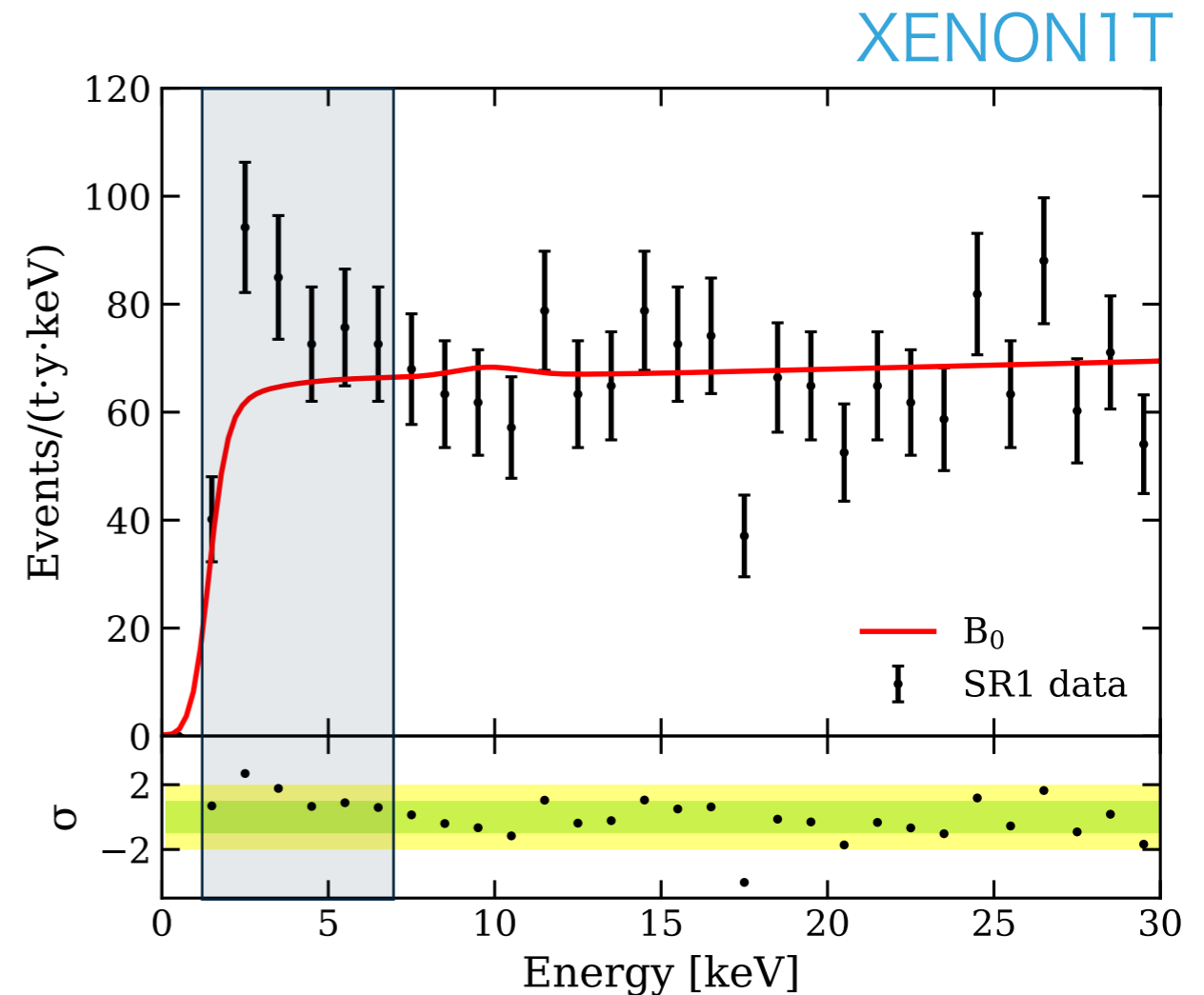
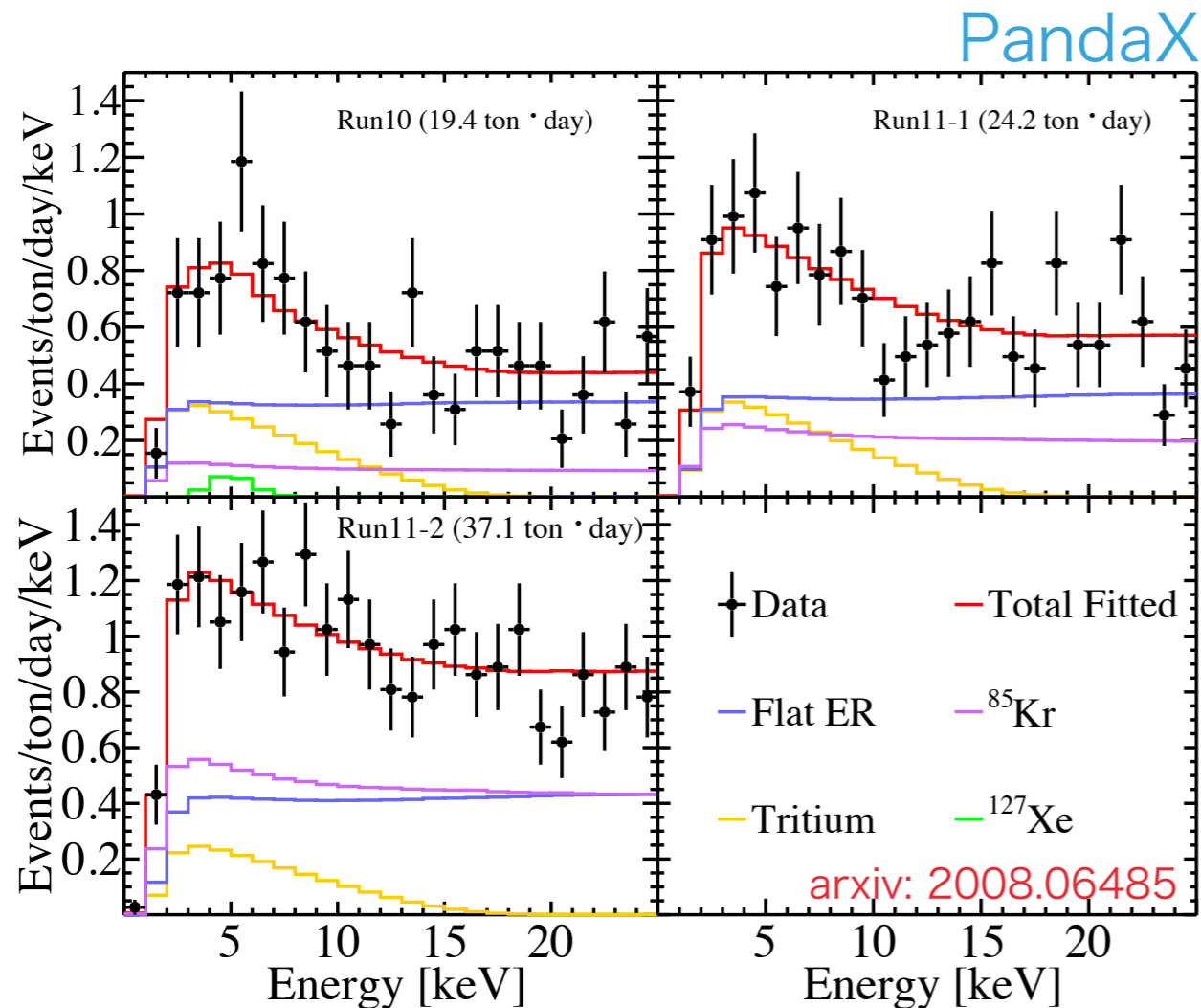
**consistent with this work  
for all 3 hypotheses**

larger upper limits, S2-only analysis is not  
sensitive to the excesses we found

# Other Experiments: PandaX?

Exposure: 101 ton-days (XENON1T: 236 ton-days)

BG level: ~5 times higher BG than in XENON1T



Note that PandaX experiment injected Tritiated Methane (CH<sub>3</sub>T) directly into the detector for low-E calibration, but they could not completely remove it with distillation.

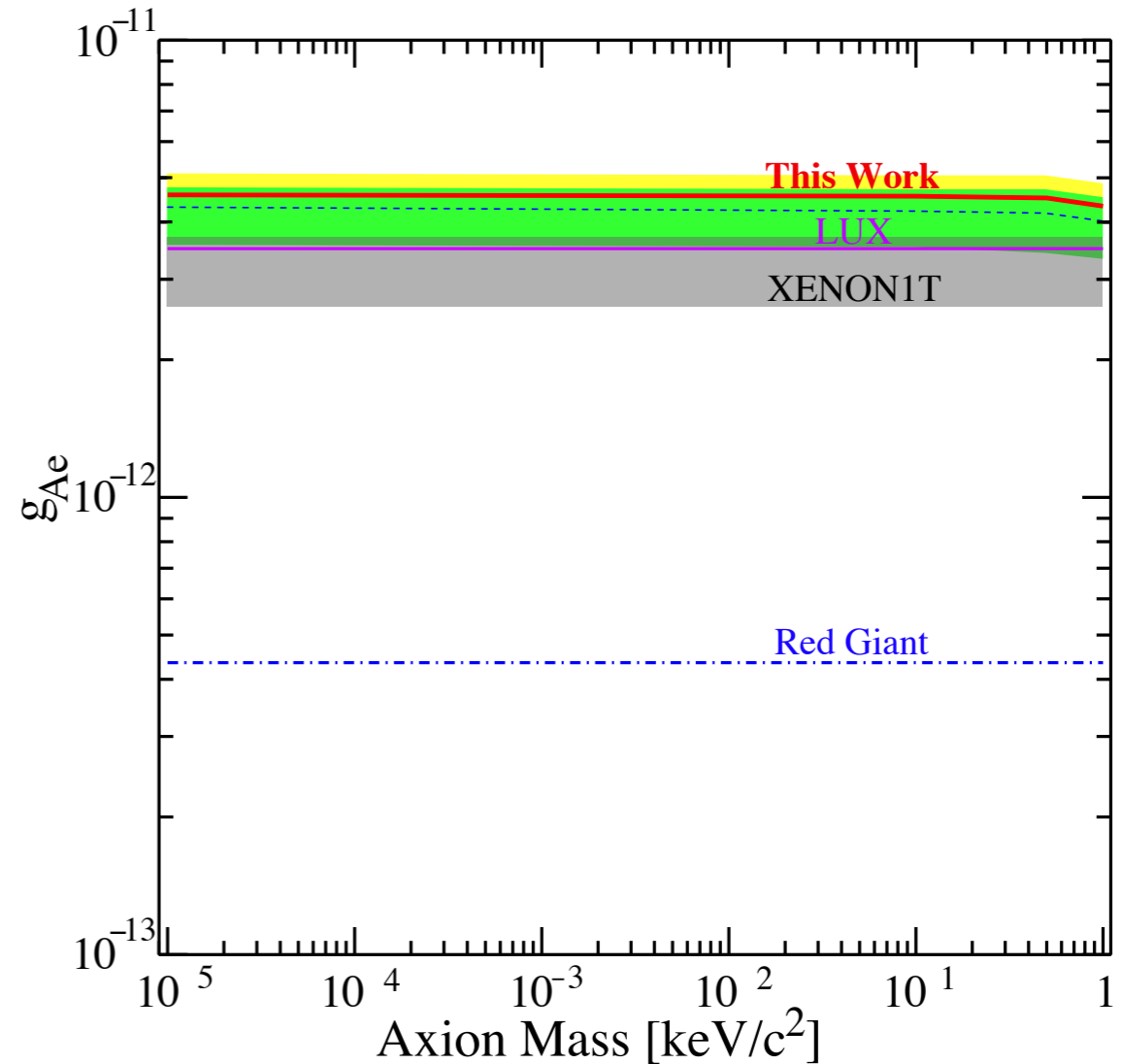
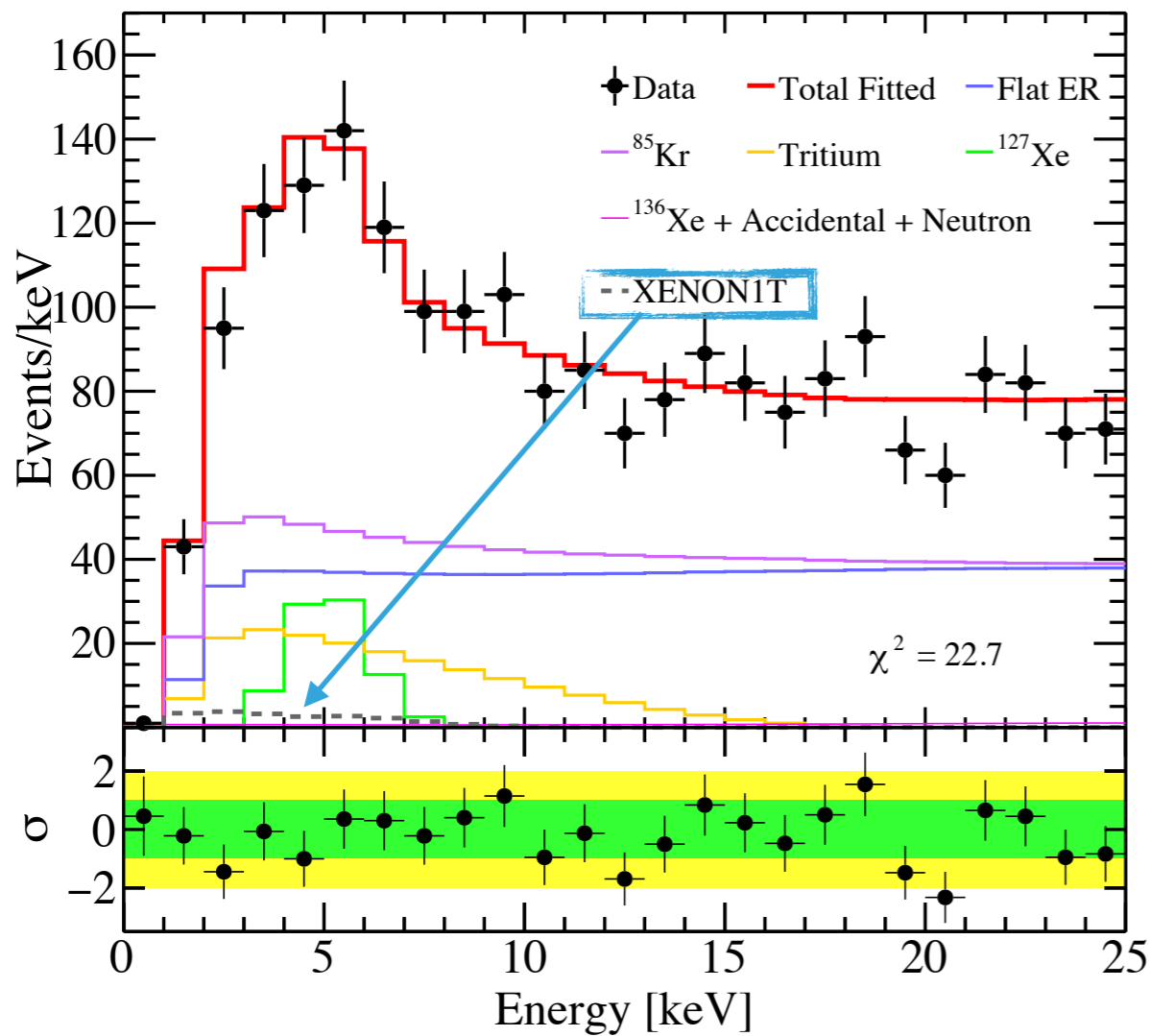
Resulting tritium concentration is ~0.04 uBg/kg ( $\sim 5 \times 10^{-24}$  mol/mol in xenon)

Best-fit for tritium hypothesis in XENON1T:  $\sim 6 \times 10^{-25}$  mol/mol (we did not inject CH<sub>3</sub>T!)

# Other Experiments: PandaX?

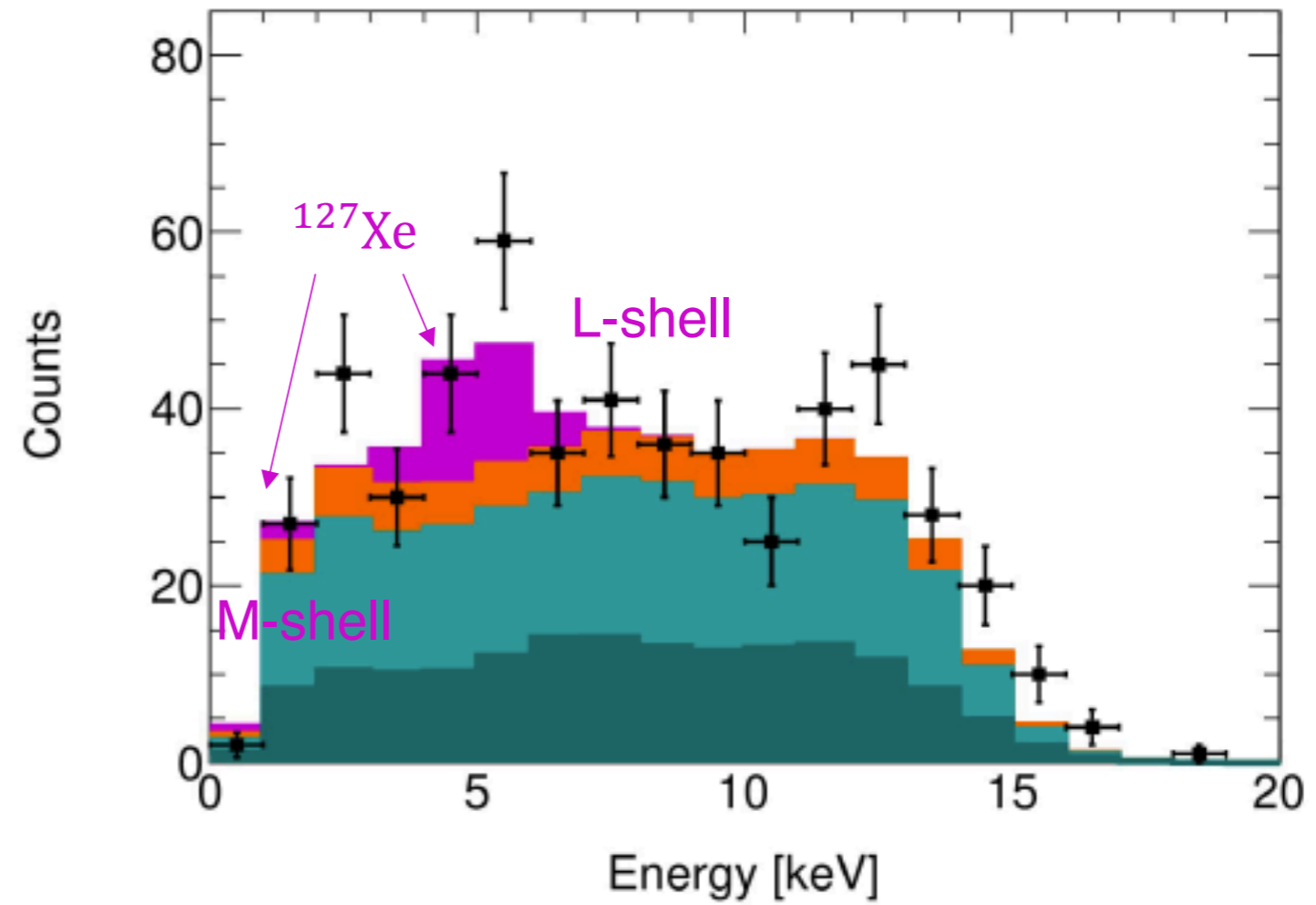
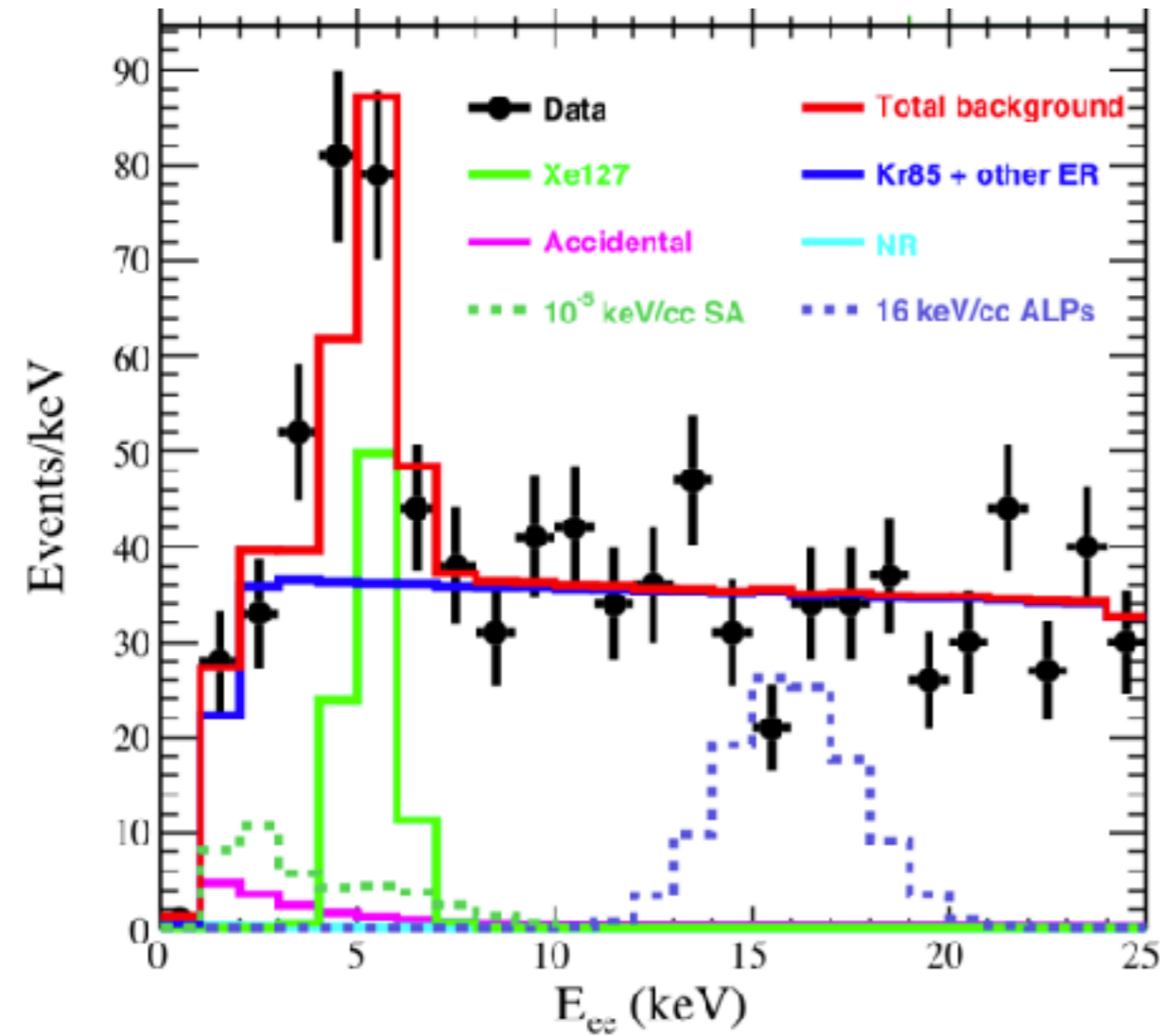
Exposure: 101 ton-days (XENON1T: 236 ton-days)

BG level: ~5 times higher BG than in XENON1T



Expected excess assuming the best fit signal strength from XENON1T is compatible with their data within uncertainties, but their data are also consistent with background- only hypothesis

→ No conclusion because of lower statistics & higher BG



LUX Collaboration, PRL 118, 261301

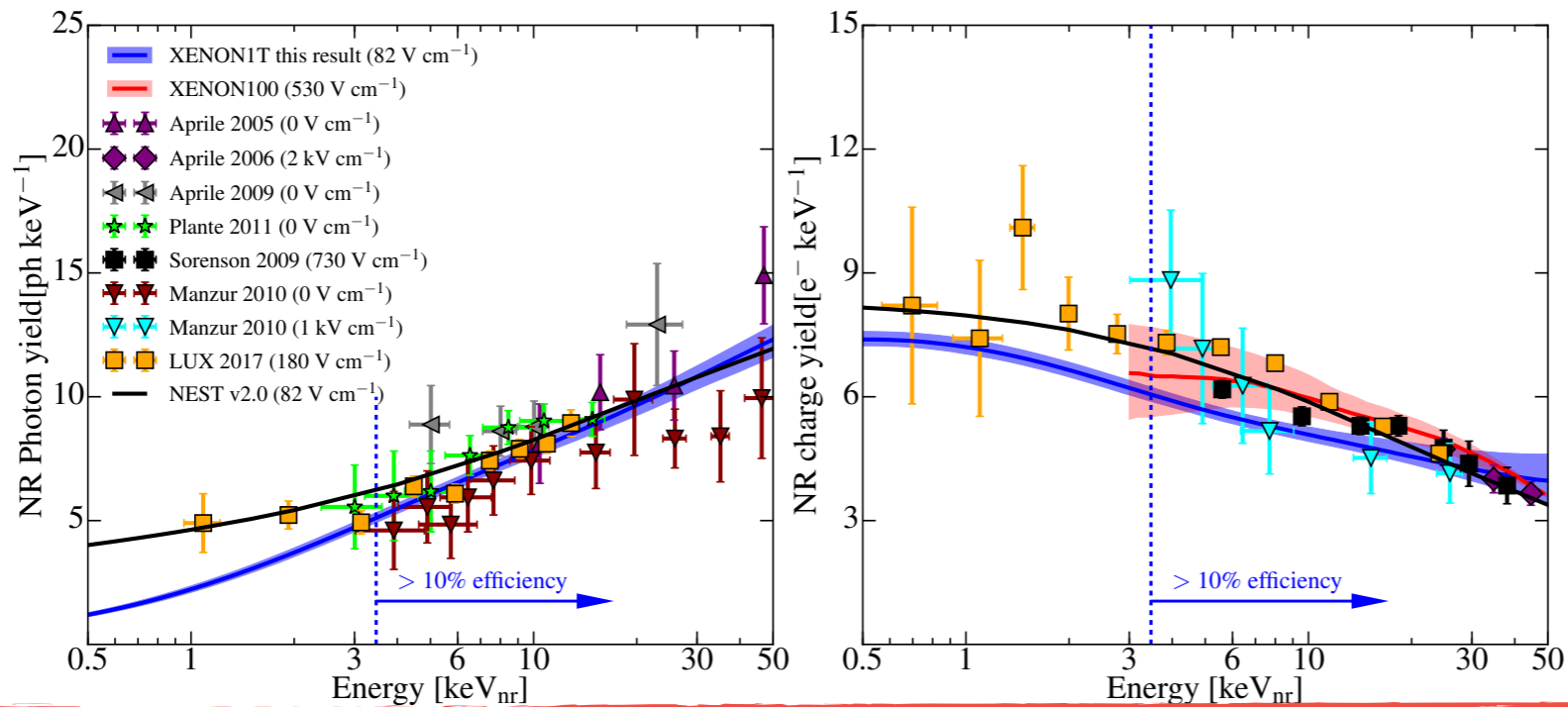
low-z-origin  $\gamma$  rays (dark green),

other  $\gamma$  rays (light green),

$^{85}\text{Kr}$  or Rn-daughter contaminants in the liquid xenon undergoing  $\beta$  decay (orange)

x rays due to  $^{127}\text{Xe}$  (purple).

NR

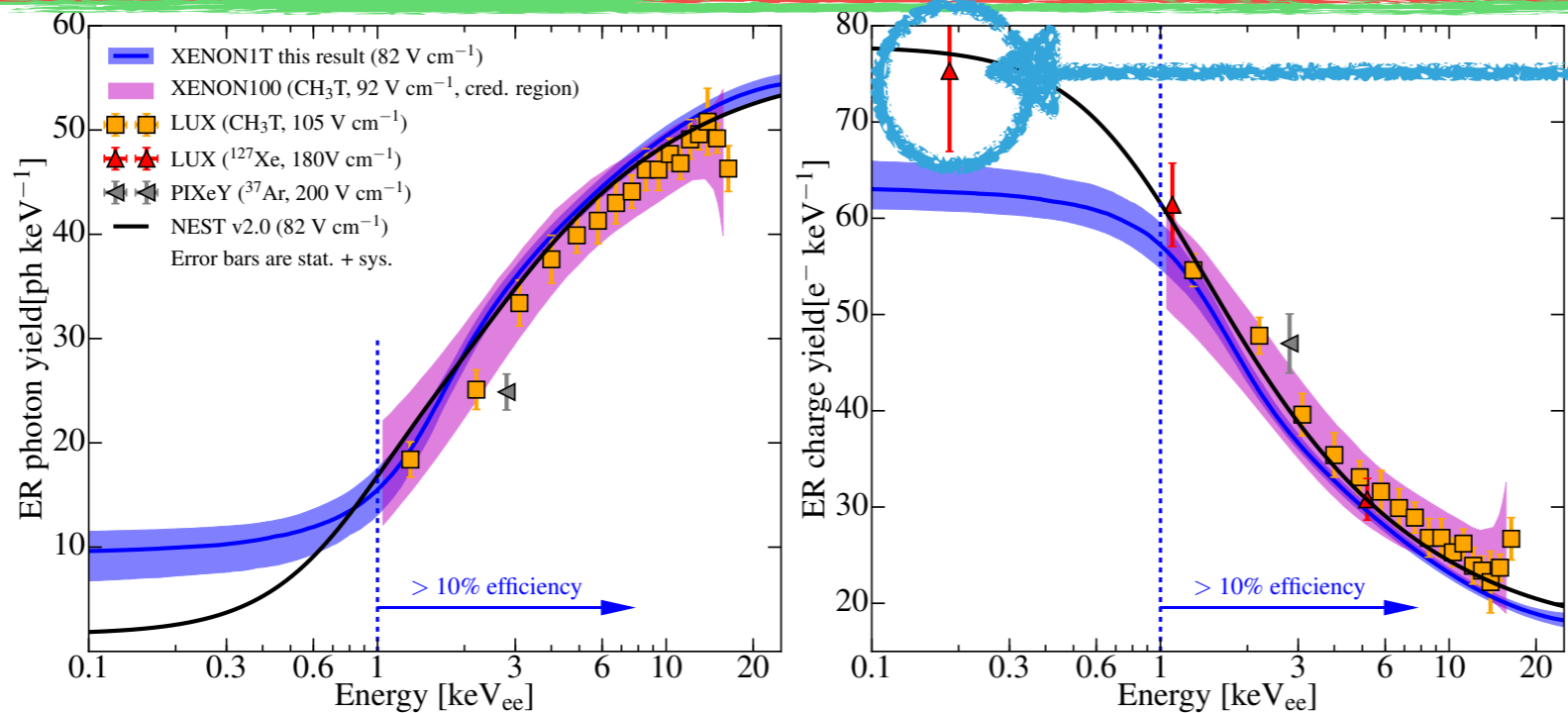


New results!

[arXiv: 1902.11297](https://arxiv.org/abs/1902.11297)

Estimated using  $^{220}\text{Rn}$  (ER) and  $^{241}\text{Am}$  / Neutron Generator (NR) calibration data.

ER



For charge yield, there is a 0.186 keV<sub>ee</sub> ( $^{127}\text{Xe}$ ) measurement from LUX

Phys. Rev. D **96**, 112011

1e<sup>-</sup> produces ~30 PEs.  
Photo-detection eff ~ 10%

- ▶ 10% efficiency for S1&S2 analysis corresponds to energy deposit of ~1 keV for ER and ~4 keV for NR
- ▶ ER signals can significantly enhance the detection efficiency → higher sensitivity for low-mass WIMPs.
- ▶ Moreover, S2-only analysis can also decrease the threshold from 1 keV<sub>ee</sub> to 0.186 keV<sub>ee</sub>

$$m_a \simeq \frac{6 \times 10^6 \text{ GeV}}{f_a} \text{ eV}/c^2$$

DFSZ: two Higgs doublets model  
couplings to leptons at tree level

KSVZ: heavy quark model  
couplings to leptons only at loop level

**quarks/electrons related by Beta**

**photons/electrons related by E/N**

Dine-Fischler-Srednicki-Zhitnitsky (DFSZ)

Kim-Shifman-Vainshtein- Zhakharov (KSVZ)

**axion-photon coupling  
same for both models**

- relative contributions from each component can allow to distinguish between models (Primakoff dominates in KSVZ models); can also constrain  $\beta_{\text{DFSZ}}$
- nuclear transition contribution always relatively small

the ABC flux is dominant in DFSZ models, while the Primakoff flux is dominant in KSVZ models.

# QCD axion models

$$m_a \simeq \frac{6 \times 10^6 \text{ GeV}}{f_a} \text{ eV}/c^2$$

DFSZ: two Higgs doublets model  
couplings to leptons at tree level  
quarks/electrons related by Beta

$$g_{ae} = \frac{m_e}{3f_a} \cos^2 \beta_{\text{DFSZ}}$$

$$\tan(\beta_{\text{DFSZ}}) = \left( \frac{X_u}{X_d} \right)^{1/2}$$

Dine-Fischler-Srednicki-Zhitnitsky (DFSZ)

KSVZ: heavy quark model  
couplings to leptons only at loop level  
photons/electrons related by E/N

$$g_{ae} = \frac{3\alpha^2 N m_e}{2\pi f_a} \left( \frac{E}{N} \ln \frac{f_a}{m_e} - \frac{2}{3} \frac{4+z+w}{1+z+w} \ln \frac{\Lambda}{m_e} \right)$$

$$z = m_u/m_d, \quad m_u/d$$

$$w = m_u/m_s$$

Kim-Shifman-Vainshtein- Zhakharov (KSVZ)

$$g_{a\gamma} = \frac{\alpha}{2\pi f_a} \left( \frac{E}{N} - \frac{2}{3} \frac{4+z}{1+z} \right)$$

axion-photon coupling  
same for both models

- relative contributions from each component can allow to distinguish between models (Primakoff dominates in KSVZ models); can also constrain  $\beta_{\text{DFSZ}}$
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# Solar axion/ALP flux

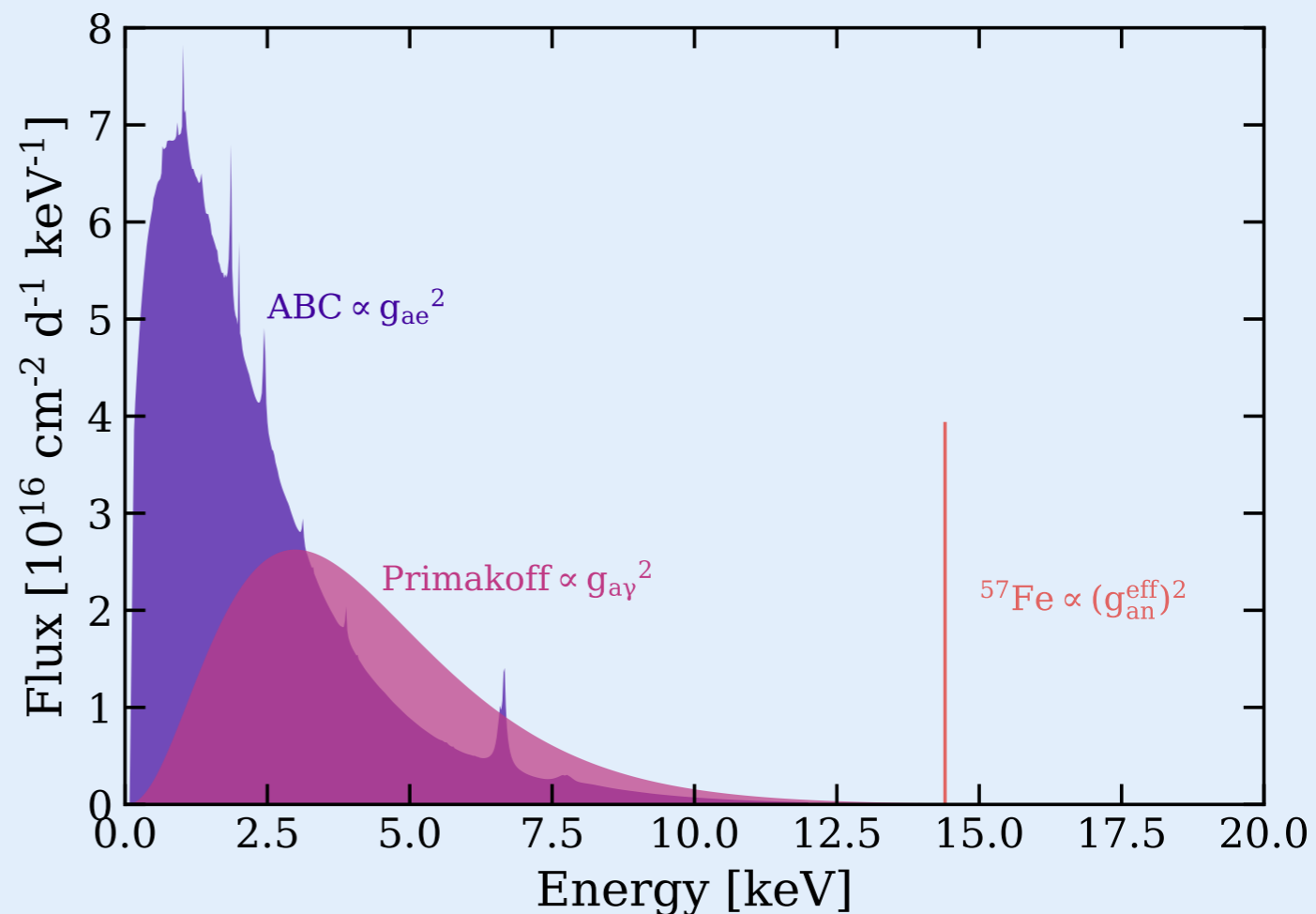
## ABC

(Atomic recombination and deexcitation, Bremsstrahlung and Compton)

$$\frac{d\Phi_a}{d\omega} = \frac{1}{4\pi R_{\text{Earth}}^2} \int_{\text{Sun}} dV \frac{4\pi\omega^2}{(2\pi)^3} \Gamma_a^{\text{P}}(\omega)$$

## Primakoff:

$$\frac{d\Phi_a^{\text{Prim}}}{dE} = \left( \frac{g_{a\gamma}}{\text{GeV}^{-1}} \right)^2 \left( \frac{E_a}{\text{keV}} \right)^{2.481} e^{-E_a/1.205/\text{keV}} \times 6 \times 10^{30} \text{ cm}^{-2} \text{ s}^{-1} \text{ keV}^{-1}.$$



## Fe-57 nuclear transition:

$$\Phi_a^{57\text{Fe}} = \left( \frac{k_a}{k_\gamma} \right)^3 \times 4.56 \times 10^{23} (g_{an}^{\text{eff}})^2 \text{ cm}^{-2} \text{ s}^{-1}.$$

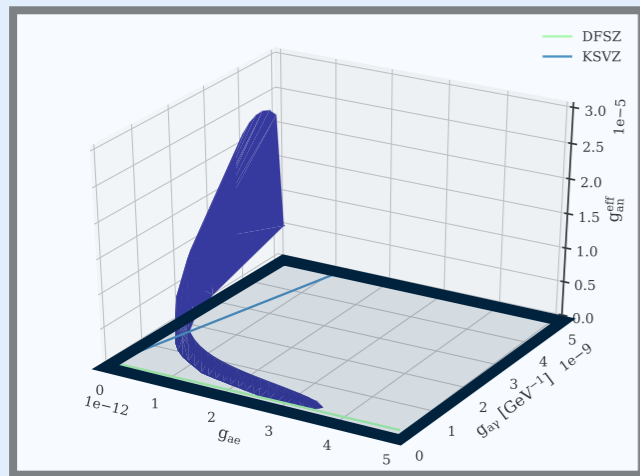
$k$ : momenta of axion, photon



# Statistical inference

3D confidence volume (90% C.L.)

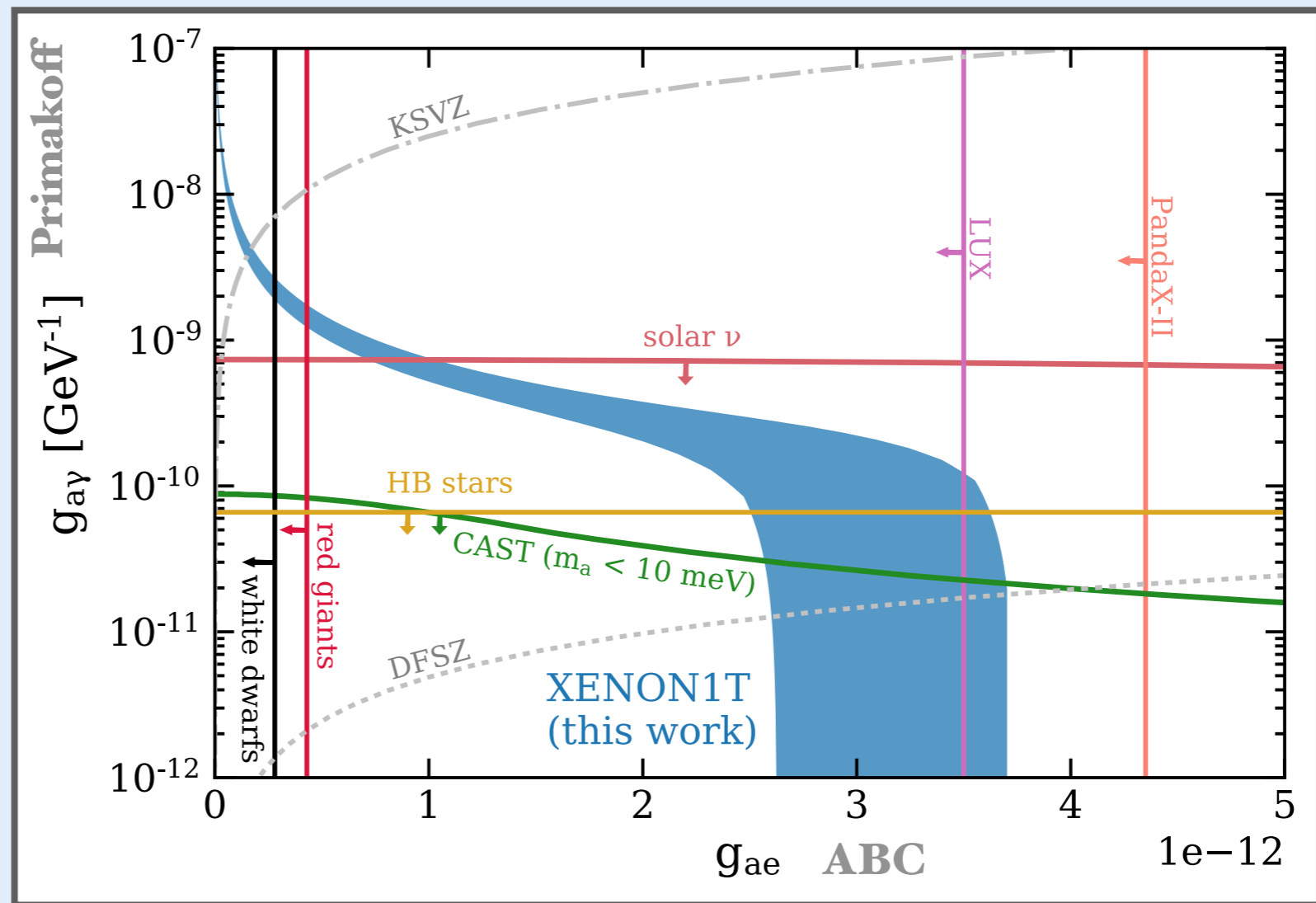
Projected onto 2D regions



$$g_{ae} < 3.7 \times 10^{-12}$$

$$g_{ae} g_{an}^{eff} < 4.6 \times 10^{-18}$$

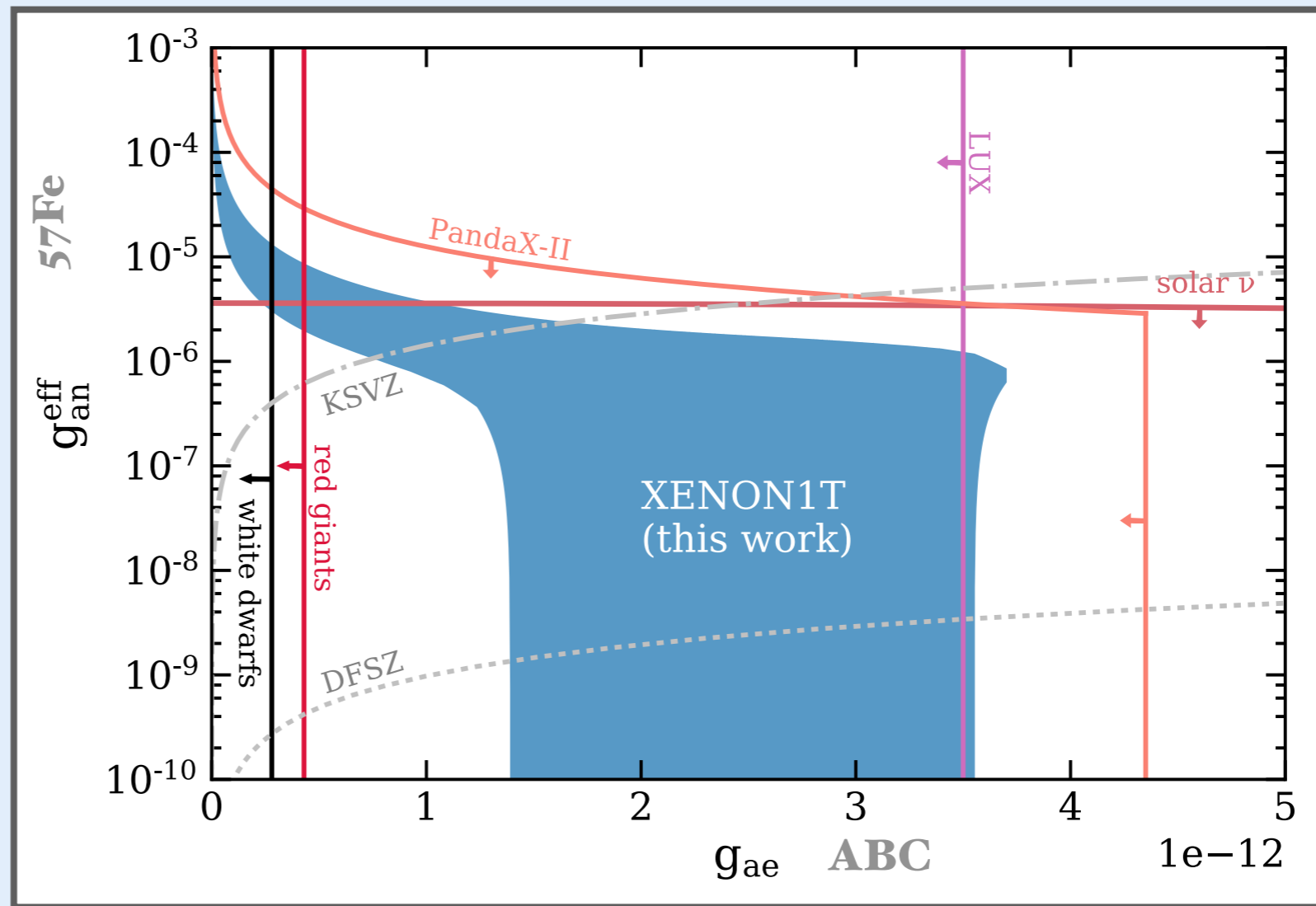
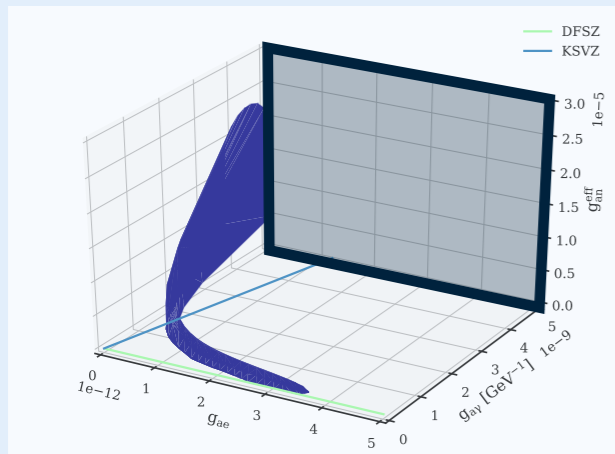
$$g_{ae} g_{a\gamma} < 7.6 \times 10^{-22} \text{ GeV}^{-1}$$



Discrepancy with astrophysical constraints from stellar cooling  
(*arXiv:2003.01100*)

# Statistical inference

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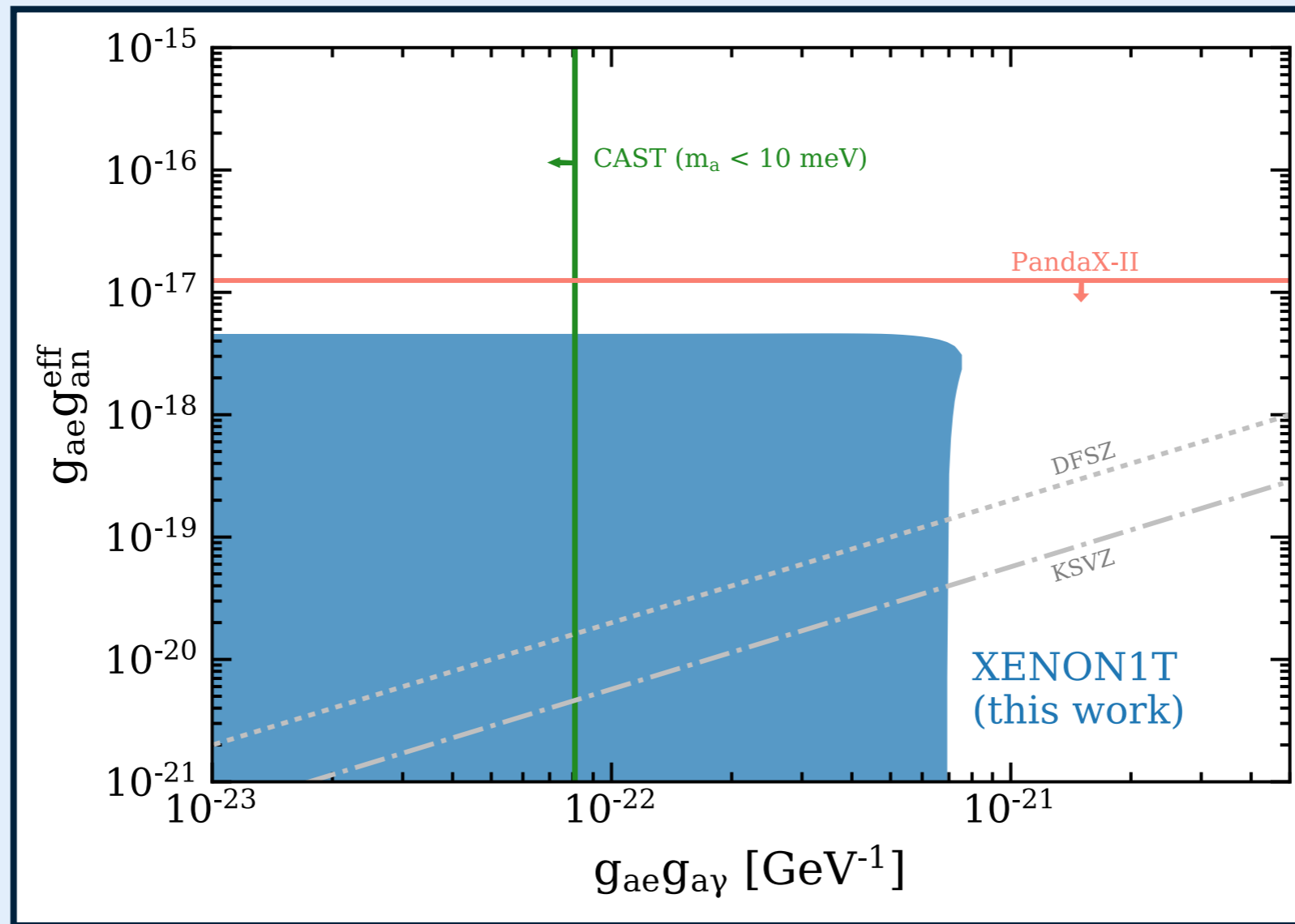
$$g_{ae} g_{a\gamma} < 7.6 \times 10^{-22} \text{ GeV}^{-1}$$

Poor fit for small ABC rate

Discrepancy with astrophysical constraints from stellar cooling  
(*arXiv:2003.01100*)

# Solar axion/ALP

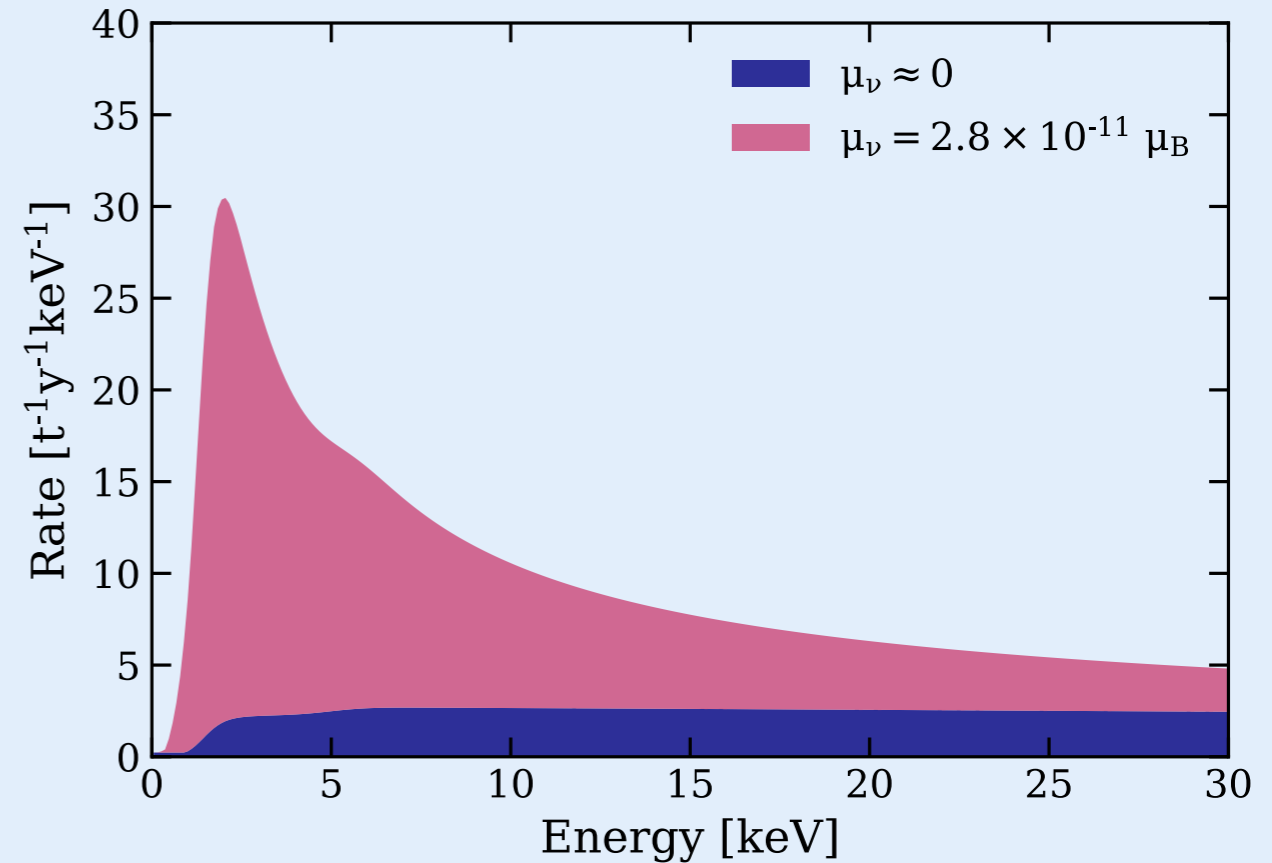
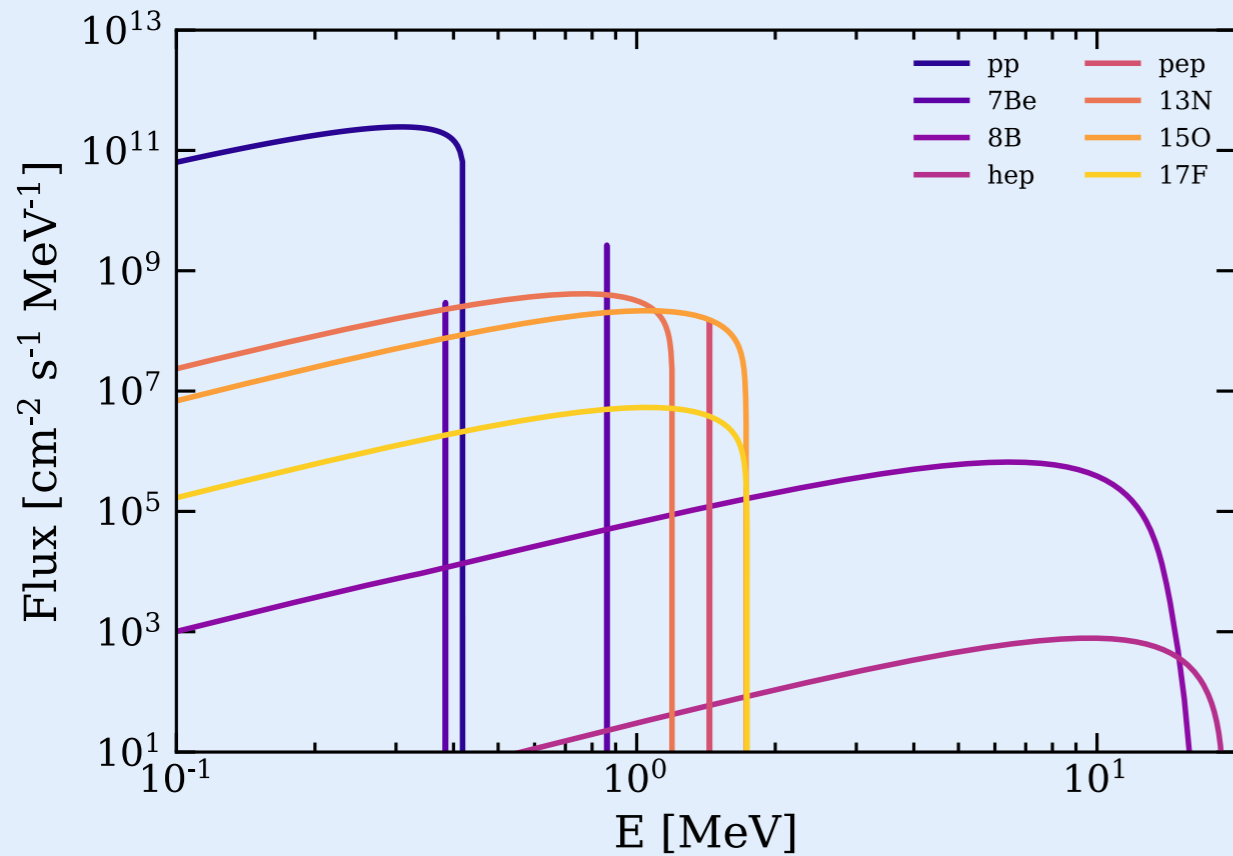
3D confidence  
volume (90%  
C.L.)  
Projected onto  
2D regions



Primakoff and  $^{57}\text{Fe}$  components can be absent *if* the ABC component is present

No statistical significance for Primakoff or  $^{57}\text{Fe}$  on their own

# Neutrino magnetic moment



$$\frac{d\sigma_\mu}{dE_r} = \mu_\nu^2 \alpha \left( \frac{1}{E_r} - \frac{1}{E_\nu} \right)$$

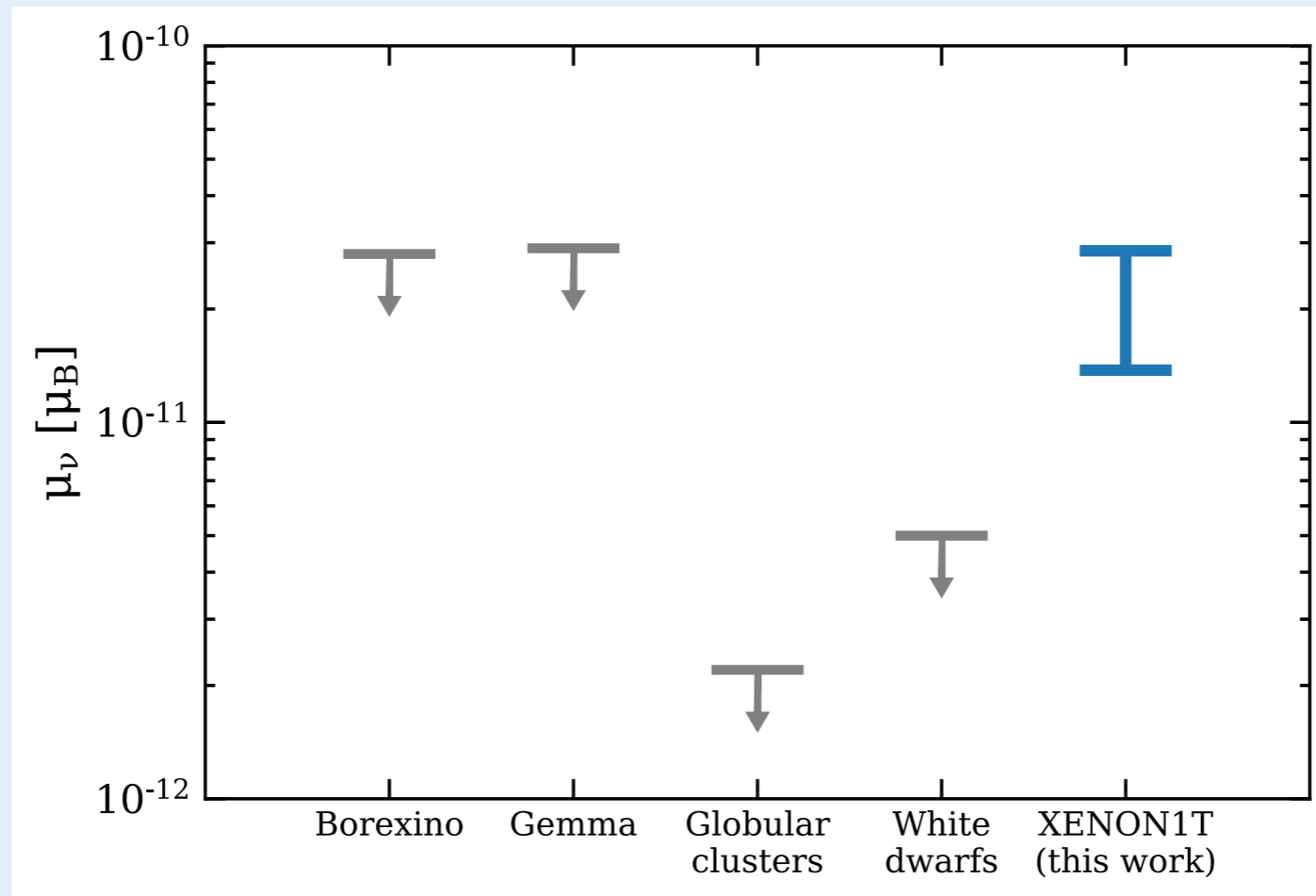
In the (extended) SM:

$$\mu_\nu \approx 3 \times 10^{-19} \left( \frac{m_\nu}{\text{eV}} \right) \mu_B$$

A larger value would imply new physics, and possibly solve Dirac vs Majorana.

$$\mu_\nu \gtrsim 10^{-15} \mu_B \rightarrow \text{Majorana fermion}$$

# Neutrino magnetic moment



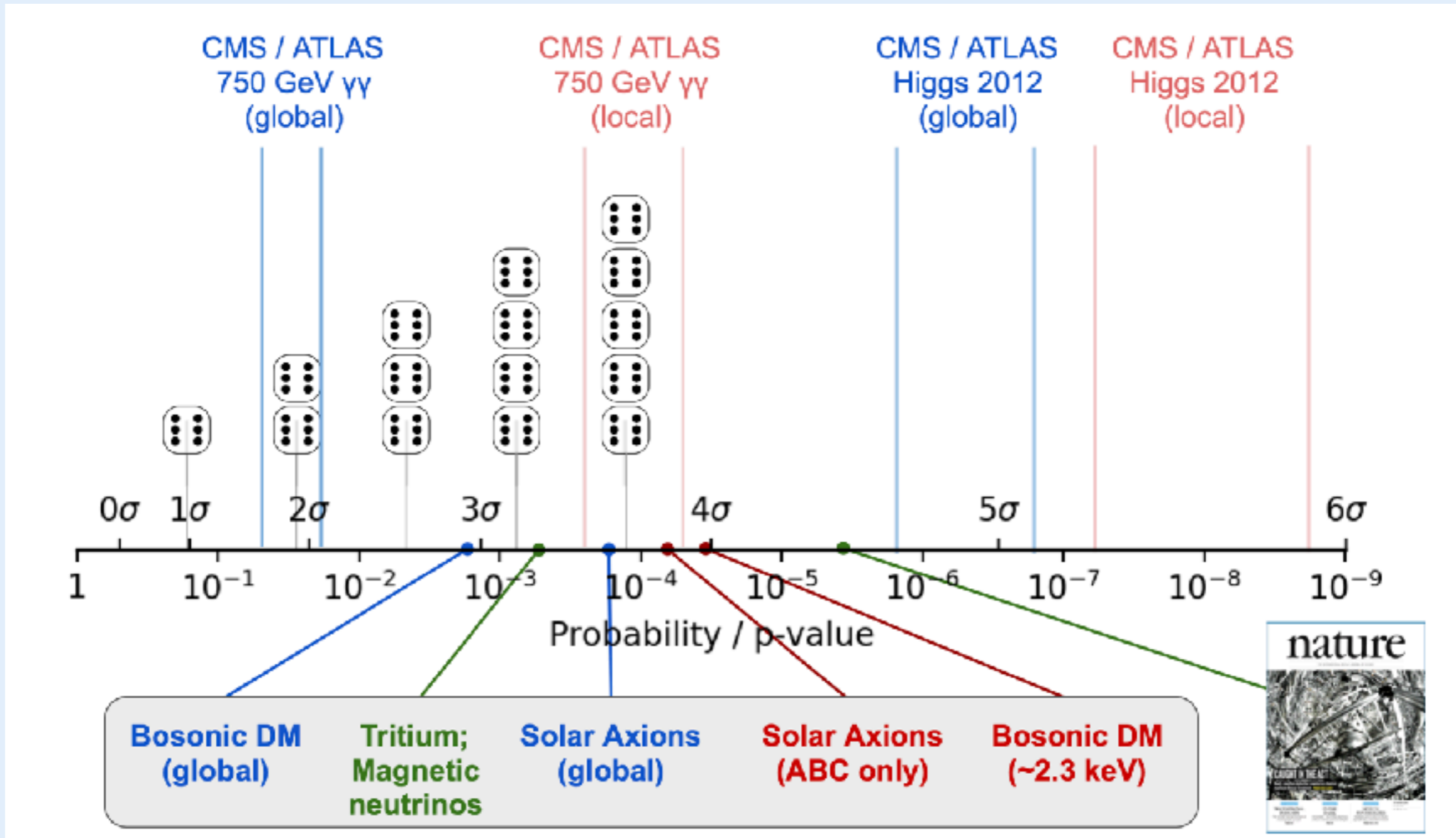
$$\mu_\nu \in (1.4, 2.9) \times 10^{-11} \mu_B$$

# Background Fit

Component		Expected Events	Fitted Events	Constant in time? (shared across partitions)
$^{214}\text{Pb}$		(3450, 8530)	7480 +/- 160	YES
$^{85}\text{Kr}$		890 +/- 50	773 +/- 80	NO
$^{136}\text{Xe}$		2120 +/- 210	2150 +/- 120	YES
$^{133}\text{Xe}$		3900 +/- 410	4009 +/- 85	NO
$^{131}\text{Xe}$		23760 +/- 640	24270 +/- 150	NO
$^{83\text{m}}\text{Kr}$		2500 +/- 250	2671 +/- 53	NO
Materials		323 (fixed)	323 (fixed)	YES
Solar neutrino		220.7 +/- 6.6	220.8 +/- 4.7	YES
$^{124}\text{Xe}$	KK	125 +/- 50	113 +/- 24	YES
	KL	38 +/- 15	34.0 +/- 7.3	YES
	LL	2.8 +/- 1.1	2.56 +/- 0.55	YES
$^{125}\text{I}$	K	79 +/- 33	67 +/- 12	NO
	L	15.3 +/- 6.5	13.1 +/- 2.3	NO
	M	3.4 +/- 1.5	2.94 +/- 0.50	NO

*unconstrained in the fit*

# Probabilities



slide credit: Jelle Aalbers